

**POLYMERISATION KINETICS AND OPTICAL PHENOMENA OF
PHOTOACTIVE DENTAL RESINS**

By

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SYNOPSIS

Globally, the use of resin based composites (RBCs) is increasing but several shortcomings remain. This may be related to problems associated with incomplete conversion (40-70%), polymerisation shrinkage (1-4% by volume) and the associated stress generated at the tooth/restoration interface. Additionally, the increased number of technique sensitive incremental steps required to fill relatively large cavities is due to inefficient light transmission at depths greater than 2 mm. Accordingly, the change in optical properties and the setting reaction of RBC materials is not well understood. An interesting approach to control the setting reaction is by the application of the exposure reciprocity law. Therefore the current investigation demonstrates the applicability of this law with regards to monomer composition, filler percentage and photoinitiator type on degree of conversion with the aim of a better understanding of the setting reaction. Although appropriate resin matrix chemistry and photoinitiator chemistry may provide potential for reduced curing time, inadequate cure depths remain without appropriate filler adaptation.

The development of techniques that will allow dynamic monitoring of optical and physical change during cure will further aid material development with the goal of improved depths of cure and will allow such restorations to be cured with one 'shot'. The current investigation has demonstrated the use of several analytical techniques (FT-IR spectroscopy, UV-Vis Spectroscopy and low coherence interferometry) which will aid such developments. UV-Vis spectroscopy and low coherence interferometry may have some significance towards a better understanding of the dynamic changes of optical properties of RBC materials. UV-Vis Spectroscopy may be used to monitor the decomposition of the photoinitiator which will improve light transmission and low coherence interferometry can be used to monitor both refractive index and physical thickness change. The current study demonstrated the complexity of optical phenomena within RBCs which are affected by material composition as well as cavity dimensions.

Whilst research continues to develop a novel RBC with reduced shrinkage and improved depths of cure, there is currently no commercially available solution to such problems. Consequently a better understanding of the setting reaction, optical properties and physical properties will aid material development.

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DEDICATED TO MY FAMILY

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DECLARATION

I hereby declare that this thesis is my own work and effort. Where other sources of information have been used, they have been acknowledged.

This thesis consists of two experimental Chapters with six sections in the form published and submitted papers with contributions from the following people other than the primary supervisors:

Chapter 2.1: High Irradiance Curing and the Applicability of the Exposure Reciprocity Law in Commercial Dental Resin-Based Materials

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Chapter 3.2: Dynamic Monitoring of Refractive Index Change Through Photoactive Resins.

Tomlins P – Named Author/ Static Refractive Index Measurements

LIST OF MAIN ABBREVIATIONS

Δb : Change in *b value* of the CIELAB system

ΔE : Colour change value

ΔT : Change in temperature

ΔT^{Max} : Maximum change in temperature

DC^{Max} : Maximum degree of conversion

R_p^{Max} : Maximum rate of polymerisation

R_p : Rate of polymerisation

ANOVA: Analysis of variance

BHT: Butylated hydroxy toluene

Bis-GMA: bisphenol A glycidyl dimethacrylate

CQ: Camphoroquinone

DC: Degree of conversion

DMAEMA: Di-methyl-aminoethyl methacrylate

DOC: Depth of cure

FT-NIRS: Fourier-transform near infrared spectroscopy

HF: High filled

LCI: Low coherence interferometry

LF: Low Filled

NIR: Near infra-red

PI: Photoinitiator

RBCs: Resin based composites

RI: Refractive Index

TEGDMA: triethylene glycol dimethacrylate

UF: Unfilled

UV-Vis: Ultra-Violet-Visible

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CHAPTER 1

Introduction and Literature Review

1.0 Historical perspectives

The need for the artificial replacement of missing body parts has existed for thousands of years in order to improve and maintain primary health as-well as primary oral health. The latter includes prevention and management of dental caries and periodontal disease and improvements to the functionality and appearance of missing, damaged or worn teeth. Attempts at filling spaces created by tooth loss or tooth alterations has been documented as far back as the early century's in B.C were natural teeth were ligated with gold or silver wire, and detached teeth or teeth carved from other materials such as animal teeth, ivory and bone were connected by wire and ligated to existing teeth (Laney, 1997). However, such materials were utilised with limited success due to the instability of the materials in the oral cavity due to the corrosive effects of human saliva and the tendency for ivory and bone to become stained over time (Kelly et al 1996).

The development of amalgam in the 18th and 19th century revolutionised restorative dentistry and gained widespread acceptance shortly after (Berry et al., 1998). Unfortunately, controversial issues regarding the use of mercury and the release of mercury vapour during mastication arose and its safety came into question, despite a lack of consensus (Osborne, 1992; Dodes, 2001; Osborne, 2004; Mutter et al., 2004; Sadowsky, 2006). More recently however, several Scandinavian countries (Norway, Denmark and Sweden) have banned the use of dental amalgam amongst fears of mercury toxicity and its environmental impact (Jones, 2008). However a significant drift from amalgam to resin based composites (RBCs) in several countries recently (UK, USA and Finland) is a result of clinical indication and the patients aesthetic demand and as a result of dentists seeking less invasive treatments (Forss and Widström, 2001; Burke et al., 2003; Haj-Ali et al., 2005; Gilmour et al., 2007) rather than environmental and toxicity issues. Consequently, the use of resin-based composites has exponentially grown over several decades and today it is the direct restorative material of choice for many dentists across the world. For example, by 2005 in the USA, resin composite was the most frequently used direct restorative material for posterior restorations (Hajj-Ali et al., 2005). The exponential growth is evident in studies performed in the UK, where in 2001,

49% of dentists questioned never or rarely used resin composites for posterior restorations (Burke et al., 2003), whereas in 2007 the number increased to a majority of restorations (Gilmour et al., 2007). Although these material types are extensively used in clinical practice there remain many unanswered questions concerned with intrinsic material properties and associated setting parameters such as shrinkage of the resin matrix, photoinitiator absorption, optical phenomena related to resin and filler interaction. This investigation explores such material and system deficiencies in an attempt to improve our understanding of RBC technology, which may hopefully assist future material development.

1.1 Resin based composites

A composite structure is a mixture of at least two phases with differing properties of one kind or another, which are intermediate in many senses to those of the components. This broad definition includes glass ionomer cements, compomers and ormocers as-well as RBCs. Resin based materials were first introduced in dentistry during the late 1940s. The earliest materials that were used were simple and based on unfilled poly(methyl methacrylate). Several shortcomings existed, amongst which polymerisation shrinkage (20-25% by volume), poor colour stability, and poor chemical, physical and mechanical properties led to modifications in order to try to improve these properties. In 1951 Knock and Glenn attempted to solve the problem of polymerisation shrinkage by including inorganic filler particles in the resin (Knock and Glen, 1951) although poor mechanical properties and significant discolouration remained a problem due to the absence of coupling agents between the filler and the resin.

In the early 1960's, resin based composites were introduced into the dental scene (Bowen, 1962) by Dr Rafael Bowen and this class of material looked a promising alternative for amalgam. Bowen developed the monomer in 1956 after attaching methacrylate groups to an epoxy monomer to form bisphenol A-glycidyl methacrylate (Bis-GMA). The introduction of a high-molecular weight, di-functional monomer (Bis-GMA or Bowens Resin) greatly facilitated the commercial development of materials containing inorganic fillers. In 1962, Bowen patented the combination of Bis-GMA resin

and silane-treated quartz particles, which is the origin of most commercial composites available today (Bowen, 1962). The first use of composites in paste/liquid form was developed by Chang in 1969 and Henry Lee in 1970 (Glenn JF, 1982). These materials were conventionally polymerised with the redox methods of the cold-cure acrylic chemistry commonly used in restorative dentistry. In RBCs the two phases are the resin matrices (i.e. the monomers) and the fillers.

The fundamental work on the use of high molecular weight epoxy and methacrylate derivatives that incorporated inorganic filler loading has received wide spread attention and since then much work has been done to reach the current level of sophistication. However clinical longevity remains limited compared to amalgam and there exists a likelihood of root canal therapy if failure occurs (Drummond, 2008). However, since their inception, the properties of composites have been greatly improved resulting in higher mechanical properties (Palin et al., 2003), lower thermal coefficient of expansion (Stevens, 1999), lower dimensional change on setting (Watts and Cash, 1991), and higher resistance to abrasion (Turssi et al., 2005), which all contribute to the increased clinical performance.

1.1.1 Aesthetic quality

One of the major advantages of RBCs is the property of aesthetic quality. The aesthetic revolution began shortly after the introduction of RBCs in the 1970's. Although acrylic resin and silicate exhibited satisfactory initial aesthetics, the introduction of RBCs meant that for the first time, clinicians were able to restore teeth for aesthetics as well as form and function without rapid bulk discoloration occurring which was inherent with acrylic resins. The popularity of RBCs increased as the ability to mimic tooth structure through anatomical stratification improved and the incorporation of dyes and opacifiers became possible (Terry, 2003).

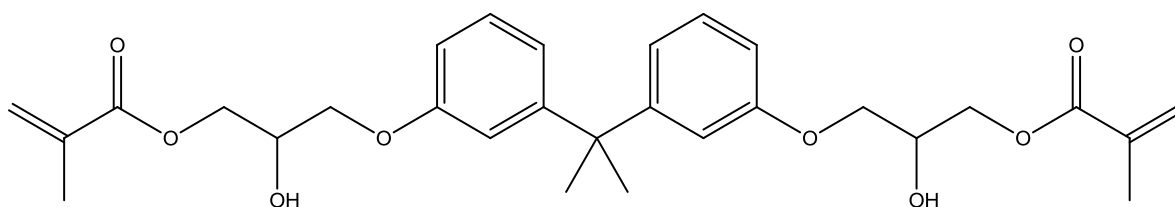
1.1.2 Composition

Composite materials generally consist of a resin-based matrix and an inorganic filler, the ratio of which affects the properties of the material. The incorporation of filler is the main strategy used to improve the poor mechanical and physical properties associated with unfilled resins. The filler gives the composite wear resistance, strength, reduced dimensional change aesthetics and opacity (Ferracane, 1995; Turssi et al., 2005). The composite also consists of a coupling agent such as silane, which enhances the bond between the filler and the resin matrix, and an initiator in order to initiate the process of polymerisation when external energy such as light or heat is applied. An inhibitor may also be added into the composite in order to prolong the shelf life of the monomer (Bowen and Marjenhoff, 1992) and improve ambient light stability. Furthermore, pigments and dyes (optical modifiers) may also be incorporated to improve aesthetics and aid colour matching with natural teeth.

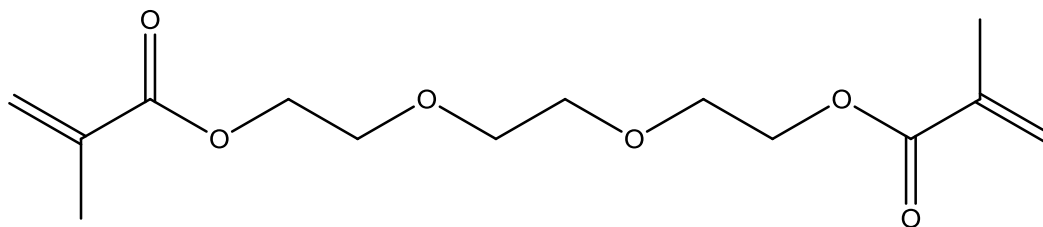
1.1.3 Resin matrix

Resin restorative composites are typically prepared from Bowen's resin, a compound of bisphenol A and two molecules of glycidyl methacrylate called 2,2-bis[4(2-hydroxy-3 methacryloyloxy-propyloxy)-phenyl] propane (Bis-GMA; Figure 1.1.1) (Bowen, 1962). Bis-GMA was the earliest base resin successfully incorporated into an RBC for direct restoration and has been the primary component of dental RBCs for more than half a century. Although it was fundamental in the introduction and development of RBCs, the relatively large methacrylate molecule has two aromatic rings and hydroxyl groups which add to its molecular weight and stiffness. Consequently, the material is very viscous (~1,000,000 mPa.s) and reactivity and degree of conversion remain low (Pfeifer et al., 2009). Furthermore, the material is too viscous to handle comfortably and the incorporation of reinforcing filler to an appropriate mass percentage for sufficient mechanical and physical properties is almost impossible. Di-functional monomers such as triethyleneglycol dimethacrylate (TEGDMA; Figure 1.1.1) and ethylene glycol dimethacrylate (EGDMA; Figure 1.1.1) which have low molecular weight are added to

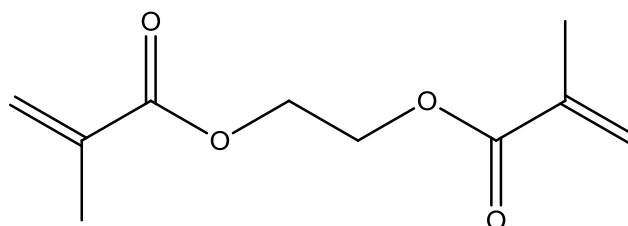
reduce the viscosity and act as diluents for the matrix system (Silikas and Watts, 1999) which facilitates the incorporation of greater amounts of inorganic filler and improves mechanical properties of the final RBC material. Additionally, these types of monomer contain reactive carbon double bonds at each end that can undergo addition polymerisation and therefore reactivity and degree of conversion are increased.



Bisphenol A glycidyl methacrylate
Bis-GMA



triethyleneglycol dimethacrylate
TEGDMA

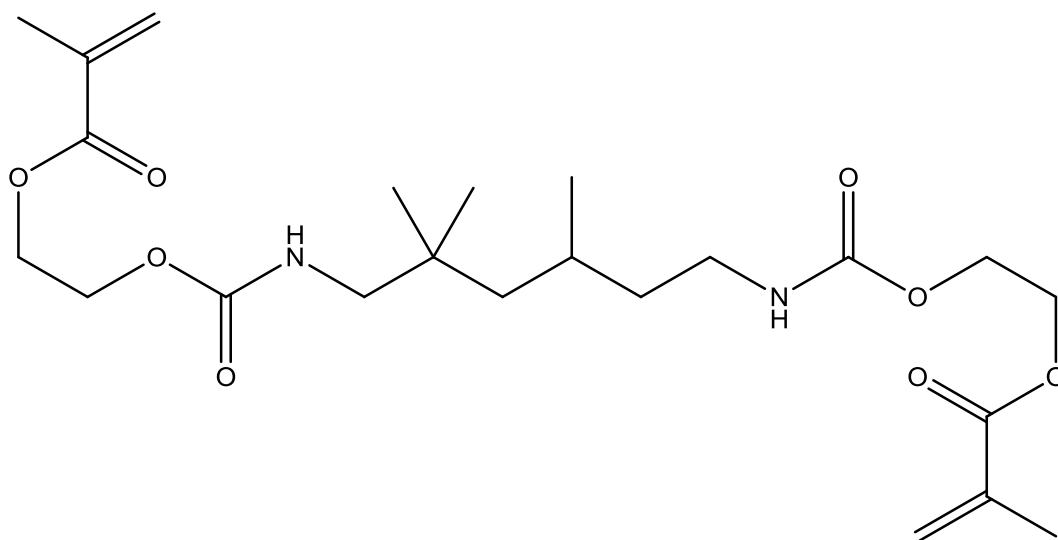


ethylene glycol dimethacrylate
EGDMA

Figure 1.1.1: The chemical structure of the base and diluent monomers used in RBC materials.

TEGDMA has a much lower molecular weight and thus the viscosity is significantly lower (~10 mPa.s). Although the reactivity increases due to the lower viscosity, the presence of ether groups (C-O-C) and the lack of aromatic rings along its structure, reduces its mechanical properties when compared to Bis-GMA. Furthermore, inferior physical properties are also reported, namely polymerisation shrinkage (TEGDMA = 12.5%, Bis-GMA = 5.2%) which is due to its lower molecular weight and increased concentration of carbon double bonds which inevitably improves conversion (Asmussen, 1982; Munksgaard et al., 1985; Braga et al, 2005). However, its low hydrophobicity may result in increased stain susceptibility and leaching of TEGDMA into the oral environment (Sideridou and Achilias, 2005).

Other methacrylate-based monomers such as urethane dimethacrylate (UDMA; Figure 1.1.2) may be used to replace Bis-GMA or maybe used synergistically with Bis-GMA to improve conversion and mechanical and physical properties (Palin et al., 2003). UDMA is derived from Bis-GMA but differs in the presence of urethane group, which offer a greater functionality to their counterparts. The functionality offered by these groups adds toughness and flexibility to the monomer backbone chain, providing the possibility for enhanced conversion and durability. Although it has similar molecular weight to Bis-GMA, the lack of aromatic groups produces a less viscous resin (~11,000 mPa.s), which significantly improves its handling properties. To date, no polymerisation system has been successful in the long term as the methacrylate based monomers in dentistry.



7,7,9-trimethyl-4,13-dioxo-3,14-dioxo-5,12-diaza-hexadecan-1,16-diol dimethacrylate
UDMA

Figure 1.1.2: The chemical structure of the lower molecular weight base monomer, UDMA compared to Bis-GMA, which contains urethane functional groups.

Other monomer systems are also available which have been used to address the problem of polymerization shrinkage. Ring opening chemistry has been available since the mid-1970s where homo-polymerisation of bicyclic ring-opening monomers were reported to result in expansion rather than shrinkage upon polymerisation (Bailey, 1975). The expansion was believed to be a result of the double spiro-cyclic ring opening of the spiro-orthocarbonate (SOC) molecule. However, it was not until the late-1970s that such chemistry was considered in dental composites by Thompson et al. (1979) who unsuccessfully attempted to combine the strength of conventional resins with the expansion properties of SOC. The incorporation of Bis-GMA with SOC was impossible and the resin mixture resulted in decreased monomer conversion, which affected mechanical and physical properties. Then in 1992, by variations in ring size and melting points, the incorporation of methacrylate monomers such as Bis-GMA and TEGDMA, became possible (Stansbury, 1992 and Stansbury 1992b). These were polymerised by a cationic and a free-radical mechanism and showed a significant improvement in polymerisation shrinkage than a methacrylate-based control group. Further developments occurred over the years to eliminate polymerisation shrinkage (Byerley et al., 1992; Moszner and Salz, 2001), however, problems of reactivity remained which resulted in

inadequate saturation of SOC rings, decreased cross-linking and therefore decreased mechanical properties.

In attempts to reduce polymerization shrinkage, thiol-ene resin chemistry has also been suggested for the use in dental restorative materials (Jacquelyn et al., 2005). The reaction of thiol-ene monomers creates reactive thiol or vinyl (ene)- functionalized oligomers. Thiol-ene resins polymerise by a step growth mechanism rather than a chain growth mechanism found in dimethacrylate based resins. As a result, the gel point is significantly delayed which enable better control of the polymerization process, i.e. molecular weight control and reduced shrinkage stress (Cramer and Bowman, 2001). Furthermore, the use of such monomers has been reported to improve shrinkage by up to 33% when compared to methacrylate systems (Patel et al., 1987; Jacquelyn et al., 2005). However, although thiol-ene resins performed well physically, i.e. reduced shrinkage and shrinkage stress, their mechanical properties remain inferior to Bis-GMA/TEGDMA resins (Jacquelyn et al., 2005). Recently composite methacrylate-thiol-ene formulations have been shown to exhibit improvements in methacrylate conversion, flexural strength, shrinkage stress, depth of cure, and water solubility (Boulden et al., 2010). In such systems, the thiol-ene acts as the reactive diluent in ternary formulations which then results in a unique polymerization kinetics and shrinkage dynamics which is a combination of both thiol-ene and methacrylate chemistry (Cramer et al., 2010). Additionally, such systems exhibit a pseudo-hybrid polymerization process whereby the first stage is dominated by methacrylate homo-polymerisation and chain transfer to thiol and the second stage is dominated by thiol-ene polymerization (Cramer et al., 2010). Consequently, the polymers produced are expected to show improved physical and mechanical properties. Research is ongoing on this topic.

The utilization of epoxy resin chemistry has also been realised for decreasing reduction in polymerisation shrinkage. 3M ESPE Dental Products (Seefeld, Germany) introduced 'Filtek™ Silorane' (c.2006), which is an epoxy based RBC developed from oxirane-based materials. Silorane (3,4-epoxycyclohexylcyclopolydimethylsioxane) is synthesised from oxirane and siloxane moieties resulting in a polymer with similar mechanical and physical properties compared to methacrylate based RBCs whilst showing increased hydrophobicity (Palin et al., 2005), decreased cytotoxicity (Schweikl

et al., 2004), increased ambient light stability (Weinmann et al., 2005) and most importantly decreased polymerisation shrinkage and stress (Ernst et al., 2004; Weinmann et al., 2005). Such composites utilise an innovative approach of ring-opening chemistry. The monomers connect by opening, flattening and extending towards each other, which results in a significantly lower polymerisation shrinkage in contrast to a methacrylate based composite in which the monomers are linear and connect by actually shifting closer together, which results in a greater loss of volume.

The initiating system in FiltekTM Silorane composes of three components, camphoroquinone, an iodonium salt and an electron donor. The electron donor decomposes the iodonium salt to form the cation which initiates the polymerisation (Weinmann et al., 2005). A significant improvement over traditional methacrylate based RBCs in polymerisation shrinkage (<1%) has been reported due to the ring opening cationic polymerisation (Weinmann et al 2005).

1.1.4 Fillers

The incorporation of fillers is a strategy whereby inorganic particles or pre-polymerised material is added in proportion to the composite, which then proportionally reduces the shrinkage (Bowen, 1962; Atai and Watts, 2006) by offsetting it with the relatively inert filler particles (Condon and Ferracane, 2000). However, even this method is reported to produce relatively high polymerisation shrinkage (Atai and Watts, 2006).

In the 1970s and 1980s, it was recognised that the size and amount of filler was critically important for the material and physical properties and thus RBCs were characterised according to the filler content. Modification in the filler component will change rheology, curing kinetics, degree of conversion and ultimately the mechanical and physical properties (Leprince et al., 2009; Beun et al., 2009).

Additionally, the type of filler is also important towards both optical and physical properties. The use of quartz provides excellent optical match to the polymer and is readily available. However, several drawbacks exist. Namely, it is not radioopaque and is very abrasive towards enamel. Furthermore, the relatively large particles reduces its polishability and therefore this was addressed with the development of amorphous silica.

However, the relatively small particles did not permit large volume fractions of filler to be incorporated. Consequently pre-polymerised fillers are added (which have larger particle size) to improve volume filler fraction. The use of such a technique improved the polishability of the material dramatically but the problem of radio-opacity remained. Most modern composites are filled with radiopaque silicate particles based on barium, strontium, zinc, aluminium, or zirconium (Ferracane, 1995).

Composites have been mainly classified into five groups, which may also have sub-group classifications:

Macrofilled: Traditional macrofilled composites were developed in the 1970's and are based on quartz, strontium or barium glass and had filler particle sizes ranging from $10\mu\text{m}$ – $40\mu\text{m}$. The quartz filler had good aesthetics and durability but suffered from absence of radiopacity and low wear resistance. Although barium and strontium glass particles are radio-opaque, they are less stable than quartz.

Microfilled: the poor polishability and relatively high wear resistance of macrofiller introduced much smaller particles in the 1970s which contained colloidal silica filler with a particle size of $0.01\mu\text{m}$ – $0.05\mu\text{m}$. The small size made it possible to polish the resin composite to a smooth surface finish but due to the small volume and high surface area of the filler, it was difficult to obtain high filler loads, which led to poor physical properties.

Hybrid: This type of filler was introduced to solve the problem of polymerisation shrinkage whilst maintaining good physical properties by combining both macrofilled and microfilled technologies together.

Modern Hybrid: This type of filler was designed to combine the advantages of macrofilled and microfilled composites but additionally contain reduced submicron fillers. However, they do not have the final surface finish potential or translucency of microfilled resin composites.

Nanocomposites: Although this type of filler existed as early as the 1970s as microfills, there was no concept of nano-technology and therefore this classification did not exist. However, in 2003, 3M ESPE marketed the first composites which they claimed to be nano but was actually manufactured with micron size agglomerated clusters of nano-particles. Since then, several manufacturers have used nano sized particles in conjunction with microfillers and have suggested improved mechanical properties as a consequence.

However, the classification of ‘nano’ creates some debate amongst researchers as to date no commercial composites truly contains only nano sized particles.

Composite Type	Filler Size (μm)	Filler Material
Macrofilled	10 - 40	Quartz or Glass
Microfilled	0.01 - 0.1	Colloidal Silica
Hybrid	15 - 20 and 0.01 - 0.05	Glass and colloidal silica
Modern Hybrid	0.5 - 1 and 0.01 - 0.05	Glass, Zirconia and colloidal silica
Nanofiller	< 0.01 (10nm)	Silica or Zirconia

Table 1.1.1: Summary of filler properties and classification of fillers in dental composite materials.

1.1.5 Inhibitors

Inhibitors are used commercially to prevent spontaneous polymer formation and extend the shelf life of RBCs. Manufacturers of dental resin composites typically add hydroquinone or butyl hydroxytoluene (BHT) to the composite in order to inhibit polymerisation (Figure 1.1.3).

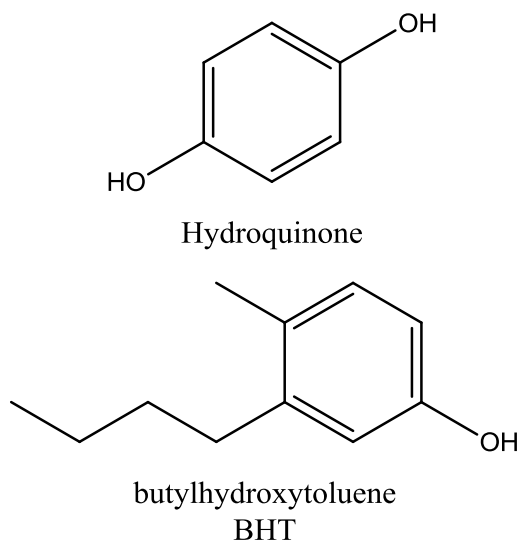


Figure 1.1.3: The chemical structures of the common inhibitors used in RBCs

The inhibitors react with the free radical thus decreasing the rate of initiation and increasing the rate of termination (Moad and Solomon, 1995). Decreasing the rate of initiation results in retardation of polymerization and an increase in the rate of termination. Although chemical inhibition may be used to extend the shelf life of composites, inhibition of polymerisation may also occur when there are large amounts of oxygen present and this could lead to partially cured restorations which may cause problems because the optimum properties are not obtained (Gauthier et al., 2005).

BHT is a commercially available primary antioxidant and is also used as an inhibitor in dental composites (Geurtsen, 1998). Reduction in reaction speed by chemical inhibition occurs as the free radicals are terminated by reacting with the phenolic hydrogen of the BHT molecule. BHT molecules are primary antioxidants which are effectively free radical scavengers, which combine with peroxy radicals and break the autocatalytic cycle. The phenoxy radicals may then inactivate another free radical by C-C or C-O coupling or loss of another hydrogen atom to form a quinone, which may react further. Therefore, each inhibitor molecule can terminate two or more polymer chains (Moad and Solomon, 1995). The conversion of monomer to polymer proceeds at a reduced rate until all the inhibitor is consumed. Inhibition may also extend the 'pre-gel phase', which will allow shrinkage forces to be dissipated before the cross-linking reaches a certain point at which molecular displacement becomes impossible (Braga and Ferracane, 2002). However, this idea has never been commercialized and BHT is only used to prolong the shelf life of a commercial composite by acting as a stabilizer to inhibit auto polymerization.

1.1.6 Pigments and UV absorbers

Pigments and dyes are used commercially to provide shading and opacity to aid colour matching to natural teeth. Typically metal oxides such as titanium and aluminium oxides are used. Titanium oxide provides opacity, whilst magnesium, copper and iron oxides provide a variety of colours and shades.

As dental composite fillings are prone to discolouration due to UV absorption (below 400nm), UV absorbers such as Benzophenone are used commercially, which absorb UV light to prevent such discolouration.

1.1.7 Photoinitiators

The first photo-polymerised composite system was developed in the 1970's under the trade name of Nuva; Dentsply/ Caulk (Rueggeberg, 2002) as an attempt to decrease the lengthy setting times. Photo-initiation polymerisation involves the use of a light to produce free radicals to start the polymerisation process and allows the ability to 'command-cure' restorations. This has become highly popular in dentistry as degree of conversion can range from 40-70% with appropriate conditions (Fonseca et al., 2004; Mendes et al., 2005; Ogunyinka et al., 2007). Systems used in current restorations and luting-agents frequently use both photo- and chemical-initiation (dual curing) as it is often difficult to expose the material to sufficient light to reach the maximum degree of conversion and thus maximum strength (Watts, 2005). The potential advantage of photo-activated composites over chemically activated composites is that they have greater handling time because the polymerisation process is initiated once the composite is irradiated with light, whereas for the chemically activated composite resins, working time was extremely limited and setting times range from 3 to 5 minutes once the material is mixed and can only be controlled by modifying the concentrations of initiator and accelerator (Fonseca et al., 2004).

1.1.7.1 Camphoroquinone

The photoinitiator system within a RBC generally consists of a two-component system, a photoinitiator and a co-initiator. Camphoroquinone is the most commonly used visible light activated free radical photo-initiator for dental resins. The absorption range corresponds to visible blue light between 400-500nm ($\lambda_{\text{max}} = 470 \text{ nm}$). The co-initiator or photo sensitizer is normally a tertiary aliphatic amine reducing agent, which reacts with camphoroquinone in its excited triplet state to generate free radicals. The absorption of

one photon of radiation promotes the non-bonding electron of the carbonyl group to an excited state based on the π^* anti bonding orbital. This excited state may return to the ground state by fluorescence or a radiation-less transition or it may decompose to another species. If the excited state is an excited singlet state, it may undergo intersystem crossing to the triplet state. The excited triplet then forms an exciplex with the amine reducing agent by charge transfer from the nitrogen lone pair to carbonyl, thus producing two radical ions. In the exciplex the amine transfers the hydrogen localized in the α carbon to diketone, resulting in a production of an amino radical and a cetyl radical (Figure 1.1.4). The amino radical is responsible for the initiation of the polymerisation reaction, while camphoroquinone cetyl radical is inactive (Watts, 2005; Alvim et al., 2007).

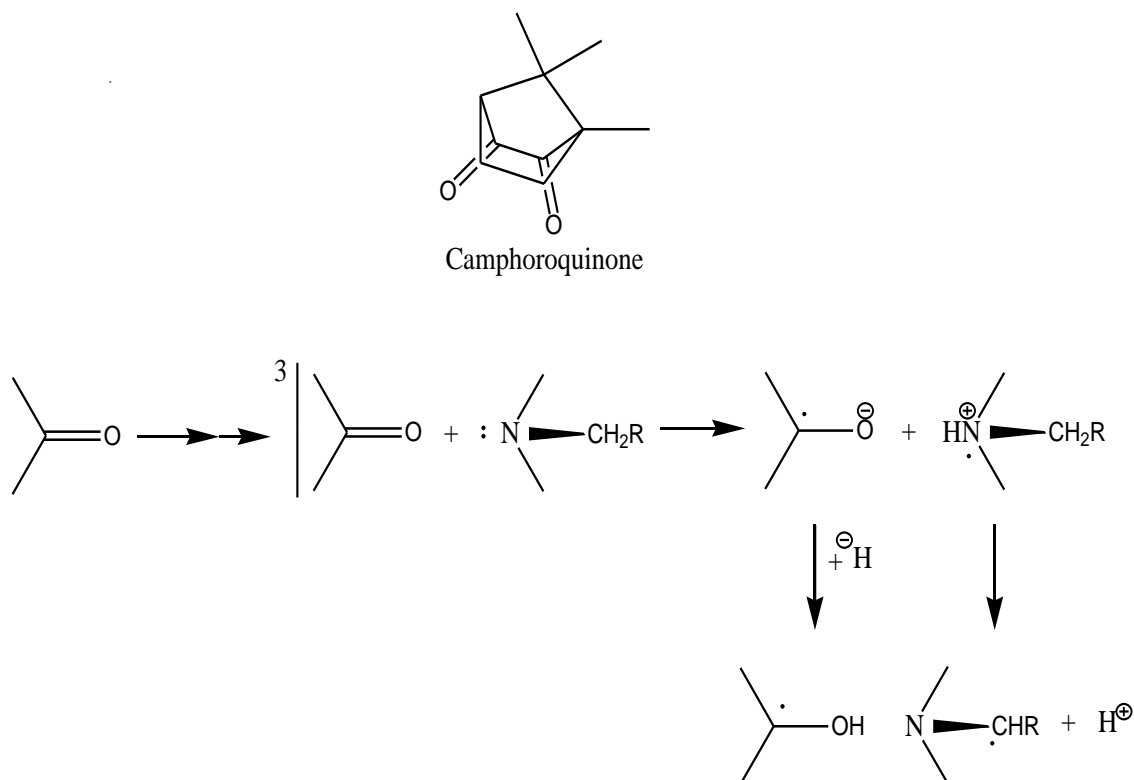
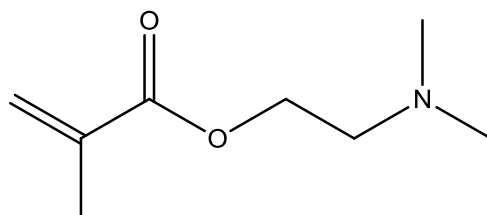


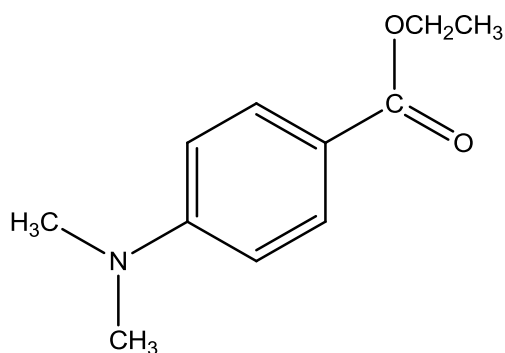
Figure 1.1.4: The production of free radicals during initiation in resins and RBCs containing camphoroquinone and an amine reducing agent.

Although several amine reducing agents (Figure 1.1.5) are capable of providing electrons for charge transfer during the initiation process, dimethylaminoethyl dimethacrylate (DMAEMA) is commonly used. The efficiency of this is dependent upon

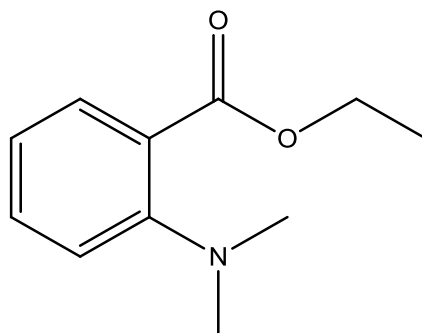
its chemical structure and functional groups and affects clinically important properties such as the rate of polymerisation, the depth of cure and the final monomer conversion (Rueggeberg et al., 1997; dos Santos et al., 2007). Aliphatic amines, such as ethyl dimethyaminobezoate and dimethylaminobenzoic acid ethyl ester may show a greater efficiency due to the aliphatic functional groups that have a greater electron density.



dimethylaminoethyl methacrylate
DMAEMA



Ethyl p-N-dimethylaminobenzoate
EDB



N,N-dimethyl-aminobenzoic acid ethyl ester
DABE

Figure 1.1.5: Chemical structures of co-initiators/ photo sensitisers used in photo-active RBCs.

Appropriate photoinitiator chemistry is essential for efficient polymerisation to achieve satisfactory mechanical and physical properties of the polymer (Ogunyinka et al., 2007). An optimum correlation between photoinitiator and co-initiator type and concentration will maximize photon absorbance efficiency which may in turn maximize depth of cure of a filled system (Chen et al., 2007; dos Santos et al., 2007). The photoinitiator concentration should be limited in order to obtain an optimum photo-curing reaction with high monomer conversion since excessive un-reacted photoinitiator, products of their photolysis, and any un-reacted monomer, may cause cytotoxicity (Pagoria et al., 2005). Furthermore, the concentration of camphoroquinone affects the aesthetics of a restoration as exceeding a critical concentration limit of camphoroquinone will lead to yellow discoloration and any un-reacted molecules may return back to the ground state (Ogunyinka et al., 2007) which will cause discolouration of the final polymer. Such discoloration may reduce aesthetic quality (Rueggeberg et al., 1997; Suh BI., 1997; Studer et al., 2001; Brackett et al., 2007) and affect light transport through material thickness. Furthermore, the discoloration may also create problems in color matching to natural teeth.

1.1.7.2 Alternative photoinitators

More recently alternative photoinitiators (Figure 1.1.6) have been used in RBCs such as phenyl proanediene (PPD), Benzil (BZ) and Norrish Type I photoinitiator systems such as mono (Lucirin TPO) and bi- (Irgacure 819) acylphosphine oxides (Neumann et al, 2005; Neumann et al, 2006; Ogunyinka et al., 2007). As some of these are non-pigmented, they have been utilized in bleach shade RBCs either synergistically with camphoroquinone or as a stand alone photoinitiator which may improve polymerisation kinetics, mechanical properties and aesthetic quality (Park et al, 1999; Weinmann et al, 2005; Neumann et al, 2006; Shin and Rawis, 2009). For such photoinitiators, it is interesting to note the absorption characteristics corresponding to the range of conventional light curing units (halogen~380-550nm, LED~440-500nm; Section 1.4), which will inevitably dictate their compatibility, efficiency and clinical acceptance (Table 1.1.2).

Photoinitiator	UV-Vis Absorption		
	Absorption Range (nm)	λ_{max} (nm)	Molar extinction coefficient at λ_{max} (L.mol ⁻¹ .cm ⁻¹)
Camphoroquinone	400-550	470	~35
Lucirin TPO	300-430	381	~550
Irgacure	300-440	370	~300
PPD	300-480	393	~150
Benzil	300-460	385	~50

Table 1.1.2: The absorption characteristics of photoinitators used in RBCs.

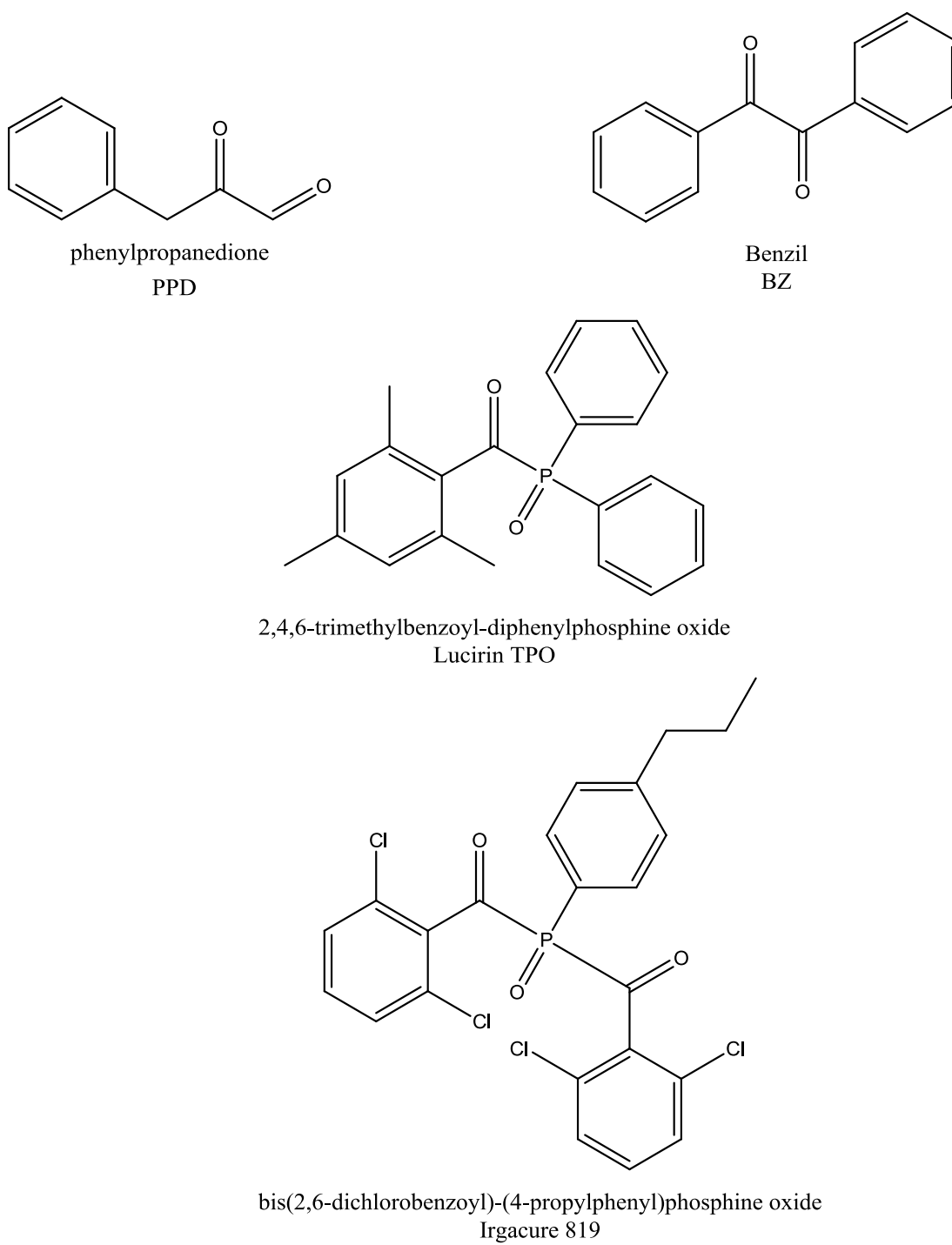


Figure 1.1.6: The chemical structures of alternative photoinitiators used in RBCs.

1.1.7.3 Lucirin TPO

Lucirin TPO seems to be an interesting molecule despite its shorter wavelength absorption as it may provide the potential to be a more aesthetic and a more efficient photoinitiator due to less absorption in the visible region and its much higher molar extinction coefficient, respectively (Table 1.1.2). The Norrish Type I photoinitiator produces free radicals by the photochemical cleavage of the carbon-phosphorus bond resulting in two initiating radicals without the need for a co-initiator or photosensitizer (Figure 1.1.7). The radicals are more reactive than the radicals created in the camphoroquinone-amine system. Ilie and Hickel (2008) demonstrated the potential for Lucirin TPO to replace camphoroquinone in adhesives with regards to degree of conversion and hardness. Other investigators have also reported the possibility of improved conversion with Lucirin TPO due to the much higher molar absorptivity (Allen, 1996; Neumann et al., 2005). However, the use of low wavelength light to initiate radicals maybe problematic to achieve sufficient depths of cure as the high molar absorptivity will reduce light transmission at depth (Neumann et al., 2005). Furthermore, interfacial scattering in RBCs is expected to be higher at lower wavelengths, which will further reduce light transmission (Shortall et al., 2008). Such problems can be overcome by adequate improvement in light curing technology to improve spectral emission/absorption mismatch, and adaptation of filler size to reduce low wavelength scattering. Additionally, using camphoroquinone or other photoinitiators synergistically, which have higher wavelength absorption, may aid curing at greater depths. Nevertheless, RBCs containing such initiators are limited to relatively thin specimens, which may not be a problem in flowable composites and adhesive layers of restorations.

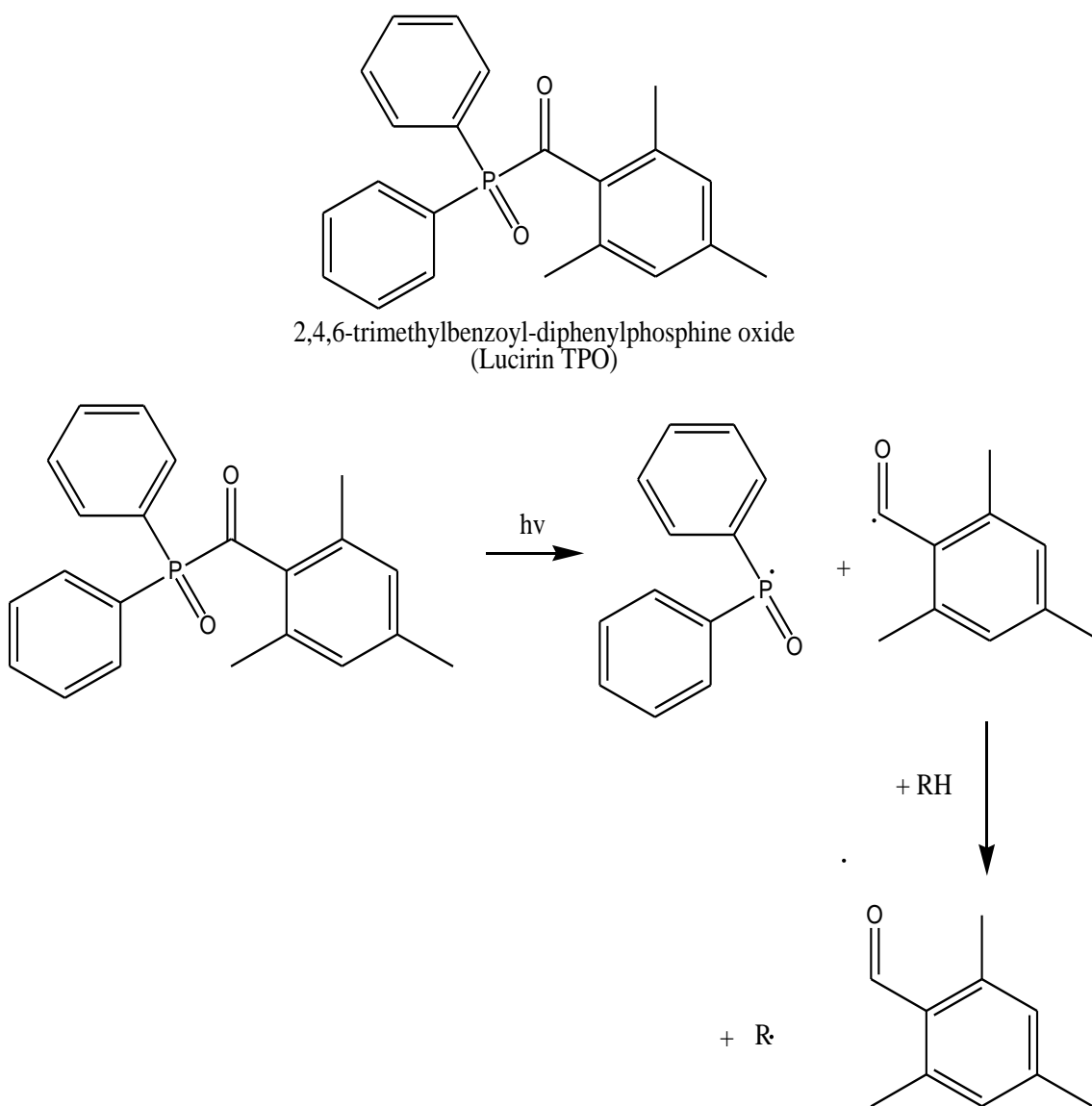


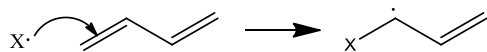
Figure 1.1.7: The production of free radicals during initiation in resins and RBCs containing Lucirin TPO.

1.2 Photo-induced Polymerisation

Resin polymerisation is characterised by three processes; initiation, propagation and termination. As described earlier, the photoinitiator is responsible for the initiation of the reaction by the production of free radicals ($\text{X}\cdot$), which can then open the double bonds at both ends of the monomer, which leads to cross-linking and propagation (Figure 1.2.1). Termination occurs as a result of either radical-radical interaction or radical entrapment in a highly cross-linked structure, which prevents further interaction with

vinyl double bonds. However, initiation, propagation and termination are not successive steps but co-exist during polymerisation, the kinetics of which are affected by several factors including monomer composition (Lovell et al., 1999), photoinitiator chemistry and filler percentage (Ogunyinka et al., 2007).

Initiation



Propagation

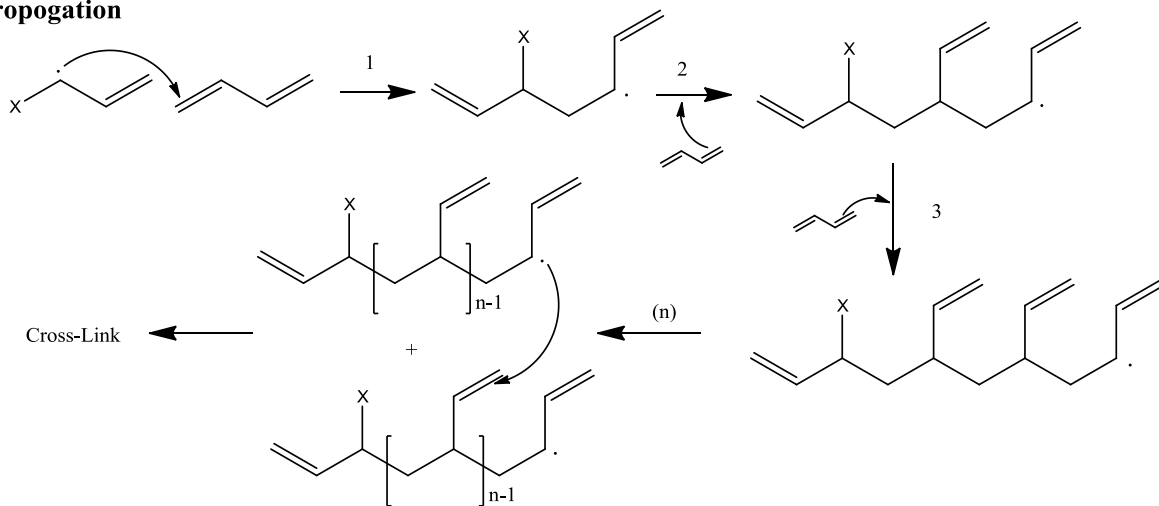


Figure 1.2.1: The schematic representation of the initiation and propagation process during photopolymerisation. (1), (2), (3) and (n) represent the theoretical steps of linear monomer addition.

1.2.1 The kinetics of polymerisation

As initiation, propagation and termination co-exist during polymerisation, reaction kinetics are complex and vary during the curing reaction, which are governed by the restriction imposed by system mobility. The system mobility is affected by material viscosity and has profound affects on cure kinetics. When system mobility is high, i.e. low viscosity, termination will mainly occur by radical-radical interaction (bimolecular). However, when system mobility is low, i.e high viscosity, then termination will mainly occur by a monomolecular pathway (radical-monomer).

$$R_p = k_p [R\bullet][M] \quad \text{Equation 1.2.1}$$

$$R_t = k_t [R\bullet]^2 \quad \text{Equation 1.2.2}$$

Equation 1.2.1 and 1.2.2 can be used to describe the polymerization process, where R_p and R_t are the rate of propagation and termination respectively, k_p and k_t are the respective rate constants and M and $R\bullet$ in square brackets are the monomer concentration and free radical concentration respectively. At the very beginning of polymerisation, conversion and therefore the extent of cross-linking is low, and it can be assumed that the concentration of free radicals remains constant (i.e. equilibrium between rate of initiation, R_i where $R_i = R_t$ and R_p) with the steps of polymerisation being chemically controlled. However, according to Anseth et al. (1996), for dimethacrylate monomers, this is only valid during the initial stage of polymerisation where conversion is less than 10%. As propagation proceeds, conversion increases, which subsequently restricts the mobility of the curing system due to increased viscosity. The cross-linked polymer network transforms from a liquid to a rubber and consists of two components; the sol (the monomer) and the gel component. The gelation point is characterised by the point at which the mobility of propagating species becomes so low that free-flow of macro-radicals becomes restricted and termination becomes diffusion limited. However, small monomer molecules can still diffuse through the cross linked networks and create new growth centres. As mobility restrictions mainly affect large growing polymer chains but not small monomer molecules k_t decreases significantly but k_p remains largely unaffected (i.e. $R_p \gg R_t$). Consequently, $[R\bullet]$ increases and results in a dramatic increase in the rate of polymerisation, which is known as “autoacceleration” and is more pronounced in the case of more viscous monomers (Goodner and Bowman, 1999; Feng and Suh, 2007). Once the mobility restrictions become too high, diffusion becomes impossible (even for small monomer molecules), k_p decreases and a dramatic decrease in the rate of polymerisation is observed (“autodeceleration”) and conversion starts to plateau as the system becomes vitrified. Vitrification occurs when the glass transition temperature (T_g) of the polymerising material reaches reaction temperature, i.e ambient

temperature plus ΔT of the reaction exotherm (n.b. above this temperature the same polymer would return from a glassy to a rubbery state if heated up). The glass transition temperature of a monomer based of 50% Bis-GMA and 50% TEGDMA is approximately 47.5°C (Dewaele et al., 2009). Although the majority of polymerization takes place during irradiation, conversion then evolves for up to 24 h post irradiation as a result of reduction in free volume bringing radicals closer together allowing further reaction (Truffier-Boutry et al., 2006).

1.2.2 The problem of polymerisation shrinkage and shrinkage stress

A major limitation of using RBCs as a dental restorative is shrinkage during polymerization, which may lead to a poor marginal seal, marginal staining and potentially an increased risk of recurrent caries. Polymerisation leads to closer packing of the molecules and consequently a loss in volume and bulk contraction occurs as the resin matrix changes from a pre-gel state to a solid. The polymerisation shrinkage increases proportionally with the amount of carbon double bonds present in the reaction environment. Consequently, increasing the proportion of low molecular weight monomers increases the concentration of double bonds and thus the degree of conversion increases with an increase in the volumetric shrinkage. Nonetheless, replacement of the bulky Bis-GMA molecule with lower molecular weight monomers such as UDMA and bisphenol-A ethoxylated dimethacrylate (Bis-EMA) have been shown to provide improved mechanical properties (Indrani et al., 1995; Asmussen and Peutzfeldt, 1998; Palin et al., 2003) although polymerisation shrinkage remains a problem.

Since the mid-twentieth century, the main strategy to reduce polymerisation shrinkage has been to increase filler load. Shrinkage is an intrinsic property of the resin matrix and it follows that a reduction in the monomer percentage would proportionally reduce the shrinkage, as well as significantly increasing stiffness, which may also act to compromise the tooth/RBC interface. Polymerisation shrinkage remains, and to date no methacrylate-based chemistry has been developed to solve this problem, which may limit the clinical longevity of a restoration. The polymerisation shrinkage of modern

methacrylate based RBCs remain in the region of 1-4% by volume (Lai and Johnson, 1993; Watts and Hindi, 1999; Schmidt et al., 2010).

As described in section 1.2, the gelation point of an RBC occurs when the viscous flow of the curing monomer stops and is unable to keep up with the curing contraction (Davidson and Feilzer, 1997). When the composite resin is in the pre-gel state, no stress is conducted to the surrounding tooth structure. As the composite cures, the material can flow from the unbound surface to accommodate for shrinkage. The flow stops once the composite becomes too rigid due to the increasing modulus, and stresses are transmitted to the surrounding cavity walls. The shrinkage strain induced during polymerisation is characterised by its magnitude, direction and strain dependence and becomes problematic when it becomes restricted by the adhesive bond at the tooth-restoration interface, which causes it to manifest into shrinkage stress. The resultant stress may compromise the synergism between the tooth and restoration interface (Davidson et al., 1984), increase the likelihood of mechanical failure (Sakaguchi et al., 1992), permit the ingress of bacteria, causing pulpal irritation (Lutz et al., 1991), result in cuspal deflection (Abbas et al., 2003), and in some cases result in micro, or even macro-cracking of the surrounding tooth (Hübsch et al., 1999; Palin et al., 2005).

The stress magnitude is not an inherent material property and is determined by characteristics of composites such as filler content and modulus (Condon and Ferracane, 2000), curing rate (Bouschilcher and Rueggeberg, 2000) and degree of conversion (Braga and Ferracane, 2002). Many authors have also reported that the final stress is governed by the configuration factor (C-factor) of the cavity; the ratio of bonded to un-bonded surfaces and the compliance of the surrounding tooth structure or test system (Feilzer et al., 1987; Loguercio et al., 2004; Bragga et al., 2006; Watts and Satterthwaite, 2008). Some authors suggest that C-factors less than 1 result in the low development of shrinkage stress and the composite remains bonded to the cavity walls (Feilzer et al., 1987). The C-factor itself is governed by the geometry of the cavity, i.e. the diameter and the height which both independently influence the magnitude of stress generated.

Over the decades many changes have occurred in order to address the problem of polymerization shrinkage and its associated stress. Major changes in light-curing units and curing modes (Section 1.4.6) have occurred in attempts to try and prolong the gel

point, which has been reported to alleviate developing stress in vitro (Sakaguchi and Berge, 1998; Hofmann et al., 2002).

Any clinical benefit of so-called, “soft-start” curing regimes (which employ lower irradiance at the outset of polymerisation) has yet to be realised, and there remains a risk of under-curing if adequate intensity is not provided to the restoration. Many factors associated with light transmission will also affect how much shrinkage occurs. Such factors include, light intensity, wavelength (Section 1.4.6) and the shade and opacity of the composite.

1.3 Applications of Resin Based Composites

The synthetic resins of RBCs have been used in dentistry for many years as restorative materials or adhesives. The fundamental characteristics that allow their successful use as restorative materials are aesthetic quality (Burke et al., 2003; Haj-Ali et al., 2005) and adequate mechanical physical properties (Peutzfeldt, 1997). Furthermore, they are relatively easy to handle which allows easy manipulation in the oral cavity. RBCs have many direct and indirect uses which include veneers, anterior restorations, posterior restorations, crowns, bridges, repair of fracture restorations, liners, pits and fissure sealants, bonding of orthodontic brackets and splinting (Ritter, 2005; Sadowsky, 2006).

‘Extension for prevention’, a principle introduced by Dr GV Black in 1895 which quickly became a foundational stone for the practice in the 20th century. The principle was based on what was known in the 19th century and the limitations of technology, and made perfect sense but was later shown in several studies that the best dentistry was a least invasive one, one with the ultimate goal of conserving tooth structure (Ericson et al., 2003; Anussavice, 2005). Nowadays, composites are placed with the principle of ‘minimally invasive dentistry’ which was indicated through the use of adhesive restorative materials and new preparation instruments of the market (Rossomando, 2007). It is often difficult to work with high and medium viscosity composite pastes as they are difficult to handle. Flowable composites are better adapted to a less invasive preparation of the cavity as the composite easily flows due to lower viscosity. This type of composite

was first introduced in the 1980's. With respect to the more viscous packable RBCs, flowable composites can be produced in two ways; either by increasing the amount of diluent monomer or by reducing the filler mass percentage. Additionally, the use of flowable materials with reduced filler content has been shown to reduce shrinkage stress due to its much lower elastic modulus (Bragga et al., 2003). Again, this has been reported in vitro and its effect on lowering stress will differ depending on the ratio of bonded to non-bonded surfaces (Feilzer et al., 1987; Loguercio et al., 2004; Braga et al., 2006; Watts and Satterthwaite, 2008), compliance (Feilzer et al., 1987; Watts and Marouf, 2000), material composition (Anseth et al., 1996) and irradiation protocol (Sakaguchi and Berge, 1998; Hofmann et al., 2002). Application of such composites include minimally invasive class III cavities, class I and II cavities without masticatory loads, pit and fissure sealing as well as extended pit and fissure sealing and class V cavities. Furthermore, they may also be used for cavity linings in class I and II cavities to alleviate stress (Braga et al., 2003).

Packable composites were designed and developed for use in the posterior region in the late 1990's. Packable composites generally contain a high proportion of fillers or specialised filler technology such as integrated matrix fillers (Leinfelder et al., 1999). Such materials facilitate contouring of the approximal contact area when using standard metal matrices. Furthermore, they do not stick to instrument and exhibit good positional stability, which inevitably aids contouring of the occlusal surface. However, a more invasive preparation is required as this type of composite is highly viscous.

1.4 Light-Curing Technology

Light curing technology has improved over the last several decades. By 2007, the use of light emitting diode (LED) light curing units became wide spread and is set to dominate the light curing industry (Krämer et al., 2008). However, to date, four main types of polymerization initiating technology are available, namely halogen bulbs, plasma arc lamps, argon ion lasers and LED.

1.4.1 The quartz-tungsten halogen bulb

Quartz tungsten halogen (QTH) lamps have been used in dentistry for several decades despite their remarkably low efficiency due to the relatively large heat generation (Strydom, 2005). Furthermore, as halogen bulbs operate with a hot filament, high operating temperatures due to long cure times limits the lifespan of the bulb to approximately 100 h with consecutive degradation of the bulb, reflector and band-pass filters which are used to limit the emission range to the absorption range of camphoroquinone (370-550 nm) (Neumann et al., 2005; Krämer et al., 2008). Recently, one high power QTH light has been developed which uses a 340 Watt bulb producing up to 3000 mW/cm² and has a water-cooling system incorporated within the lamp housing to prevent the unit overheating and to limit the degradation of the bulb. The latter problem reduces the curing efficiency of the unit over time (Uhl et al., 2006). Typically for such lights, 5% of the total energy is visible light, 12% is heat, and 80% is light emitted in the infra-red region (Mills et al., 2002; Burgess et al., 2002).

1.4.2 Plasma-arc lamps

The reduction of curing time, which may have financial implications for a dentist, motivated researchers to find an alternative. This led to the introduction of plasma arc lamps (PAC) which emit at higher intensities (Rueggeberg, 1999). The emission of light occurs through the glowing plasma, which is composed of gaseous mixtures of ionized molecules such as xenon molecules and electrons. In comparison to QTH lamps, the spectral emission range is comparable but have a higher intensity focused centrally around 470nm (the maximum absorption wavelength of camphoroquinone).

1.4.3 Argon ion-lasers

Argon-ion laser technology has also been used which emit blue-green light of activated argon ions in selected wavelengths (between 450 and 500 nm) (Meniga et al., 1997). When compared to QTH lights, argon lasers show improved conversion rates and polymerization depths (Fleming and Maillet, 1999). Although these showed improved physical properties compared to conventional QTH systems, heat generation during polymerization combined with considerably high initial shrinkage stresses have been reported to be problematic (Christensen et al., 1999; Strydom, 2005). Furthermore, when compared to QTH lights, argon lasers show improved conversion rates and polymerization depths (Fleming and Maillet, 1999). However, these lights remain expensive due to the narrowness of the beam and their only interesting feature over other lights is their very low light dispersion over a distance (Krämer et al., 2008).

1.4.4 Light emitting diodes

Solid state light emitting diodes (LED) were later introduced as an attempt to solve the problems associated with conventional halogen lights as describe earlier. The efficiency of these lights was greatly improved as the power consumption was significantly lower. Light is generated through junctions of doped semiconductors (p-n junctions). In gallium nitride LEDs under forward biased conditions, electrons and holes recombine at the LED's p-n junction leading to the generation of blue light which is collimated through a polymer lens. The light emitted covers the maximum absorption range of camphoroquinone despite its lower bandwidth (440-500 nm) so no filters are needed (Nakamura et al., 1994). Despite the lower bandwidth, higher intensity light is produced by such LEDs and consequently higher photo polymerization efficiency has been reported (Neumann et al., 2006). Additionally, the advantages include a greater lifespan as LED lamps with a lifetime of several thousands of hours without a significant loss in intensity. Fifteen percent of the total energy produced is light energy and the remaining eighty five percent is emitted as heat energy, although the heat produced is in the opposite direction to the curing tip (Mills et al., 2002).

Over the last decade, LED technology has changed significantly with the development of recent “generations” of high power LEDs comparable to advances in high tech computer technology. The first generation LED lights were unable to compete with conventional halogen technology despite their greater efficiency. This was due to the lower power density, which forced manufacturers to build complicated arrangements of ten to fifteen diodes into one lamp (Burgess et al., 2002; Hofmann et al., 2002). Despite the higher intensity, the expectations toward a possible reduction in curing times with first generation LEDs could not be confirmed (Besnault et al., 2003).

Second generation LEDs (or power LEDs) were later developed which were capable of delivering rather high outputs ($>1500 \text{ mW/cm}^2$) with one single diode. These were shown to produce material properties of polymers which were comparable to QTH lights (Soh et al., 2004; Shortall, 2005) and the potential for reduced curing time without loss of material properties was recognized (Yap and Soh, 2005; Schattenberg et al., 2008). However, maximum light intensities were limited to approximately 2000 mW/cm^2 as the heat generation created clinical concerns regarding gingival and pulpal tissue damage caused by the so called photodynamic effect and therefore the reappearance of cooling devices occurred (Krämer et al., 2008).

More recently, the introduction of new photoinitiators created compatibility issues with narrow band LED light, which brought about the development of third generation (or poly-wave) LED lights (Price and Felix, 2009). This generation of LED lights comprises of two peaks, one in the visible region and one covering the shorter wavelength UV-Vis region. The use of such LED lights allows effective overlap of the emission spectra of the lights with the absorption spectra of other photoinitiators such as Lucirin TPO, Iragcure and phenyl propanedione (PPD), which absorb at shorter wavelengths, as well as the most commonly used initiator, camphoroquinone (Neumann et al., 2005; Neumann et al., 2006).

1.4.5 The importance of spectral emission/absorption overlap

For light curing units to be clinically relevant, they must adequately excite photoinitiators (Section 1.1.7). The type of curing light is important as the spectral emission varies (Figure 1.4.1) and consequently the absorption-emission overlap is affected. Although most dental light curing units are designed to provide maximum efficiency with camphoroquinone, it becomes problematic when manufacturers of RBCs incorporate alternative photoinitiators into materials which may cause a spectrum mismatch when first and second generation LEDs or even some halogen and argon-laser lights are used (Price and Felix, 2009). To obtain maximum polymerisation efficiencies, an appropriate selection of light curing unit must be considered which will be compatible with the photoinitiator system (Neumann et al., 2006).

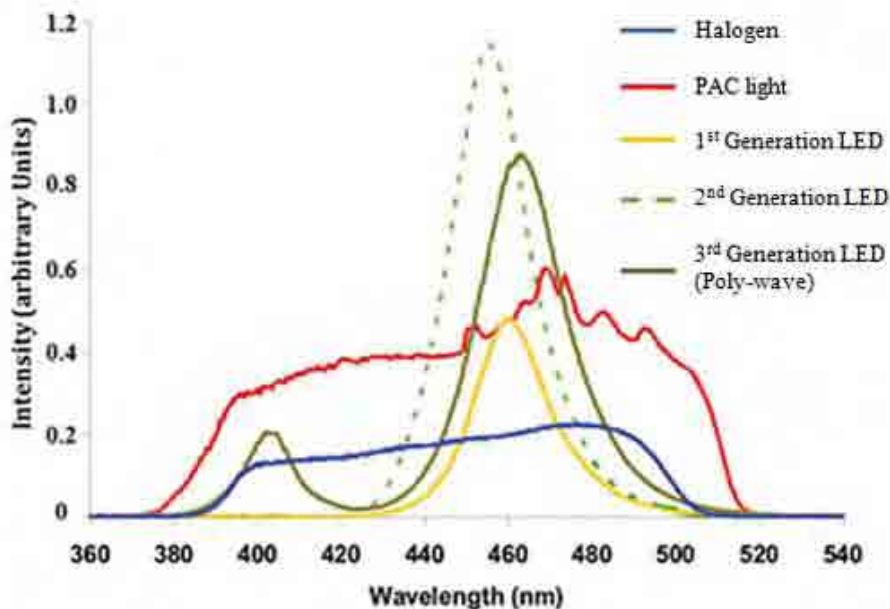


Figure 1.4.1: Schematic representation of the spectral irradiance of the different types of light curing units. (Modified from Price et al., 2010)

1.4.6 The use of alternative curing protocols to reduce shrinkage stress

Several curing protocols have been developed in order to address the problem of polymerisation shrinkage and its associated stress. Such techniques include, step curing, soft start and pulse-delay methods which have shown to improve conversion, material properties and marginal quality through the reduction in internal stresses (Sakaguchi and Berge, 1998; Hofmann et al., 2002).

The first attempt to reduce initial shrinkage stress was by the step curing method, which initiates polymerisation initially at 100 mW/cm^2 for 10s followed consecutively by a higher intensity dose (700 mW/cm^2). This technique may extend the pre-gel phase, whereby any developing stress can be alleviated through composites flow. After the gel-point, flow ceases and developing stresses cannot be compensated for. Although initial results indicated lower polymerisation stresses than conventional halogen lights, improved marginal adaption has not been reported (Muangmingsuk et al., 2003).

The soft start method is based on similar principles to the step curing method. However, the increase from low intensity (100 mW/cm^2) to higher intensity (800 mW/cm^2) occurs exponentially. However, the use of such a technique has not been unanimously clarified, although researchers have reported a reduction in stress, improved marginal quality of RBC restorations has not been proven in vivo (Krämer et al., 2008).

Pulse activated polymerisation methods use short impulses of high intensity irradiations. For example, ten pulses for at two seconds at 750 mW/cm^2 exposes the material to a total of twenty seconds of irradiation. The short pulses may extend the pre-gel phase and allow flow of the composite, which may alleviate stress. Pulse-delay methods have also been introduced, which initially irradiates the restoration with short pulses of light energy (pre-polymerisation at low intensity, e.g. 2 seconds or 20 seconds at 100 mW/cm^2). After a short waiting period of about three minutes, the final polymerisation is carried out for 30 seconds at high intensity. However this method remains controversial due to the low degree of conversion at the base of the cavity and furthermore no significant advantages have been reported (Yap et al., 2002).

1.5 The Exposure Reciprocity Law

The introduction of high powered curing lights allows the reduction in curing time with the advantage of minimised chair side procedure time, which has financial implications (Christensen, 2000). Furthermore, a reduced curing time seems attractive to a patient, who are understandably reluctant to keep their mouth open for a long time.

Curing time can be reduced by optimizing the spectral overlap between the light curing unit and the photoinitiator to increase the number of absorbed photons, however the general trend amongst manufacturers of light curing units has been to produce lights with increased irradiance with the assumption that curing time can be reduced proportionally (Burgess et al., 2002; Hofmann et al., 2002; Yap and Soh, 2005; Schattenberg et al., 2008). This is primarily based on the assumption that the degree of conversion and mechanical properties of photoactive materials are dependent upon the total energy of the irradiation, i.e. radiant exposure (J/cm^2), the product of irradiance (I , mW/cm^2) and irradiation time (t , s; Equation 1.5.1).

$$It = \text{constant}$$

Equation 1.5.1

The concept of total energy has been used extensively in other fields including photography and radiography and more recently has been applied to photo-active dental materials (Morgan, 1944; Halvorson et al., 2002; Peutzfeldt and Asmussen, 2005). The idea of total energy originated from the work of Robert Bunsen and Henry Roscoe, who in 1859 reported that all photochemical reactions are dependent upon the total absorbed energy and is independent of irradiance and time. However, failures were observed and a modification in the law was proposed by Scharzchild in 1900 (Equation 1.5.2):

$$I^p t = \text{constant}$$

Equation 1.5.2

where p is the Schwarzschild coefficient that he believed was a constant, which was latter shown to vary from material to material and in some cases with a change in irradiance. In 2003 it was concluded that the law proposed by Schwarzschild adequately modeled the

photo-response of a wide range of materials as a function of irradiance, as p lies between 0.5 and 1.0 for 97% of materials where the latter corresponds to complete conformity to exposure reciprocity (Martin et al., 2003).

In photo-active dental materials, the total energy principle promotes the common assumption that varying combination of curing irradiance and exposure time provides similar material properties for the same radiant exposure, irrespective of how radiant exposure is obtained (by different combinations of irradiance and time). This is the principle known as the “exposure reciprocity law”. This law has been tested with regards to several properties such as degree of conversion, flexural strength, flexural modulus, hardness and depth of cure (Feng and Suh, 2007) and in some cases the law holds quite well with regards to modulus (Sakaguchi and Ferracane, 2001), degree of conversion (Halvorson et al., 2002; Emami and Söderholm, 2003) and hardness (Price et al., 2004), whilst clearly in others it does not for similar outcomes (Peutzfeldt and Asmussen, 2005; Asmussen and Peutzfeldt, 2005; Feng and Suh, 2007). Although radiant exposure plays an important role, irradiance and time independently influences polymer chain length, extent of cross linking and mechanical properties and deviations from the law are reported. The investigation by Musanje and Darvell (2003) is of particular interest, as they tested highly filled commercial materials and their work highlighted the importance of particular composition and reactivity. Out of the four products tested at different combinations of irradiance and time, only one satisfied reciprocity with regards to flexural strength, modulus and total energy to failure. Additionally, a recent study by Feng and Suh (2009) also highlighted the inferior properties of some resin based composites and adhesives when cured for short exposure times at high irradiance and the same group earlier reported the importance of resin viscosity on degree of conversion and the validity of exposure reciprocity law (Feng and Suh, 2007). Other studies have also highlighted the inferior properties of lower viscosity systems, i.e. dentin bonding agents, compared with higher viscosity resin composites when using high power light curing units when compared to other lights (D’alphino et al., 2007).

The efficiency of the photoinitiators are dependent upon their absorption characteristics (Neumann et al., 2005) and the compatibility with the light curing unit (Section 1.4.5), which will determine the amount of free radicals created (Neumann et al.,

2006). As degree of conversion is directly related to the efficiency of the photoinitiator system (Emami and Soderholm, 2003; Ogunyinka et al., 2007;) and the amount of light absorbed by the photoinitiator, it should be expected that reciprocity should hold at a similar radiant exposure as the same total amount of efficient radicals are created by different irradiation protocols (Chen et al., 2007). However, at very high irradiance, photoinitiators with very low extinction coefficients i.e. camphoroquinone ($\sim 35 \text{ L.mol}^{-1}.\text{cm}^{-1}$), the photoinitiator system may become saturated resulting in no improvement in the initiation rate (Musanje and Darvell, 2003) and thus failures in the law may occur.

The introduction of high powered lights and other photoinitiators may allow the possibility of reduced curing time with increased irradiance although to date there is no real consensus on an adequate radiant exposure to produce sufficient polymerization, with radiant exposures ranging from 16-24 J/cm² (Rueggeberg et al., 1993; Koran and Kurschner 1998; Vandewalle et al., 2004) where the latter is usually considered a standard for less than 2 mm increments cured with a halogen lamp. As mentioned earlier, the combination of irradiance and time at the same radiant exposure will independently influence polymer chain length, extent of cross linking and mechanical properties. Peutzfeldt and Asmussen (2005) concluded that increasing radiant exposure increases the degree of conversion, flexural strength and flexural modulus. Furthermore, the same study also showed the trend of decreasing degree of conversion with high irradiance and short time irradiation protocols for each radiant exposure whilst maximum flexural strength and modulus occurred at an intermediate irradiance (Peutzfeldt and Asmussen, 2005). With high irradiance and short time irradiation protocols, degree of conversion and the kinetic chain length may decrease and the frequency of cross-linking increase, which may explain the parabolic relationship between irradiance, flexural strength and flexural modulus. Dewaele et al. (2009) suggested that degree of conversion, elastic modulus and T_g of a light curing polymer-based material are mainly and positively influenced by radiant exposure, whilst a negative correlation exists with respect to irradiance. Leprince et al. (2010) reported the impact of irradiation mode on the radical entrapment in photoactive resins. When similar radiant exposure were compared, higher concentrations of trapped free radicals were measured for longer times at lower irradiance. The authors suggested that since the premature termination of free radicals is

proportional to their squared concentration, which is higher for high irradiance protocols, early termination of free radicals would result in a lower degree of conversion. Furthermore, the concentration of trapped free radicals would also be influenced by the composition of the reaction medium as polymerisation leads to highly cross-linked networks where mobility restrictions would favour monomolecular termination (Section 1.2.1; Andrzejewska, 2001). Thus it is clear that the validity of exposure reciprocity law is complex and several factors will influence whether the law is upheld.

1.6 Optical and Physical Properties of Resin Based Composites

The electromagnetic spectrum consists of photons of varying wavelengths and energy (Figure 1.6.1). The interaction of matter with photons in the ultraviolet and visible regions of the electromagnetic spectrum is known as UV-Vis spectroscopy, whereas the spectroscopy of photons in the infra-red region is known as infra-red spectroscopy. Spectroscopy is the study of the interaction between radiation or light with matter as a function of wavelength, which can be used to assess the amount of a given substance or for the identification of a substance. Ultraviolet-visible (UV-Vis) spectroscopy has been used to calculate the quantum yield of conversion of photoinitiators in RBCs (Chen et al., 2007). Degree of conversion of the monomers used in RBCs has been analysed using mid- and near infra-red (MIR, NIR) spectroscopy (Antonucci and Toth, 1983; Ferracane and Greener, 1984; Urabe et al., 1991; Stansbury and Dickens, 2001).

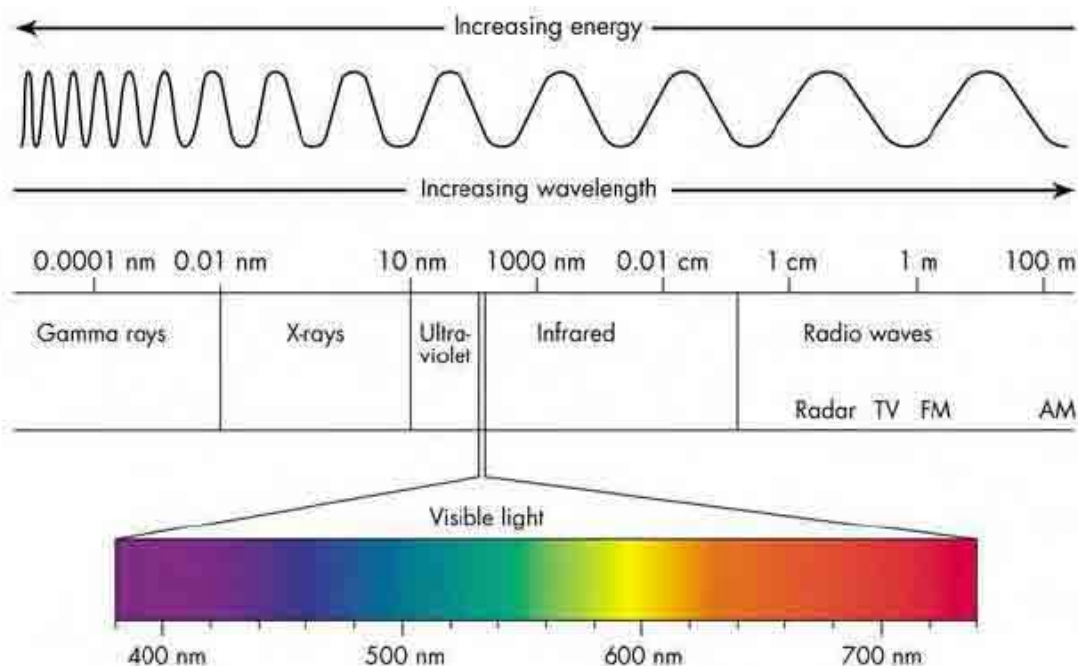


Figure 1.6.1: The electromagnetic spectrum and the associated wavelength and energies. (Adopted from www.antonine-education.co.uk, 10/12/2010)

1.6.1 Degree of conversion

The degree of conversion is an important property to characterise as it ultimately relates to physical and mechanical properties of RBCs. Conversion occurs as carbon double bonds of monomers are converted to extended networks of carbon single bonds (Section 1.2). Although it is expected that dramatic changes in mechanical and physical properties will occur when an RBC converts from a liquid to a solid, at higher conversion however, relatively small changes in conversion will also dramatically improve some polymer properties as they are primarily affected by the cross linking density of the network. Vitrification of the polymer network consequently limits the degree of conversion of RBC in the range of 55-75% (Ruyter and Øysæd, 1987; Ferracane, 1995) and the extent of polymerisation is also affected by other factors such as, composition, specimen geometry, photoinitiator chemistry, irradiation protocol, temperature and the presence of oxygen. Although the majority of the conversion takes places during irradiation, conversion of RBCs typically evolves for 24h post irradiation, with a

logarithmically related decline in ‘post cure’ (Vankerckhoven et al., 1982; Kildal and Ruyter, 1994; Truffier-Boutry et al., 2006).

Several analytical techniques have been used to assess the degree of conversion in RBC materials, amongst which a temperature based differential scanning calorimetry (DSC) and infra-red spectroscopy are the most common (Antonucci and Toth, 1983; Ferracane and Greener, 1984; Urabe et al., 1991; Stansbury and Dickens, 2001).

1.6.1.1 Fourier transform infra-red spectroscopy

Infra-red (IR) spectroscopy techniques are mainly based on absorption spectroscopy in the IR region of the electromagnetic spectrum (Figure 1.6.1) and is divided into three main regions; the near-, mid- and far- IR according to their relation to the visible spectrum. The absorption at specific frequencies is characteristic of chemical structure and can be used to identify and quantify specific functional groups in a molecule.

Fourier transform infra-red (FT-IR) spectroscopy has evolved as a technique which is used to obtain the IR spectrum of absorption, emission, photoconductivity or Raman scattering of a solid, liquid or gas, and is capable of simultaneously acquiring spectral data in a wide spectral range. The name ‘Fourier transform’ simply means a mathematical algorithm, which is used to convert the raw data into the actual spectrum. The spectrometer is based on a Michelson interferometer (Section 1.6.3.1), where light from a polychromatic IR source is collimated and directed to a beam splitter. The light is then split equally onto a fixed mirror and a moving mirror, which is then reflected back to the beam splitter. Ideally, fifty percent of this light is passed into the sample compartment, which is then focused through the sample onto a detector. An interferogram is produced by creating a ‘time-delay’ with the moving mirror and recording the signal at various points. This is then computed and the Fourier transform is applied by a computer programme to obtain the IR spectrum.

In RBCs, IR spectroscopy directly measures the conversion as a function of the un-reacted methacrylate groups. Several methods within this analytical technique exist, which utilise a transmission method (Ruyter and Györösi, 1976), a multiple internal

reflection method (Ruyter and Svendsen, 1978), and an absorption method (Stansbury and Dickens, 2001). The transmission method and the multiple internal reflection methods measure the decrease in the intensity of the carbon double bond stretching mode at 1637cm^{-1} as the methacrylate monomer is converted into polymer. However, these methods require reference bands that are stable during polymerisation which are typically based on aromatic groups of Bis-GMA (1608cm^{-1} or 1583cm^{-1}) or the N-H or C=O groups of urethane monomers. Conversion for a Bis-GMA/TEGDMA system is then calculated from the intensity at each band respectively (Equation 1.6.1):

$$\text{Conversion \%} = \frac{1638\text{cm}^{-1} / 1608\text{cm}^{-1} \text{ cured}}{1638\text{cm}^{-1} / 1608\text{cm}^{-1} \text{ uncured}} \times 100$$

Equation 1.6.1

Conversion has been characterised by FT-IR using other methods, namely potassium bromide (KBr) pellets for bulk conversion and Raman spectroscopy, which involves scattering rather than absorption (Stansbury and Dickens, 2001). However, most of these techniques are limited by the fact that only the conversion in thin specimens was possible or the sample had to be ground and destroyed to measure bulk conversion.

Near infra-red (NIR) spectroscopy has been used in medical and pharmaceutical industries for the non-destructive analysis of products and over the last decade has evolved as a technique to measure bulk conversion in RBCs having geometries of clinical relevance (Burns and Ciurezak, 1992; Murray and Cowe, 1992; Stansbury and Dickens, 2001). In RBCs, NIR monitors the absorption of the =C-H group of monomers at 6165cm^{-1} and 4743cm^{-1} (Bis-GMA/TEGDMA) which corresponds to conversion of the methacrylate groups. The band at 4743cm^{-1} is not typically used to characterize methacrylate conversion because the sharp baseline drop near its =C-H peak makes reliable peak area measurement difficult. Therefore the degree of conversion at 6165cm^{-1} is the calculated from equation 1.6.2:

$$\text{Conversion \%} = 1 - \left(\frac{\text{Final absorbance}}{\text{Initial absorbance}} \right) \times 100$$

Equation 1.6.2

It is also possible to measure the conversion of RBCs based on other monomers as the majority of the peaks lie between 6163 and 6167 cm^{-1} (Stansbury and Dickens, 2001). Additionally, such measurements do not require the presence of an internal reference peak and the non-destructive nature allows real-time measurement of conversion by the continuous measurement of the change in absorbance of the =C-H group. It is also then possible to interrogate curing kinetics, such as autoacceleration, rate maximum and time to reach maximum rates of polymerisation

1.6.1.2 An alternative to infra-red spectroscopy

Differential scanning calorimetry (DSC) provides a measure of the methacrylate conversion based on the enthalpy of the exothermic polymerisation reaction. This thermo-analytical technique uses the amount of heat required to increase the temperature of a sample and a reference as a function of temperature. The technique was originally developed by Watson and O'Neil (1960) and then later became commercially available. It is now used in a variety of applications including the measurement of polymerisation in dental polymers (Lee et al., 2004; Tanimoto et al., 2005).

The underlying principle of this technique is that when a sample undergoes a physical transformation from one phase to another, more or less heat will need to flow to it than the reference to maintain both of them at the same temperature. For example, when ice melts, it will require more heat flowing to the sample to increase the temperature at the same rate as the reference. This is an endothermic process in which heat is absorbed by the sample to undergo its phase transition. Conversely with RBCs, an exothermic polymerisation process (due to the conversion of the carbon double bonds to single bonds and photoinitiator decomposition) will require less heat to raise the sample temperature. Thus, by observing the heat flow between the sample and reference, it is possible to measure the amount of heat absorbed and released during phase transitions. By assuming the heat produced during the polymerisation is proportional to the number of monomer units converted under isothermal conditions, it is possible to determine the degree of conversion and kinetics of polymerisation (Tanimoto et al., 2005).

Differential Thermal Analysis (DTA) has been shown to be a reliable method for measuring the conversion of dental composites (McCabe 1985; Imazato et al. 2001). Other methods to measure temperature change during polymerisation have also been employed. Some methods are based on thermocouples and others are based on thermistors (Shortall and Harrington, 1998). Although polymerisation temperatures can be measured, several limitations exist. Firstly, the non-isothermal conditions prevent reaction kinetics from being determined and secondly the positioning of the temperature sensors (on the material surface or embedded within its bulk) and the volume of the test sample will dictate how much heat is registered. For this reason it seems DSC is more reliable, but thermocouple and thermistor based methods remain adequate for simple temperature measurements during polymerisation of RBCs.

1.6.2 Light transport through photocurable resins

The transmission of light and number of photons absorbed by the photoinitiator is critically important to the extent of polymerisation and is a desired parameter to fully characterise since it ultimately relates to the final mechanical and physical properties of the polymer. Inefficient light transmission and limited cure depths are associated with surface reflection (Harrington and Wilson, 1995; Watts and Cash, 1994), absorption (Chen et al., 2007; Ogunyinka et al., 2007), scattering by filler particles (Enami et al., 2005; dos Santos et al., 2008) and interfacial filler/resin refraction (Shortall et al., 2008). Upon irradiation, the photoinitiator will decompose (Section 1.1.7) and reduce its concentration. Polymerisation shrinkage of the curing resin (Section 1.2.2) will reduce the optical path length with the outcome of increased light transmission through the sample according to Beer-Lamberts law (Equation 1.6.3).

$$A = \epsilon Cl = -\log\left(\frac{I}{I_0}\right)$$

Equation 1.6.3

where, A is the absorbance, l is the optical path length, C and ϵ are the concentration and molar extinction coefficient of the absorbing species (i.e. the photoinitiator) and I and

I_0 are the light transmission and initial light transmission, respectively. Interestingly, varying patterns of light transmission in RBCs have been reported even using the same method. Harrington et al. (1996) developed a method based on a photo-diode which allowed the actual consumption of light energy by the activator to be demonstrated and reported an increasing light intensity with radiation time for most commercial composites. Using the same technique, Ogunyinka et al. (2007) reported decreasing transmission during polymerisation of experimental composite formulations. Then, Shortall et al. (2008) described the effect of monomer and filler composition on light transmission and reported that the refractive index mis-match between the filler and resin will affect its light transmission profile where maximum transmission was observed when the refractive index of the resin and filler are equal. This underlines the complexity of the optical properties of RBC materials when cured at depths as problems associated with attenuation and inherently non-uniform curing rates profiles exist which will inevitably affect the extent of monomer conversion.

1.6.3 Refractive index

As the resin matrix (typically, co-monomer mixtures of Bis-GMA and TEGDMA; Section 1.1.3) cures, its optical properties change and refractive index rises due to a rapid increase in cross-link density and viscosity. As the refractive index of the resin approaches that of the filler, interfacial filler/resin scattering is reduced and light transmission is increased (Shortall et al., 2008). The refractive index (n) of a medium at optical wavelengths is defined as the ratio of phase velocity of light in a vacuum (c) to the phase velocity (v_p), in a given medium (Shoemaker et al., 1989) and the velocity at which it propagates is wavelength dependent (Equation 1.6.4).

$$n = \frac{c}{v_p} \quad \text{Equation 1.6.4}$$

When light passes through the interfaces of two media, it is accompanied a change in velocity and a directional change. Hence the refractive index can also be

calculated by calculating the ratio of the sine of the angle of incidence to the sine of the angle of refraction (Figure 1.6.2).

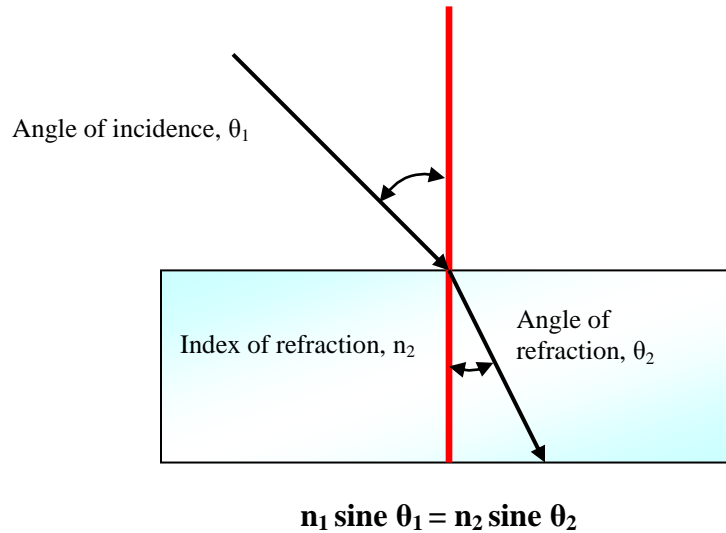


Figure 1.6.2: Schematic representation of the refraction of radiation through two media and the refractive index.

The relationship between the phase velocity in two media (v_1 and v_2) the angle of incidence (θ_1) and refraction (θ_2) and the refractive index of the two media (n_1 and n_2) can be predicted by Snell's law (Equation 1.6.3):

$$\frac{v_1}{v_2} = \frac{\sin \theta_1}{\sin \theta_2} = \frac{n_2}{n_1} \quad \text{Equation 1.6.3}$$

The most commonly employed technique to measure refractive index uses an Abbé refractometer. The refractive index is measured at the standard yellow Fraunhofer line of sodium (589nm) at 20 °C (n_{D20}) and is wavelength and temperature dependent. The measurement is based on observing the boundary line of the total reflection and measuring the angle of refraction. It is then possible to determine the refractive index of an unknown sample according to Snell's law if the index of refraction of the layer that is in contact with the sample is known.

1.6.3.1 Contribution of low coherence interferometry

Characterising the optical and physical properties of RBCs remains an important aspect in the development of restorative materials. Static and theoretical measurements of refractive index are well established for dental polymers (Shortall et al., 2008; Lehtinen et al., 2008). Additionally, polymerisation shrinkage of dental composites has been assessed by several methods amongst which dilatometry (Archimedes principle) (de Gee and Davidson, 1981) and the bonded disc method (Watts and Cash, 1991) are the most commonly employed although devices based on laser interferometry and optical magnifiers are available (Fogleman et al., 2002; Demoli et al., 2004; Arenas et al., 2007).

Low coherence interferometry (LCI) is an optical technique which has been applied to several fields related to industrial surface metrology, biological analysis, medical systems and more recently has been applied in the physical and optical characterisation of dental polymers to obtain thickness and topography (Tomlins et al., 2007; Morel and Torga, 2009). It is well suited to such applications as there is no direct contact between the sample and the fiber output and the possibility of measuring samples at high temperatures exists. In RBC materials, the ability to measure refractive index and thickness change simultaneously during curing photoactive resins may allow for the design of composite formulations which exhibit optimum evolution of translucency (as the refractive index of the curing resin approaches that of the filler) and enable precise control of aesthetic quality and light transmission through material bulk to cure at depths.

A ‘coherence grating’ provides precise axial positioning of an object in the direction of light propagation and therefore it is possible to obtain micron range accuracies. Information on finite volume and structural information as a function of depth may be obtained by focusing the light to obtain good transverse resolution (perpendicular to the optical beam) so that backscattered light is coupled back into the interferometer. This then mixes with a reference beam to form interference fringes, which is then detected by a photo-detector.

The most commonly used interferometer is the Michelson interferometer and was invented by Albert Abraham Michelson. The interferometer consists of two highly polished mirrors, a half silvered mirror and a monochromatic light source. The half

silvered mirror is partially reflective and therefore the beam can be transmitted in one direction and reflected in another. The principal is that when a beam of light coming from a monochromatic light source is incident on a half silvered glass plate, it is divided into equal intensities by partial reflection and transmission, i.e. both beams are coherent. Then if the two paths differ by a whole number, i.e by a moving mirror, there is constructive interference. Subsequently, when both beams recombine at the half silvered mirror, it produces the interference fringes that is detected at the detector (Figure 1.6.3)

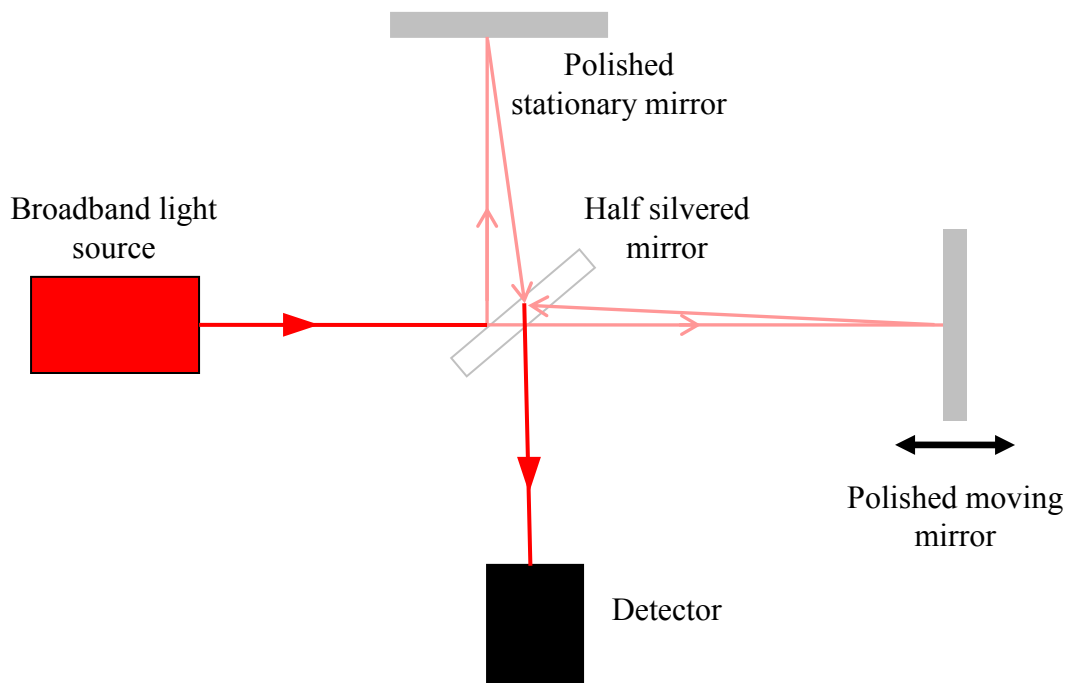


Figure 1.6.3: Schematic representation of a Michelson Interferometer

1.6.4 Colour

The matching of restorative materials to natural teeth is complicated as the human perception of colour varies according to differences in illuminant, background colour, positioning, translucency and human eye fatigue (van der Burgt et al., 1990; Dagg et al., 2004; Joiner, 2004). Additionally, moisture, skin tone and the surface characteristics of the restoration, i.e. gloss matt and roughness, further complicates matters. As described earlier (Section 1.1.1) the ability of RBCs to mimic tooth structure is one of its greatest assets and therefore a standardized method of quantifying any colour change is essential for a restoration to be aesthetically acceptable.

The composition within the RBC determines its ability to colour match to natural teeth, i.e the photoinitiator and pigments and dyes. One of the greatest problems associated with the use of pigmented photoinitators such as camphoroquinone is the presence of a partly unbleachable chromophore group (Jakubiak et al., 2003), which reduces its compatibility in bleach shade composites. Additionally, such photoinitators are typically used with amine photo-sensitizers which are also known to produce coloured reaction by products (Asmussen, 1985). This ultimately limits the concentration of camphoroquinone/amine that can be incorporated in to RBCs with concentrations ranging from 0.17 mass percentage to 1.03 mass percentage (Taira et al., 1988). Moreover, colour matching is further complicated by the fact that photoinitators decompose during irradiation with a subsequent loss of colour. However, the use of alternative photoinitators (Section 1.1.7.2) which absorb short wavelength visible light may show improved aesthetics and better colour matching properties particularly with bleach shade composites.

Clinically, during colour matching, RBC colour is quantified by the use of shade guides (van der Burgt et al., 1990) but remains subjective and operator sensitive. The same operator may even visualize the same colour differently at different times (Okubo et al., 1998). Despite many inherent disadvantages associated with visual matching, it remains the most common method for determining tooth colour and shade matching due to its cost effectiveness over instrumental colour measurement devices (Brewer et al., 2004).

1.6.4.1 Colorimetric spectrometry

The colour of RBCs relates to optical properties of opacity and translucency and may be measured using a colorimeter or spectrophotometer. The sample under test is irradiated with light, which is transmitted, reflected, refracted and diffused through the sample (van der Burgt et al., 1990). The reflectance pattern is then measured which is unique (fingerprint) for each colour and therefore can be numerically quantified. Although such measurements are more reliable than visual assessment, its accuracy is limited due to the translucent nature of tooth structure and the occurrence of edge loss, in which tooth structure may be affected by background colour (Seghi et al., 1990; Bolt et al., 1994; Brewer et al., 2004). Spectrophotometers are typically considered to be more accurate than colorimeters, however recent advances in modern colorimeters have greatly improved their reliability (Seghi, 1990).

For accuracy, such measurements are usually made with standard illuminants such as D65, which corresponds to average daylight conditions and illuminant A which represents incandescent lighting based on typical domestic tungsten filaments. The illuminant is irradiated onto the test specimen at a specific orientation (45° for dental materials) and the reflected light is then captured and analysed, the spectrum of which represents the materials colour. However, as colour is not static and differs with illuminant and viewing orientations, the Commission Internationale de l'Eclairage (CIE) have recommended four illuminating/viewing configurations to be used for materials with diffuse reflecting surfaces such as RBCs (Seghi, 1990) (Figure 1.6.4).

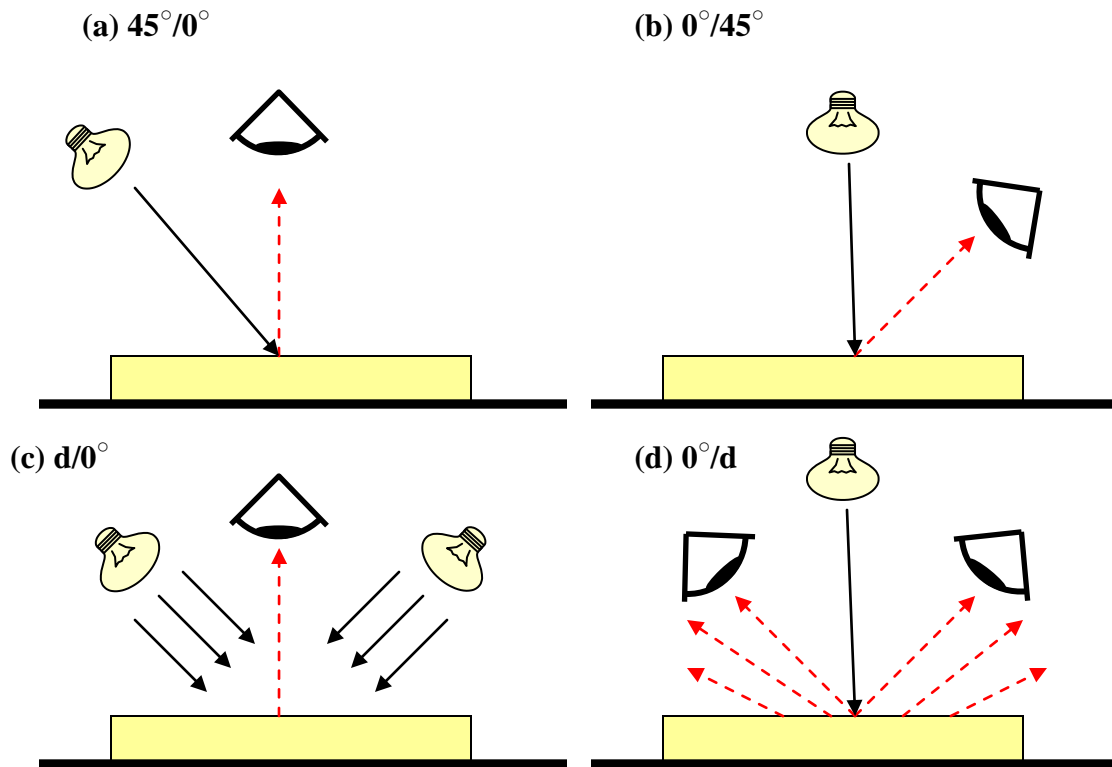


Figure 1.6.4: The illuminating/viewing configurations as recommended by CIE for materials with diffuse reflecting surfaces. (a) and (b) are illuminated by a standard illuminant and detected at 45° . (c) and (d) are illuminated by a standard illuminant and detected at 0° .

The CIE $L^*a^*b^*$ (CIELAB) measuring system is the most frequently used numerical quantification system of colour (Stober et al., 2001; Guler et al., 2005; Lee and Powers, 2005). L^* , a^* and b^* refer to the three coordinates which are taken during the measurement, where L refers to white ($L=100$) or black ($L=0$), a^* measures the green (negative value) and red (positive value) spectrum, and b^* measures the blue (positive value) and yellow (negative value) spectrum according to CIE (1976).

As mentioned above, the photoinitiator decomposition and translucency changes during cure will affect the final colour of the material. It is possible to determine the colour change (ΔE) of RBCs if the CIE $L^*a^*b^*$ co-ordinates are measured before and after cure using Equation 1.6.4.

$$\Delta E = \sqrt{(L_2^* - L_1^*)^2 + (a_2^* - a_1^*)^2 + (b_2^* - b_1^*)^2}$$

Equation 1.6.4

where the subscript numbers 1 and 2 represent before and after cure, respectively for each co-ordinate. Subsequently, the acceptability of differences in colour change is a matter of debate in dental literature (Jonston and Kao; 1989; Guler et al., 2005; Lee and Powers, 2005) whereby some authors suggest that a ΔE value less than 3.3-3.7 is clinically acceptable (Johnston and Kao, 1989; Lee and Powers, 2005) and some suggest a ΔE value less than 5.4 is clinically acceptable (Cook and Chong, 1985) for RBCs. This value is lower for resins and a value not exceeding 2 is considered acceptable (Ferracane et al., 1985). Such differences are presumably due to the many factors affecting the perception of colour, which underlines the necessity to quantify and measure colour change to improve clinical acceptability.

1.7 Summary

Since the introduction of RBCs in the 1960's (Bowen, 1962) there have been significant developments, which continue to improve both mechanical and physical properties, with the ultimate goal of improved clinical longevity. Although RBCs are becoming ever popular with increasing clinical indications, many inherent limitations remain, which limits their application. Polymerisation shrinkage (Davidson and Feilzer, 1997) and limited cure depths (Fan et al., 2002) of RBCs remain major drawbacks when compared to amalgam fillings. Much research has been conducted in order to try and improve the setting reaction through improved technology and chemistry, but polymerization shrinkage and limited depth of cure remain major concerns of clinicians when using RBCs.

1.7.1 Aims of the present investigation

The kinetics of polymerisation is complex and is dependent upon several factors including monomer composition (Feng and Suh, 2007), photoinitiator chemistry (Emami and Soderholm, 2005; Ogunyinka et al., 2007), filler content and irradiation protocol (Musanje and Darvell, 2003; Asmussen and Peutzfeldt, 2005) which all affect the setting reaction. The application of the so-called “exposure reciprocity law” seems an interesting approach to control polymerization kinetics and polymer properties. Consequently, to better understand polymerization kinetics, the associated setting reaction and the applicability of the exposure reciprocity law, Chapter 2 aims to investigate anomalies of the law in high viscosity commercial materials and their low viscosity counterparts (Section 2.1). However, the analysis of commercial materials remains limited due to the fact that manufacturers typically do not divulge the exact composition of RBCs, and thus the investigation of the exposure reciprocity law and the setting reaction in experimental materials for further understanding will also be conducted (Section 2.2). Additionally, the chemistry of the photoinitiator (i.e. the type and concentration) will play a remarkable role in the kinetics of polymerization and the setting reaction and therefore the impact of photoinitiator type in filled and unfilled photoactive resins on the exposure reciprocity law will also be investigated (Section 2.3).

Additionally, as the optical properties of RBC are not well understood, Chapter 3 aims to improve the understanding of such properties within RBCs. The composition of the RBC, i.e. photoinitiator, pigments/dyes, resin matrix and filler fundamentally control any optical change during cure. The photoinitiator is the primary absorber of light and its efficiency is dependent upon its absorption characteristics and will ultimately affect reaction kinetics. However, the use of other photoinitators such as Lucirin TPO are increasing in popularity mainly due to aesthetic benefits, therefore to better understand the use of Lucirin TPO in RBCs, the absorption characteristics of Lucirin TPO will also be investigated with regards to conversion, kinetics, physical properties and its potential for improved aesthetics (Section 3.1).

Limited cure depths are primarily associated with inefficient light transmission at depth. Optical change is mainly governed by Beer-Lamberts law, however, this is

complicated by the fact that there is a complex interaction between photoinitiator decomposition, shrinkage due to monomer conversion and an increase in refractive index due to increase in density by cross linking of polymer chains. As the monomer cures, the refractive index increases and maximum light transmission is expected at the point where refractive index of the monomer matches that of the filler (Shortall et al., 2008). However, the change in refractive index is dependent upon the resin matrix composition and is not well understood dynamically. Therefore to better understand dynamic changes in refractive index the simultaneous dynamic monitoring of refractive index change through and physical thickness change of photo-active resins will be demonstrated using a novel bespoke low coherence interferometer (Section 3.2) with the goal of a better understanding of optical and physical properties. Additionally the affect of shrinkage strain on light transmission will also be investigated (Section 3.3), with the aim of improving the setting reaction through better understanding of light transmission during light curing of various specimen geometries

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CHAPTER 2

Polymerisation Kinetics of Resin Based Composites

2.0 Polymerisation Kinetics of Resin Based Composites

Resin based composites (RBCs) have become the material of choice in restorative dentistry for the direct restoration of damaged teeth. Although the materials have been significantly improved since their introduction several decades ago, several limitations remain. Most notably, inadequate cure depth and polymerisation shrinkage continue to be of primary concern to clinicians. From a material development viewpoint, such problems have been addressed by investigating various material aspects such as inorganic filler morphology and content, resin formulation and curing light technology, with some success. Clinically, RBCs are placed incrementally and should be cured according to manufacturer's guidelines for the material in an attempt to overcome the potential negative consequences of material shrinkage and inadequate cure depth. The incremental layering technique employed is designed to allow for adequate cure of a given increment and also to reduce bulk contraction stress of the composite restoration (Mackenzie et al., 2009). However, for large restorations this can be a complex time consuming procedure and is undesirable both from the viewpoint of operator efficiency and patient comfort. Consequently, the goal of reduced chairside procedure time of large resin composite restorations requiring multiple increments is pursued globally. Accordingly, several investigators and manufacturers of light curing units support the "exposure reciprocity law" and claim that the radiant exposure (the product of irradiance and time) is the main determining factor of degree of conversion and mechanical properties of photoactive RBCs (Sakaguchi and Ferracane, 2001; Halvorson et al., 2002; Price et al., 2004; Emami et al., 2006). Consequently, due to the supposed reciprocity between irradiance and time, some dentists use high-power light curing units to reduce curing exposure. However, failures in the law are reported and it has been suggested that irradiance and time independently influence polymer chain length, extent of cross linking and mechanical properties (Musanje and Darvell, 2003; Asmussen and Peutzfeld, 2005; Peutzfeldt and Asmussen, 2005; Dewaele et al., 2009). Therefore, the aim of the current Chapter is to gain a greater insight into such failures by further investigating the physical properties of high intensity curing.

2.1 High Irradiance Curing and the Applicability of the Exposure Reciprocity Law in Commercial Dental Resin-Based Materials*

2.1.1 Abstract

The aim of this work was to investigate the effect of high irradiance curing on degree of conversion (DC) of ‘flowable’ resin composites and their counterpart higher viscosity restorative materials. Five commercial flowable materials (Venus; Heraeus Kulzer, Synergy D6; Coltene, Premise; Kerr, Grandio; Voco and Gradia; GC Corp) and their counterpart restorative versions were tested. Specimens were cured with a halogen Swiss Master Light (EMS, Switzerland) at similar radiant exposures (18 J/cm^2): 400 mW/cm^2 for 45 s, 900 mW/cm^2 for 20 s, 1500 mW/cm^2 for 12 s, 2000 mW/cm^2 for 9 s or 3000 mW/cm^2 for 6 s. DC was measured in real time by Fourier transform near infrared spectroscopy (FT-NIRS). Two way ANOVA testing revealed significant differences ($P \leq 0.02$) with respect to “composite type” and “cure protocol” for DC for all 5 product comparisons. Supplementary one-way ANOVA also revealed significant differences between curing protocols ($P < 0.05$). The majority of higher viscosity resin composite paste materials exhibited similar DC regardless of curing protocol. However, a significant decrease in DC for specimens cured at 3000 mW/cm^2 for 6 s compared with 400 mW/cm^2 for 45 s was observed for the flowable materials, Grandio (41 ± 0.36 and 62 ± 1.15 %, respectively) and Venus (44 ± 0.44 and 67 ± 0.44 %, respectively). Conversely, other flowable materials exhibited little or no significant differences between curing modes. Generally, flowables showed higher DC than their counterpart, except at high irradiance for those materials were reciprocity failed. The validity of exposure reciprocity law and final DC depends on several factors, among which resin viscosity and filler content were important. Information on material composition and appropriate radiation sources by manufacturers may assist practitioners with the selection of appropriate curing protocols for specific material / light curing unit combinations which may reduce the incidence of under-cured restorations and the clinical impact thereof.

* Hadis M, Leprince JG, Leloup G, Devaux J, Shortall AC, Palin WM. High irradiance curing and anomalies of exposure reciprocity law in resin-based materials. Submitted to Journal of Dentistry.

2.1.2 Introduction

The placement of multi-increment resin composite restorations is time consuming and there exists a significant demand for reducing cure time to minimize chairside procedure and its potential financial implication (Christensen, 2000). Therefore, the use of high power curing light sources is commonplace since they play a crucial role in polymerization efficiency. Some reduction in curing time can be achieved by optimizing the spectral overlap between the curing unit and the photoinitiator to increase the efficiency of photon absorption, which has been improved in the case of camphorquinone initiated materials by the development of blue gallium nitride LED technology (Krämer et al., 2008; Jiménez-Planas et al., 2008; Santini, 2010). The general trend amongst curing unit manufacturers has been to produce lights with increased irradiance with the presumption that curing time can be reduced proportionally. This is primarily based on the assumption that the degree of conversion and mechanical properties of photoactive materials are dependent upon the total energy of the irradiation, i.e. radiant exposure (RE , J/cm^2), the product of irradiance (I , mW/cm^2) and irradiation time (t , s).

For photo-active dental materials, the total energy principle promotes the common assumption that varying combinations of curing irradiance and exposure time provide similar material properties at constant radiant exposure. This is the principle known as the “exposure reciprocity law” and holds quite well in certain cases with regards to degree of conversion (Halvorson et al., 2002; Emami and Söderholm, 2003), modulus (Sakaguchi and Ferracane, 2001) and hardness (Price et al., 2004). Other studies have reported that although radiant exposure plays an important role, irradiance and time independently influence polymer chain length, extent of crosslinking and mechanical properties, and consequently exposure reciprocity does not hold under all conditions (Musanje and Darvell, 2003; Peutzfeldt and Asmussen, 2005; Asmussen and Peutzfeldt, 2005; Feng and Suh, 2007; Dewaele et al., 2009).

The kinetics of polymerisation is complex and is dependent upon several factors including monomer composition (Feng and Suh, 2007), photoinitiator chemistry (Cook, 1992; Emami and Söderholm, 2005), filler content and irradiation protocol (Asmussen and Peutzfeldt, 2005; Feng and Suh, 2007). For example, the investigation by Musanje

and Darvell (2003) on highly filled commercial materials highlighted the importance of particular composition and reactivity of different commercial products. They tested four commercial products at different combinations of irradiance and time and examined flexural strength modulus and total energy required for failure and reported that out of the four tested materials only one satisfied the exposure reciprocity law for all the properties considered (Musanje and Darvell, 2003). A recent study by Feng & Suh (2009) using a plasma-arc and a halogen curing unit suggested that some RBCs and adhesives exhibit inferior cure using short exposure times at high irradiance (Feng et al., 2009). The same group also highlighted the importance of resin viscosity on degree of conversion and the validity of the exposure reciprocity law (Feng and Suh, 2007). Other studies have also highlighted the inferior properties of lower viscosity systems (dentin bonding adhesives) compared to higher viscosity resin composites when using high power light curing units compared to other lights (D'alphino et al., 2007).

Clinically, the impact of viscosity might have important consequences on the extent of cure of light-cured materials in dental practice. It is problematic when commercial materials are used where exact compositions remain unknown. Notably, commercial flowable materials can either be produced by reducing the filler content or increasing the diluent monomer, either of which may affect the setting reaction and final polymer conversion for high intensity curing protocols. As a result, this work aims to compare the efficiency of different curing protocols in terms of degree of conversion in flowable composite resin materials and their highly filled counterparts by testing two hypotheses:

- i) If exposure reciprocity law is upheld for the 'parent' restorative composite (ie. similar degree of conversion regardless of curing protocol) then it will also hold for its counterpart lower viscosity, 'flowable'.
- ii) An increase in the filler content will result in a decrease in the degree of conversion regardless of curing mode.

2.1.3 Materials and Methods

Five commercial resin composite materials (Venus; Heraeus Kulzer, Synergy D6; Coltene, Premise; Kerr, Grandio; Voco and Gradia; GC Corp) and their counterpart flowable versions in the A3 shade were tested for their compatibility with exposure reciprocity law. Table 2.1.1 describes the commercial materials and their resin matrix constituents. Specimens were cured with a halogen Swiss Master Light (EMS, Switzerland; 11 mm diameter light guide) chosen for its capability to provide a range of irradiance outputs and exposure times. Spectral emission was measured using a UV-Vis spectrometer (USB4000, Ocean Optics) calibrated with a deuterium tungsten light source (Mikropack DH2000-CAL, Ocean Optics, Dunedin, USA) with known spectral output in the UV-Vis and NIR range in accordance to NIST standards. This allowed absolute measurements of irradiance for the light curing unit at different irradiances. Reflective neutral density filters (Thorlabs, UK) were stacked to control the output of the high power curing light which produced a 6-fold reduction in irradiance over the whole spectrum during absolute irradiance measurements. Five different curing protocols giving a total of 18 J/cm² of energy were applied to compare the applicability of exposure reciprocity law; 400 mW/cm² for 45 s was chosen as the reference protocol and compared to protocols of the same radiant exposure (18 J/cm²): 900 mW/cm² for 20 s, 1500 mW/cm² for 12 s, 2000 mW/cm² for 9 s and 3000 mW/cm² for 6 s. The irradiance values were based on the in-built radiometer of the curing-unit and verified by the spectrometer. Real time measurement of degree of conversion (*DC*) and rate of polymerization (R_p) were measured using a method developed by Stansbury and Dickens by which a Fourier Transform near infrared Spectrometer (FT-NIRS; Figure 2.1.1) continuously monitors (integration time = 0.85 s) the decomposition of the -CH₂ peak at 6164 cm⁻¹ (Stansbury and Dickens, 2001). Curing kinetics of the RBC samples held in a 12 x 1.4 mm cylindrical white Teflon mould placed on a glass slide and covered by a thin glass cover slip (n=3) were measured for 190 s. Although *DC* evolves for up-to 24 h post-irradiation (Truffier-Boutry et al., 2006), the major part of polymerization occurs during light activation. For that reason, *DC* at 180 s after the start of measurement was chosen for comparison and will be referred to as DC^{Max} .

	PASTE			FLOWABLE		
	Resin	Filler mass % (manufacturer)	Filler mass % (Ashing)	Resin	Filler mass % (manufacturer)	Filler mass % (Ashing)
Venus (Heraeus)	Bis-GMA based	78	73.3 (0.6)	Bis-EMA, TEGDMA	62	61.6 (0.1)
Synergy D6 (Coltene)	Bis-GMA, TEGDMA, UDMA	80	74.0 (0.2)	Bis-GMA, TEGDMA	63	57.9 (0.1)
Premise (Kerr)	Bis-EMA, TEGDMA	84	73.8 (0.3)	Bis-EMA, TEGDMA	79	58.7 (0.9)
Grandio (Voco)	Bis-GMA, TEGDMA	87	84.0 (0.2)	Bis-GMA, TEGDMA, HEDMA	80	76.7 (0.1)
Gradia (GC)	UDMA-based	77	63.3 (0.2)	Dicarbamate/TEGDMA	60	45.5 (0.0)

Table 2.1.1: Resin Matrix Composition and filler mass percentage (reported by the manufacturers) and mass percentage (experimentally determined) of commercial materials.

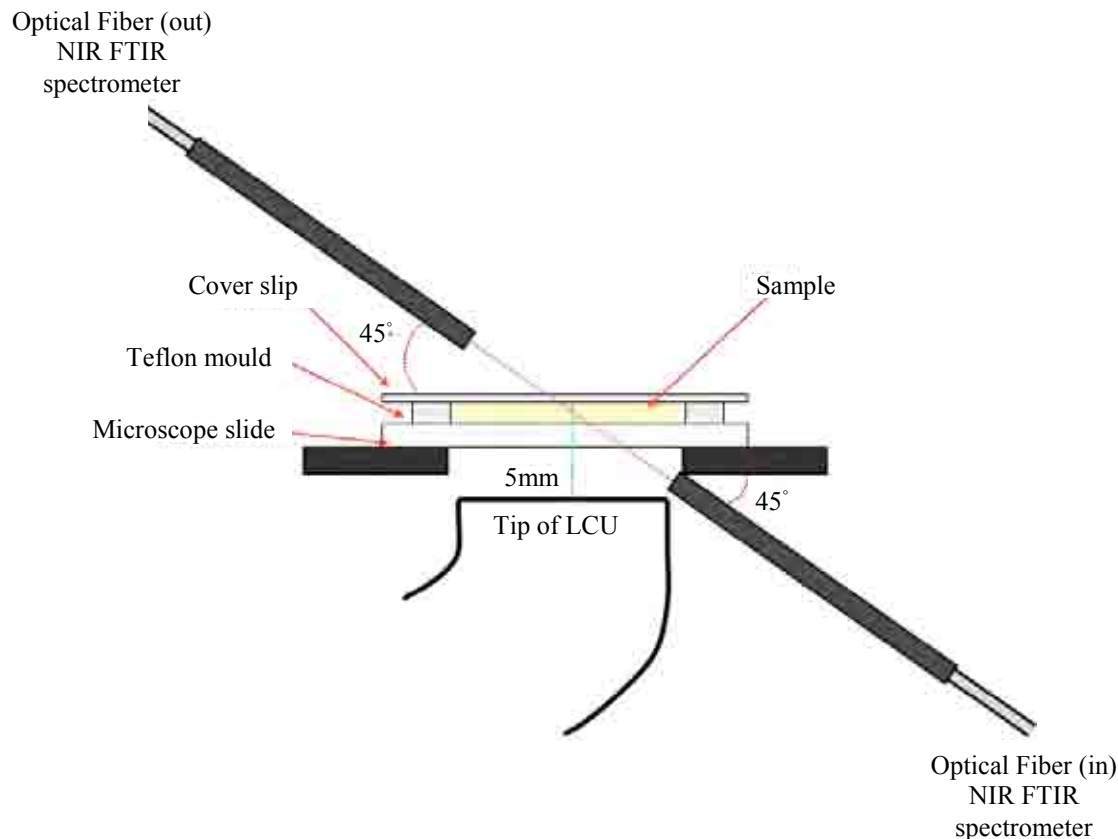


Figure 2.1.1: Experimental setup for real time FT-NIRS measurements.

Additionally, temperature rise measurements were made for the longest and the shortest irradiation modes (i.e. 400 mW/cm^2 for 45 s and 3000 mW/cm^2 for 6 s) to compare heat energy received by the curing composites from the light source at the two extremes. The thermistor bead set up consisted of a miniature bead thermistor (256-045 RS Components, Corby, England), which was imbedded into a brass jacket that was used as a stage. The recess between the brass jacket and the bead was filled with high density polysynthetic silver thermal compound (Arctic Silver 5, Visalia, CA) and then a piece of aluminium foil (BacoFoil, UK; thickness = $0.0182 \pm 0.0008 \text{ mm}$) was placed on top to aid conductivity and protect the bead from damage. Rubber alignment rings were used to ensure concentric alignment of the cavity and thermistor bead before filling the cavity (similar dimensions to the FT-IR measurements) with model composites directly onto the aluminium foil. The composites were either made with 70 wt% Bis-GMA and 30 wt% TEGDMA or 30 wt% Bis-GMA and 70 wt% TEGDMA and contained the photoinitiator/amine camphoroquinone and dimethylaminoethyl methacrylate (0.2 wt%

and 0.8 wt% respectively). The reagents were supplied by Sigma Aldrich, UK and were used as received. The resins were filled to 40 vol% with 3 μm barium glass filler (Schott, Germany) and 14 nm fumed silica particles (Aerosil 150, Evonik Industries, Germany) at a volume percent of 30 and 10%, respectively. This enabled the comparison of a high viscosity version and a low viscosity version on the amount of heat received. A cover slip (0.15 mm thick, 22 mm diameter) was pressed on top and the excess was extruded and scraped away before curing the samples ($n=3$) from above the cover slip. The heat evolved during curing was measured continuously (10 s^{-1}) and logged using bespoke software. Additionally, as reaction exotherm will affect the temperature rise, the specimens were cured several times until the temperature no longer increased. This was used as the control to allow an approximation of the heat delivered to the material due to light minus the reaction exotherm.

The measurements of viscosity were performed at a room temperature ($23 \pm 1^\circ\text{C}$) using a Physica MCR301 rheometer (Anton Paar Inc., Austria) with 25 mm parallel plates separated from each other by 0.5 mm. The complex viscosity (Pa.s) was measured at a constant strain amplitude of 0.06 and at two different angular frequencies: 0.01 and 10 rad/s. Three measurements were performed for each material tested.

As exact composition of commercial materials are unknown, the filler mass percentage of commercial materials were determined by an ashing technique which compares the difference in weight before and after ashing of 0.5 g ($n=3$) of each material at 900°C . The crucible was introduced into the furnace (Carbolite CWF 1300, UK; Max temperature: 1300°C) for an hour once the temperature of the oven had reached 900°C . The porcelain crucibles were weighed using an analytical balance which was accurate to 0.001 g (Acculab Atilon, Satorius Group) before and after ashing.

Statistical analyses were performed on the DC^{Max} , R_p^{Max} and time to R_p^{Max} data sets for the commercial materials on the ‘flowable’ and ‘restorative’ resin composites using two-way analysis of variance (ANOVA) with composite type (2 levels) and curing protocol (5 levels) as the independent variables. Additionally, one-way ANOVA and post-hoc Tukey tests ($p=0.05$) were performed on DC^{Max} , R_p^{Max} and time to R_p^{Max} data sets versus the cure protocol to confirm the validity or otherwise of the exposure reciprocity law (Table 2).

2.1.4 Results

Figure 2.1.2 represents a comparison between the emission spectra of the curing light at each irradiance against the absorption spectra of camphoroquinone. The increase in irradiance resulted in a greater spectral overlap with camphoroquinone although the spectral emission range for the light remained comparable at low irradiance.

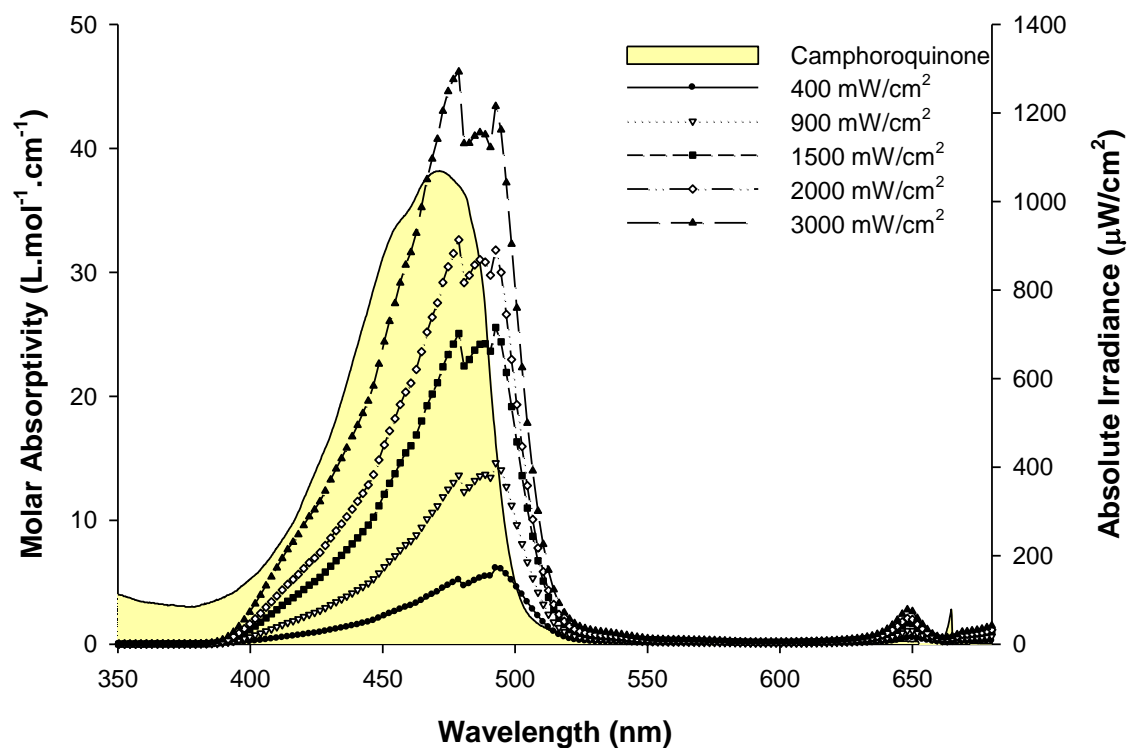
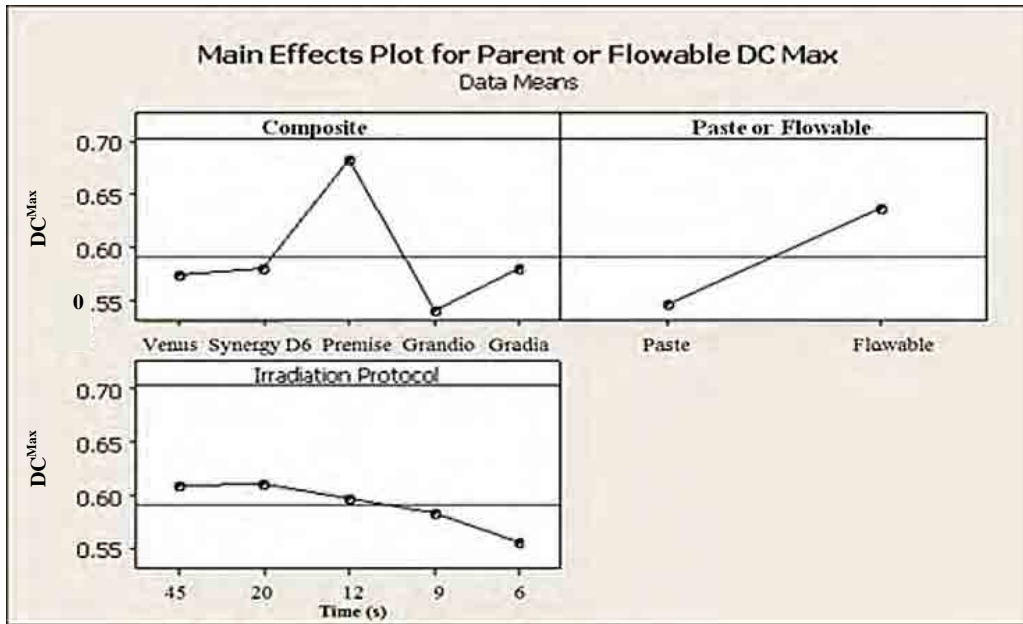


Figure 2.1.2: Absorption spectra (Left axis, in $\text{L.Mol}^{-1}\text{cm}^{-1}$) of CQ (Light yellow fill) compared to emission spectra of the Swiss Master Light (right axis, in $\mu\text{W}/\text{cm}^2.\text{nm}$) at different irradiances, i.e. 400, 900, 1500, 2000 and 3000 mW/cm^2 .

Two-way ANOVA revealed significant differences ($P \leq 0.05$) with respect to “composite type” and “cure protocol” for DC^{Max} for all 5 product comparisons. For any product that did not obey exposure reciprocity law, DC^{Max} was significantly greater for the low irradiance and longest exposure curing protocols (400 mW/cm² for 45 s) ($P < 0.05$; Figure 2.1.3). Generally R_p^{Max} increased and time to R_p^{Max} decreased as higher curing intensities were employed (Table 2.1.2).

a)



(b)

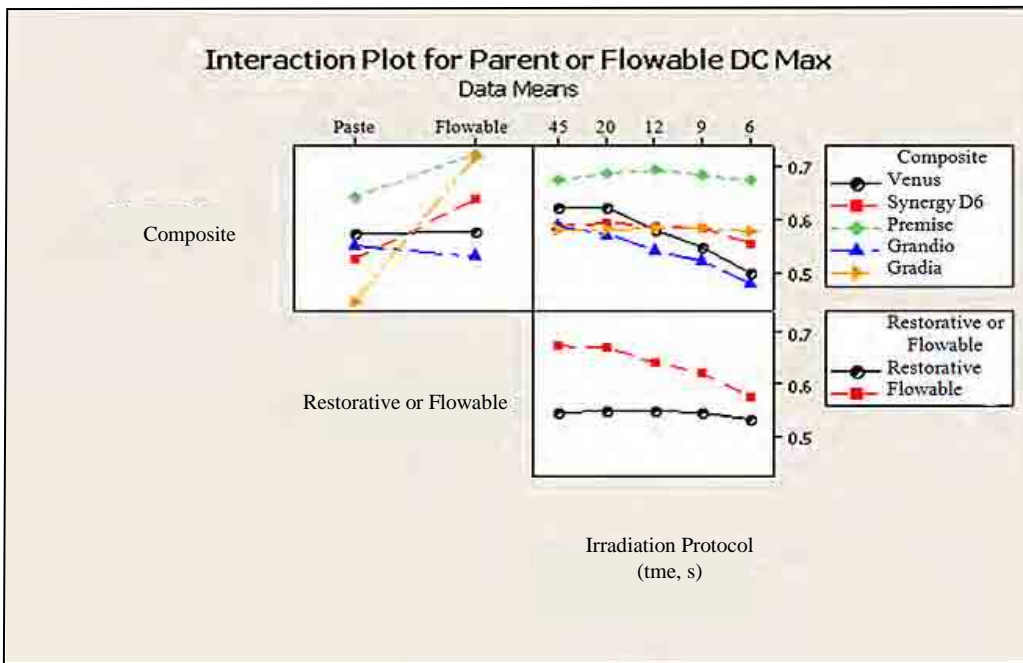
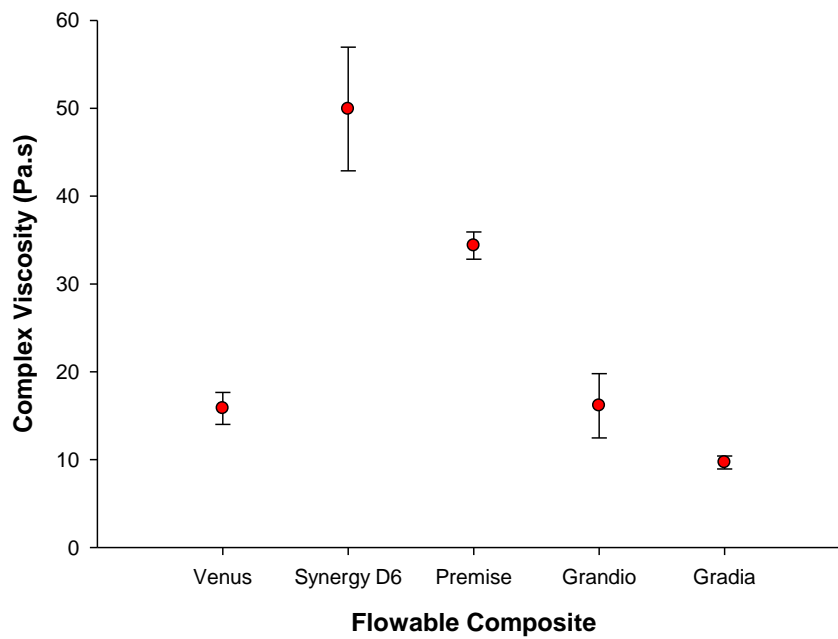


Figure 2.1.3: The main affects plot for manufacturer, composite type (restorative or flowable) and irradiation protocol for the commercial materials on DC^{Max} (a). The interactions of manufacturer, composite type and irradiation protocol of commercial materials on DC^{Max} (b).

For the ‘flowable’ materials, DC^{Max} was significantly reduced for the high irradiance protocol in all cases ($P < 0.05$; Table 2.1.2). DC^{Max} of Venus Flow and Grandio Flow were more dependent upon cure protocol and exhibited a significantly decreased DC^{Max} even at intermediate compared with low intensity curing regimes ($p < 0.05$; Table 2.1.2). Figure 2.1.4 represents the complex viscosities for the flowable composites at high and low angular frequencies. It was noted that although some flowables exhibited similar or even reduced viscosity at different angular frequencies (Gradia at 0.1 s^{-1} , Synergy D6 and Gradia at 10 s^{-1}) compared with Venus and Grandio, DC^{Max} of the former material types was less dependent upon curing protocol (Figure 2.1.3; Table 2.1.2).

(a)



(b)

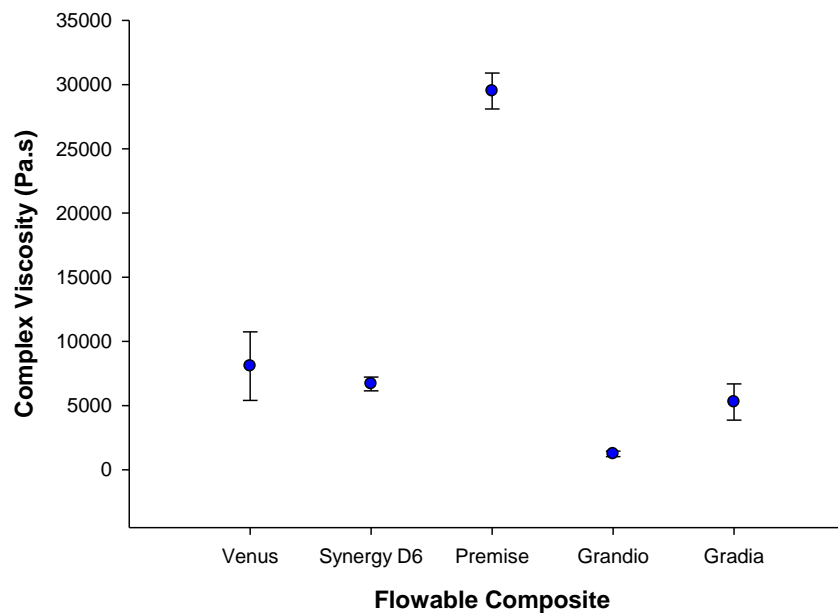


Figure 2.1.4: The complex viscosities of the flowable composites. The error bars represent the standard deviations (n=6) for each flowable material, angular frequency of 10 rad/s and (a) angular frequency of 0.01 rad/s (b).

For high viscosity resin composites no significant differences in DC^{Max} were recorded regardless of the curing protocol ($p>0.05$) except for Venus, which exhibited a significantly reduced DC^{Max} at the highest intensity curing regime ($p<0.01$, Table 2.1.2). In most cases DC^{Max} of the paste materials was significantly lower than that recorded in their flowable counterparts ($p<0.05$; Figure 2.1.3). However, for intermediate and high curing intensities, DC^{Max} of Grandio and Venus pastes were either similar or greater than the associated respective flowable material (Table 2.1.2).

	CURE PROTOCOL		DC ^{MAX} (%)		R _p ^{MAX} (%/s)		Time to R _p ^{MAX} (s)	
Resin Composite	Time (s)	Irradiance (mW/cm ²)	PASTE Mean (+/- SD)	FLOWABLE Mean (+/- SD)	PASTE Mean (+/- SD)	FLOWABLE Mean (+/- SD)	PASTE Mean (+/- SD)	FLOWABLE Mean (+/- SD)
Venus (Heraeus)	45	400	57 (0.10) a	67 (0.44) a	4.03 (0.09) d	2.99 (0.23) d	6.5 (0.5) a	11.9 (1.5) a
	20	900	58 (0.10) a	66 (1.32) a	6.39 (0.63) c	4.43 (0.33) c	4.0 (0.5) b	8.2 (0.5) b
	12	1500	58 (0.31) a	58 (0.70) b	7.58 (0.33) bc	4.77 (0.10) c	3.7 (0.5) b	6.2 (0.5) bc
	9	2000	57 (0.02) ab	53 (0.28) c	8.52 (0.20) b	5.30 (0.05) b	3.4 (0.0) b	4.8 (0.5) c
	6	3000	55 (0.26) c	44 (0.44) d	9.91 (0.83) a	5.91 (0.05) a	2.8 (0.5) b	4.3 (0.0) c
Synergy D6 (Coltene Waledent)	45	400	53 (1.16) a	64 (0.24) a	3.44 (0.11) c	4.42 (0.17) d	6.2 (1.3) a	4.5 (1.0) a
	20	900	53 (0.79) a	65 (0.47) a	6.57 (0.81) b	7.70 (0.97) c	2.3 (1.0) b	2.0 (1.3) b
	12	1500	53 (2.59) a	64 (0.93) a	10.51 (1.47) a	10.53 (0.60) b	1.4 (0.5) b	0.9 (0.0) b
	9	2000	53 (1.29) a	64 (0.39) a	12.13 (0.31) a	12.11 (0.67) ab	0.9 (0.0) b	0.9 (0.0) b
	6	3000	49 (1.48) a	61 (1.19) b	12.40 (0.89) a	13.78 (1.10) a	0.9 (0.0) b	1.1 (0.5) b
Premise (Kerr)	45	400	63 (1.19) a	72 (0.49) a	8.25 (0.76) d	7.65 (0.81) d	4.5 (1.3) a	3.4 (2.6) a
	20	900	65 (0.14) a	73 (0.15) a	11.96 (0.50) c	12.94 (1.00) c	3.4 (0.0) ab	2.0 (1.0) a
	12	1500	65 (0.79) a	74 (1.33) a	15.71 (0.50) b	15.99 (0.97) b	2.3 (0.5) b	1.7 (0.9) a
	9	2000	64 (1.37) a	73 (0.83) a	18.01 (1.07) a	18.01 (1.07) a	1.7 (0.9) b	1.7 (0.9) a
	6	3000	64 (1.22) a	71 (1.41) b	19.95 (1.71) a	20.00 (1.66) a	2.3 (0.5) b	2.0 (0.5) a
Grandio (Voco)	45	400	55 (0.35) a	62 (1.15) a	5.19 (0.13) e	2.97 (0.07) c	4.0 (0.5) a	9.1 (0.5) a
	20	900	55 (0.39) a	59 (0.74) b	8.17 (0.06) d	4.35 (0.19) b	2.6 (0.0) b	4.8 (0.5) b
	12	1500	55 (0.43) a	53 (0.76) c	10.19 (0.43) c	5.07 (0.23) a	2.6 (0.0) b	2.6 (0.0) c
	9	2000	55 (0.09) a	49 (0.67) d	11.40 (0.44) b	5.48 (0.31) a	2.0 (0.5) bc	2.8 (0.5) c
	6	3000	54 (0.41) a	41 (0.36) e	13.80 (0.65) a	5.11 (0.12) a	1.7 (0.0) c	3.1 (0.5) c
Gradia (GC Corp)	45	400	44 (0.86) a	72 (0.45) a	2.97 (0.34) d	7.23 (0.20) d	6.8 (0.9) a	4.5 (0.5) a
	20	900	44 (0.24) a	72 (0.08) a	5.09 (0.17) c	9.24 (0.10) c	4.3 (0.0) b	4.0 (0.5) ab
	12	1500	44 (0.30) a	72 (0.32) a	7.52 (0.60) b	11.59 (0.40) b	2.8 (0.5) c	3.4 (0.0) bc
	9	2000	45 (0.12) a	72 (0.25) a	10.12 (1.06) a	12.48 (0.59) b	2.0 (0.5) c	3.1 (0.5) bc
	6	3000	44 (0.37) a	71 (0.32) b	10.86 (0.41) a	15.80 (0.74) a	2.0 (0.5) c	2.3 (0.5) c

Table 2.1.2: DC^{MAX} , R_p^{MAX} and Time to R_p^{MAX} of commercial materials. Similar letters within boxes indicate no significant differences. (p<0.05).

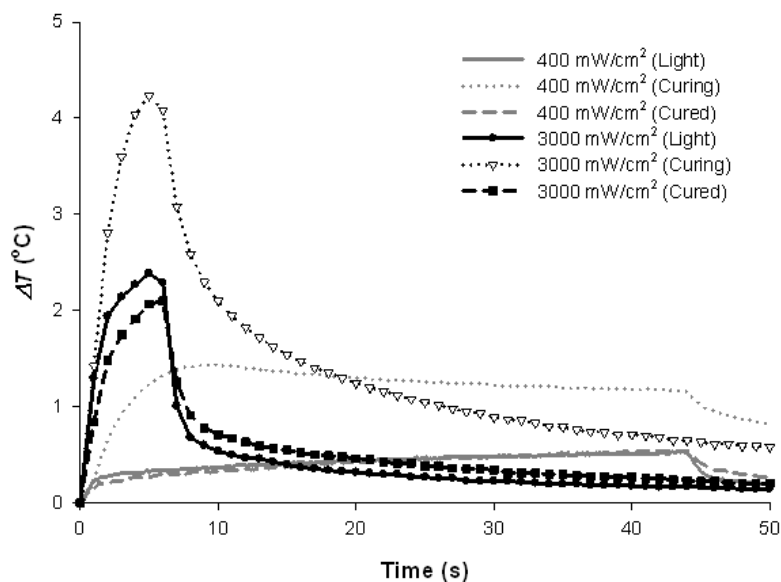
Following ashing, the majority of materials showed a decreased filler mass % in flowable materials compared with their parent restorative versions, although the difference observed between Grandio products was far less pronounced (Table 2.1.1).

Table 2.1.3 represents the total heat energy registered at the thermistor bead, calculated from area under the respective curves in Figure 2.1.5 which represents the temperature measurements in model composites and the temperature rise due to the light at the highest and lowest irradiation modes. In both cases the high irradiation mode produced a significantly higher temperature rise regardless of resin matrix than the lower irradiation mode for all three series ($P < 0.05$). The lower irradiation mode gives a more sustained temperature rise throughout the irradiation period and therefore the total heat energy at the thermistor is greater than for the short irradiation mode (Table 2.1.3).

Resin Matrix	Irradiation Series	Area Under Curves (°C)	
		45s at 400 mW/cm ²	6s at 3000 mW/cm ²
70/30	Curing specimen	53.2 (3.7)	17.8 (2.1)
	Light	18.6 (0.8)	10.9 (1.3)
	Cured Specimen	18.0 (1.6)	8.9 (0.8)
30/70	Curing specimen	74.6 (4.7)	17.3 (0.8)
	Light	17.7 (2.0)	12.5 (1.4)
	Cured Specimen	17.9 (3.8)	10.7 (0.5)

Table 2.1.3: Total heat energy detected by the thermistor during irradiation calculated as the integrals of the curves respectively (Figure 2.1.5).

(a)



(b)

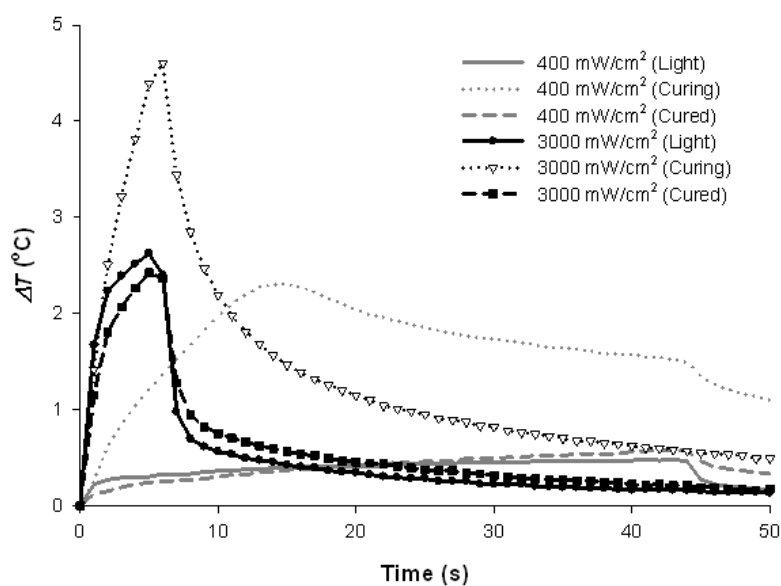


Figure 2.1.5: ΔT during polymerization (curing) at low irradiance (45s at 400 mW/cm²) compared to high irradiance (6s at 3000mW/cm²) for high viscosity resins (a) and low viscosity resins (b) filled with 40 vol% filler in comparison to the light on its own (light) and temperature rise in a fully cured specimen (cured).

2.1.5 Discussion

The first hypothesis (if exposure reciprocity law is upheld for a paste, it will also hold for its respective flowable material) was rejected as for all pastes reciprocity was achieved or nearly achieved (Venus), while no such relationship was observed for the respective flowables with much larger differences in DC^{Max} of Grandio and Venus across curing regimes. An interesting point to note is that flowable composites are manufactured in different ways, i.e. some have more diluent monomer and some have less filler percentage in comparison to their parent high viscosity versions (the effects of which will be investigated in Section 2.2). Nonetheless, previous work has indicated that above a specific (unfilled) resin viscosity ($\geq 6:4$, Bis-GMA/TEGDMA), exposure reciprocity law is upheld (Feng and Suh, 2007). Viscosity of the resin matrix is known to play a significant role in the polymerization rate, with the highest maximum rates observed for mixtures containing 50-75 wt% Bis-GMA (Emami and Söderholm, 2003). Consequently, it may be expected that flowables composites manufactured with $\geq 6:4$, Bis-GMA/TEGDMA or high viscosity resins may allow for reduced curing time clinically when high irradiance protocols are used. However, when manufacturers use low viscosity resins, it is expected that high system mobility will increase the probability of a bimolecular (radical-radical) termination pathway during the initial stages of polymerization resulting in a greater loss of radical growth centres by early termination. When high irradiance protocols are applied to composites containing such resins, radical loss is expected to be greater as a multitude of radical growth centres exist and as termination is proportional to the squared concentration of free radicals at low conversion (Leprince et al, 2010a). Consequently DC^{Max} will significantly decrease compared with similar radiant exposures achieved using a lower irradiance for longer time for composites containing low viscosity resins. Therefore, a correlation between the applicability of the exposure reciprocity law and the viscosity of the flowable materials may have been expected. However, this was not the case since Gradia exhibited the lowest viscosity (measured at 10 rad/s; Figure 2.1.4) and achieved similar DC^{Max} for most curing regimes whereas DC^{Max} for the more viscous Grandio and Venus flowable materials was clearly more dependent upon the irradiance and curing time combination (Figure

2.1.4 and Table 2.1.1). Although at lower angular frequency the rank order of viscosity changed (presumably due to differences in pseudoplastic behaviour), similar effects were seen. Here, Gradia and Synergy D6 exhibited lower viscosity compared with Venus, but the former still obeyed the exposure reciprocity law for most cure modes (Figure 2.1.1 and Table 2.1.2). Therefore, it is evident that the impact of viscosity on composite reactivity is more complicated than that observed for unfilled resins.

Broadly speaking, flowable resin composites can be manufactured either by reducing the filler mass percentage or by increasing the volume of diluent resin, maintaining a high filler load. In most cases, manufacturers do not divulge the exact composition of commercial resin composites and typically will only report resin matrix composition and filler percentage approximately. Material ashing may reveal useful information regarding filler content, however such a technique is limited since the exact percentage of organic filler types, which will decompose upon heating, are usually unknown. Since the filler mass percentages of Grandio and Grandio Flow are close (Table 2.1.1), it might be assumed that a higher quantity of diluent resin is used in the formulation of the latter. At constant volume percentage of filler, partially substituting Bis-GMA with a low molecular weight monomer (HEDMA) will serve to decrease the viscosity of the material. Although a reduced viscosity at a constant filler load is not apparent at first glance in the case of Venus/Venus Flow (78/62 and 73/62 mass% reported by the manufacturer and from ashing, respectively; Table 2.1.1). Closer examination of the resin constituents reveals substitution of the bulky Bis-GMA monomer with a lower molecular weight monomer, Bis-EMA. Again, this may serve to decrease resin viscosity in the ‘flowable’ composite and thus reduce DC^{Max} at high irradiance and short irradiation times. Commercial ‘flowable’ materials, which exhibited similar DC^{Max} and a reciprocal relationship with respect to irradiance and time, show a much greater reduction in filler percentage as opposed to any change in the resin matrix components (Table 2.1.1).

In filled systems, the interaction between complex viscosity and the applicability of exposure reciprocity law may not be directly related to bulk or macroscopic viscosity. A first possible explanation could be that macroscopic viscosity does not reflect ‘local’ viscosity of the resin matrix, as a result of inter-particulate spacing between fillers. From

the current results it appears that local viscosity is critical in determining polymerization kinetics of curing photopolymer composites. However, the situation is complex since changes in viscosity in the early stages of polymerization may differ with material composition. The effect of changing rheology throughout cure is also an important consideration and should be subject of further investigations. Additionally, RBC manufacturers are increasingly using other photoinitiator(s), the type(s) or concentration of which remain unknown to some extent. Since photoinitators have different absorption characteristics, i.e. molar extinction coefficients and absorption bandwidths, their quantum yield of conversion vary (Neumann et al., 2005; Neumann et al., 2006). As such, the type and concentration of photoinitiator may influence the reciprocal relationship between irradiance and exposure time, which also warrants further investigation and is presented in Section 2.3.

The second hypothesis (that an increase of filler content results in a decrease in degree of conversion regardless of curing protocol) must also be rejected since the flowable versions of Venus and Grandio exhibited lower DC^{Max} at high irradiance than their respective higher viscosity counterparts. For unfilled, or less filled systems vitrification is delayed, which allows further radical propagation and increased DC is usually observed for curing protocols that employ long exposure duration (D'alphino et al., 2007; Feng and Suh, 2007; Feng et al., 2009). Such an effect explains why all flowables tested here show increased DC^{Max} at the low irradiances for the longest cure times (45 s at 400 mW/cm² and 20 s at 900 mW/cm²; Table 2.1.2). An inferior cure of Grandio and Venus flowable compared with their counterpart paste materials at higher intensities might be explained by the effects of a marked reduction in parent resin viscosity as previously explained. The addition of fillers will increase the local viscosity of the reaction medium and reduce the mobility of the propagating species and therefore more highly filled systems are expected to approach a reciprocal relationship with respect to irradiance and time at constant radiant exposure. This effect is supported by previous findings, which reported a 3-fold increase in higher concentration of trapped free radicals in the organic matrix of a filled composite than in its corresponding parent resin due to spatial confinement within filled composites preventing bi-radical termination (Leprince et al., 2009). The shape, size and content of filler as well as interlocking and interfacial

interactions between filler and the resin matrix will also affect the viscosity and mobility of the reaction medium (Lee et al., 2006). For flowable composite materials with low viscosity resins, the mobility restrictions due to fillers will lessen and a reduced local viscosity will increase the probability of bi-radical termination and DC will become more dependent upon irradiance and time (Table 2.1.2).

The effect of temperature, either from the curing unit, polymer exotherm or insulating effects of filler inclusions should not be neglected. Absorption and scattering of light energy within the volume of composite may cause an increase in temperature during photopolymerisation (Panankis and Watts, 2000). The heat delivered to the material will increase proportionally with increased irradiation time and intensity (Lloyd et al., 1986; Hansen and Asmussen, 1993). An increase in temperature increases mobility, not only of the radicals but also of the whole reaction environment, delays auto-acceleration and increases the reaction rate parameters (an Arrhenius-type dependence) thus increasing polymerization rate and final degree of conversion (Lovell et al., 2003). The lower irradiance curing protocol (45s at 400mW/cm^2) is expected to have a slower polymerization process, which may result in a delayed liberated energy from the material and therefore a lower temperature increase. Furthermore, it has been reported that the use of the high irradiance mode on a Swiss Master light for a short time does not produce a significant temperature increase when compared to a standard halogen or LED light curing unit when activated for a longer time (Uhl et al., 2006) but this should be treated with caution as spectral emissions are not comparable between lights and dissimilar radiant exposures are used. Indeed, complementary temperature measurements recorded here in model RBCs using the Swiss Master Light revealed differences in the maximum temperature rise as well as total temperature rise between low (45s at 400mW/cm^2) and high irradiance (6s at 3000mW/cm^2) curing protocols. Model composites were used as manufacturers may use different monomers, fillers and photoinitiators, which will affect exothermicity during polymerization. The heat measured in a fully cured sample estimates the heating of a composite material due to the light minus the reaction exotherm. The temperature rise was higher at high irradiance, however, a longer cure duration at low irradiance exposes the materials to a more sustained temperature rise resulting in a higher total energy with respect to heat (Figure 2.1.5, Table 2.1.3) and may

contribute to the increased DC^{Max} of materials cured using these protocols. It must also be underlined that the use of experimental materials (without pigments) probably underestimates the heating due to the light because of lower light absorption. Additionally, several limitations of the experimental procedure exist. Firstly, the measurements were carried out under non-isothermal conditions and secondly the protective foil and the conductive paste may have also significantly attenuated the heating effects from the light and exotherm of the curing materials. Further work regarding the synergistic effect of temperature (under isothermal conditions) on curing kinetics during the application of exposure reciprocity law is required to fully elucidate this phenomenon.

DC data should not be considered the “gold-standard” of material properties since it is well known that irradiance and time independently influence polymer chain length and crosslinking, which in turn may affect the mechanical properties such as flexural strength and modulus, water sorption and cure depths (Musanje and Darvell, 2003; Peutzfeldt and Asmussen, 2005; Asmussen and Peutzfeldt, 2005; Ferracane, 2006; Feng and Suh, 2007; Dewaele et al., 2009). Moreover, flowable materials that are manufactured by reducing filler content rather than those which maintain higher filler loads and increase diluent resin concentration will exhibit differences in flexural strength, modulus and polymerization shrinkage, which may affect their performance in situ.

2.1.6 Conclusion

The reduced viscosity of commercial flowables can be acquired in at least two ways; either by decreasing filler load or increasing diluent monomer, which directly affects the cure kinetics and of the effect exposure reciprocity on conversion. From this work, it appears that the validity of the exposure reciprocity law may depend upon the method with which “flowable” viscosity is achieved. This underlines the necessity to adapt curing protocols according to the material filler content as well as the resin matrix composition, which is of prime importance to clinicians. Inevitably, the design of flowable materials may be optimized to ensure equivalent polymer conversion regardless of curing protocols.

2.2 Further Anomalies of the Exposure Reciprocity Law in Model Photoactive Resin Based Materials

2.2.1 Abstract

The aim of this section was to investigate the anomalies of exposure reciprocity in model resin composite materials. Resins containing camphoroquinone (CQ; 0.2 wt%) and dimethylaminoethyl methacrylate (DMAEMA, 0.8 wt%) were prepared having different Bis-GMA/TEGDMA percentages: 70/30; 60/40; 50/50; 40/60 and 30/70 mass%, respectively and were tested either unfilled (UF) or filled to 40 vol% (low fill, LF) or 60 vol% (high fill, HF) with silanated barium glass fillers and fumed silica. Specimens were cured with a halogen Swiss Master Light (EMS, Switzerland) at the same radiant exposures (18 J/cm^2): 400 mW/cm^2 (reference), 900 mW/cm^2 for 20 s, 1500 mW/cm^2 for 12 s, 2000 mW/cm^2 for 9 s and 3000 mW/cm^2 for 6 s. Degree of conversion (DC) was measured by FT-NIRS in real time. For the resins and composites two way ANOVAs for DC^{Max} , R_p^{Max} and time to R_p^{Max} revealed significant differences ($P < 0.001$) in all cases for each of the two independent variables “resin type” and “cure protocol” as well as their interactions. In the case of the UF resin, irradiation protocol was more influential ($df = 4$; $F = 9328$) for DC^{Max} than resin type ($df = 4$; $F = 3455$) whereas the opposite occurred with the LF and HF counterparts. Supplementary one-way ANOVA also revealed significant differences between curing protocols ($P < 0.05$). Only composites made with $\geq 6:4$ Bis-GMA/TEGDMA completely obeyed reciprocity. For the resins and other composites, a significant decrease in DC for specimens cured at 3000 mW/cm^2 for 6 s compared with the reference was observed. Generally the DC was lower for the composites cured with the reference protocol compared to UF resin, but for the higher irradiation mode (3000 mW/cm^2 for 6 s), addition of filler improved conversion. System viscosity as well as the irradiation protocol significantly influence polymerization rate and the onset of reaction-diffusion controlled termination, with the highest maximum rates observed for mixtures containing high viscosity parent resins cured at high irradiance. The design and development of new composites that allow high irradiance curing may reduce curing time, which is desirable to both dentists and patients.

2.2.2 Introduction

Clinically, low viscosity materials or ‘flowables’ have many uses including pit and fissure sealants, base or liners for large restorations, routine Class I and V restorations, small Class III restorations and restoration repairs. The reduction in curing time for such procedures has some inherent advantages for the dental practice, i.e. increased efficiency, as-well as the patient, i.e. comfort. The reduction in curing time has been implemented namely by the application of the exposure reciprocity law (varying combinations of irradiance and time) and the introduction of high powered light curing units (Peutzfeldt and Asmussen, 2005; Feng and Suh, 2007). However, although the law is applicable in some scenarios (Sakaguchi and Ferracane, 2001; Emami and Söderholm, 2003; Price et al., 2004), the kinetics of polymerisation is complex and is dependent upon several factors including monomer composition (Feng and Suh, 2007), photoinitiator chemistry (Cook, 1992; Emami and Söderholm, 2005), filler content and irradiation protocol (Emami and Söderholm, 2005; Dewaele et al., 2009) and as such failures in the law are reported (Musanje and Darvell, 2003; Peutzfeldt and Asmussen, 2005; Emami and Söderholm, 2005; Asmussen and Peutzfeldt, 2005; Dewaele et al., 2009).

Previously, it was shown that some low viscosity commercial flowable composites do not obey reciprocity whilst others do (Section 2.1). Reciprocity failures can be attributed to the existence of bi-radical terminations in relatively low viscosity materials. However, from the previous section it appeared that the matter is complicated as macroscopic viscosity does not reflect local viscosity (i.e. that of the resin). Additionally, the importance of particular composition on reactivity was also highlighted by other researchers using commercial products, where only one out of four products satisfied reciprocity with regards to mechanical properties (Musanje and Darvell, 2003). Section 2.1 and a previous study have reported the inferior properties of some resins and RBCs when cured with high irradiance/short duration cure protocols (Feng et al. 2009). Although previous works are relevant, the conclusions remain limited by the fact that commercial materials (with unknown specific compositions) and/or different light sources to achieve higher intensity (creating differences in emission spectra) are used (Musanje and Darvell, 2003; D’alphino et al., 2007; Feng et al., 2009). Consequently, it

is important to fundamentally understand how certain low-viscosity flowable composites obey exposure reciprocity law whilst others do not (Section 2.1). Therefore, this section of the thesis is intended to investigate and obtain a better understanding of such failures using model photoactive resins and resin based composites. Consequently, the aim of this work is to investigate the following hypotheses:

- i) For similar radiant exposure using the same light curing unit, if exposure reciprocity law does not hold for the parent resin, then it will not hold for the counterpart composite regardless of filler loading.
- ii) An increase in the filler content will result in a decrease in the degree of conversion regardless of curing mode.

2.2.3 Materials and Methods

Resin formulations containing the photoinitiator camphoroquinone (CQ; 0.2 wt%) and the co-initiator, dimethylaminoethyl methacrylate (DMAEMA, 0.8 wt%) were prepared having different Bis-GMA/TEGDMA percentages: 70/30; 60/40; 50/50; 40/60 and 30/70 mass% respectively. Reagents were supplied by Sigma Aldrich, UK and were used as received. The resins were tested either unfilled or filled with silanated barium glass fillers (Average particle size: $3 \pm 1 \mu\text{m}$; Schott, Germany) and fumed silica (14 nm; Aerosil, Evonik Industries, Germany) to 30/10 vol% (Low Fill, LF), and 50/10 vol% (High Fill, HF) respectively. Specimens were cured with a halogen Swiss Master Light (EMS, Switzerland) chosen for its capability to provide a range of irradiance outputs from 400 mW/cm^2 to 3000 mW/cm^2 without altering the spectral emission as described previously. The UV-Vis spectrometer (USB4000, Ocean Optics) was calibrated using a deuterium tungsten light source (Mikropack DH2000-CAL, Ocean Optics, Dunedin, USA) according to NIST standards, as described previously in Chapter 2.1.2. Five different curing protocols giving a total of 18 J of energy were applied to compare the applicability of the exposure reciprocity law, 400 mW/cm^2 for 45 s was chosen as the reference protocol and compared to protocols of the same radiant exposure (18 J/cm^2): 900 mW/cm^2 for 20 s, 1500 mW/cm^2 for 12 s, 2000 mW/cm^2 for 9 s and 3000 mW/cm^2 for 6 s. The irradiance values were based on the in-built radiometer of the curing-unit and

verified by the spectrometer.

The degree of conversion was measured throughout cure using a method developed by Stansbury and Dickens (2001) where a Fourier Transform Infra-red spectrometer is used to continuously measure -CH₂ peak (6164 cm⁻¹). The measurement protocol was identical to the one used in Chapter 2.1.2, where curing kinetics of resins and resin composite samples confined in a white Teflon ring mould were monitored for 180 s (n=3) (Figure 2.1.1).

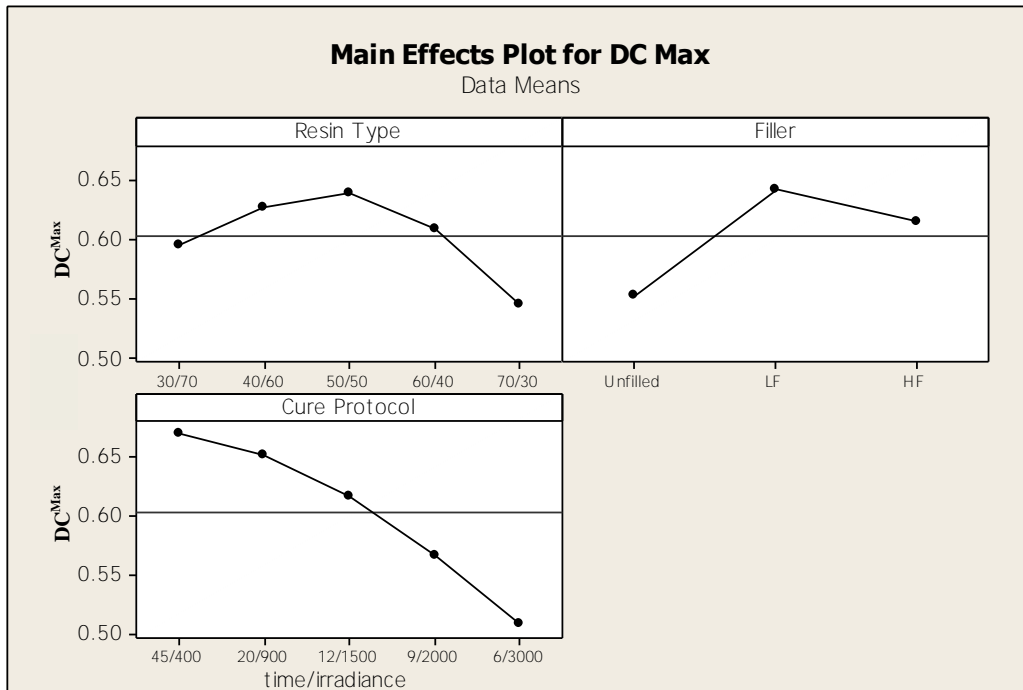
Statistical analyses were performed on the DC^{Max} , R_p^{Max} and time to R_p^{Max} data sets for the experimental materials on the unfilled resins and the 40% and 60% filled resin composites using two-way analysis of variance (ANOVA) with resin type (5 levels). and curing protocol (5 levels) as the independent variables. Additionally, one-way ANOVA and post-hoc Tukey tests (p=0.05) were performed on DC^{Max} , R_p^{Max} and time to R_p^{Max} data sets versus the cure protocol to confirm the validity or otherwise of the exposure reciprocity law.

2.2.4 Results

Figure 2.1.2 may be used to compare the emission spectra of the five irradiation modes used on the Swiss Master curing light against the absorption spectra of camphoroquinone. From this it is evident that high irradiance results in a greater spectral overlap with camphoroquinone.

Two way ANOVAs for DC^{Max} , R_p^{Max} and time to R_p^{Max} revealed significant differences ($P < 0.001$) in all cases for each of the two independent variables “resin type” and “cure protocol” as well as their interactions. In the case of the unfilled resin series irradiation protocol was more influential ($df = 4$; $F = 9328$) for DC^{Max} than resin type ($df = 4$; $F = 3455$) whereas the opposite occurred with the 40% and 60% filled counterparts (Figure 2.2.1).

(a)



(b)

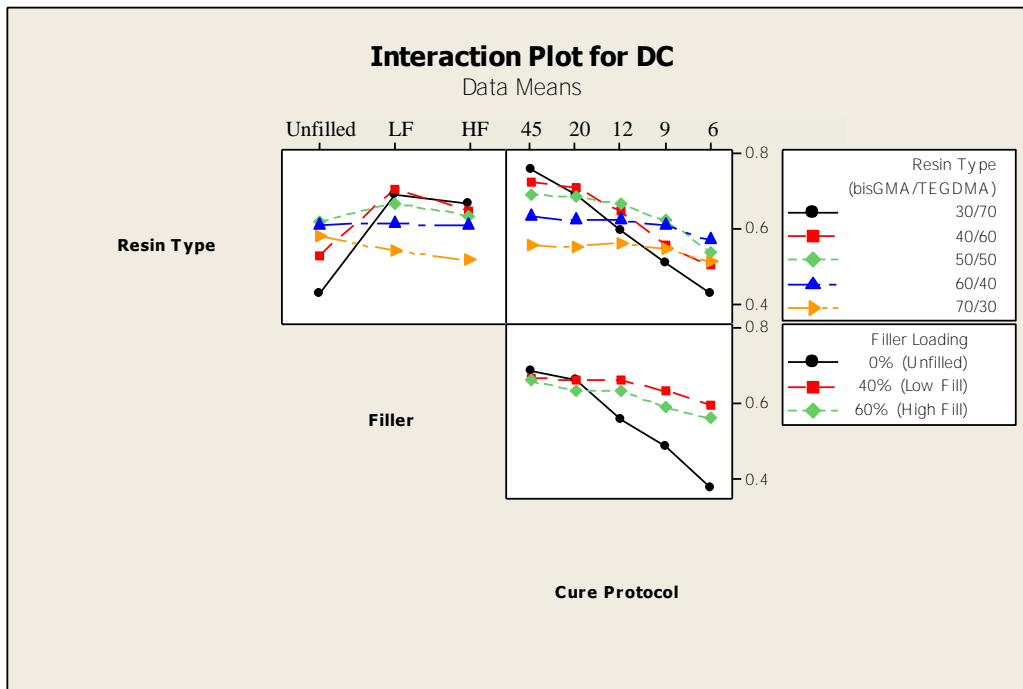


Figure 2.2.1: The main effects plots for resin type, filler percentage and cure protocol for DC^{Max} (a). The interactions of resin type, filler and cure protocol for DC^{Max} (b).

Supplementary one-way ANOVAs also revealed significant differences ($P < 0.05$) between curing protocols (Table 2.2.1). For materials which did not obey the exposure reciprocity law, DC^{Max} was significantly greater for the low irradiance/longest irradiation time (45s at 400mW/cm^2) cure protocol compared to the high irradiance/shortest irradiation time (6s at 3000mW/cm^2) cure protocol ($P < 0.05$). For composite materials, only the high viscosity resins ($>59\%$ Bis-GMA) completely obeyed the exposure reciprocity law and showed no significant differences between cure protocols ($P > 0.05$) (Table 2.2.1). R_p^{Max} was significantly greater and time to R_p^{Max} was significantly lower for the high irradiance/short irradiation time (6s at 3000mW/cm^2) cure protocol ($P < 0.05$) compared to the other protocols for all materials regardless of whether the exposure reciprocity law was obeyed.

Resin	CURE PROTOCOL		DC ^{MAX} (%)			R _p ^{MAX} (%/s)			Time to R _p ^{MAX} (s)		
	Time (s)	Irradiance (mW/cm ²)	RESIN (Unfilled)	40 Vol% (Low Fill)	60 Vol% filler (High Fill)	RESIN (Unfilled)	40 Vol% filler (Low Fill)	60 Vol% filler (High Fill)	RESIN (Unfilled)	40 Vol% filler (Low Fill)	60 Vol% filler (High Fill)
			Mean (+/- SD)	Mean (+/- SD)	Mean (+/- SD)	Mean (+/- SD)	Mean (+/- SD)	Mean (+/- SD)	Mean (+/- SD)	Mean (+/- SD)	Mean (+/- SD)
30/70	45	400	73 (0.2) a	76 (1.9) a	77 (2.5) a	1.87 (0.04) c	3.00 (0.17) c	3.70 (0.14) b	31.7 (0.5) a	19.3 (0.5) a	12.8 (1.7) a
	20	900	62 (0.8) b	74 (1.2) a	71 (1.7) a	2.52 (0.04) b	4.35 (0.25) bc	3.99 (0.31) b	19.8 (0.5) b	13.0 (1.0) b	13.6 (0.0) a
	12	1500	34 (0.2) c	72 (2.2) a	71 (1.8) a	2.51 (0.01) b	5.60 (0.26) b	6.57 (0.44) a	11.1 (0.0) c	9.6 (0.5) c	6.8 (0.9) b
	9	2000	26 (0.2) d	65 (0.6) b	60 (2.3) b	2.60 (0.01) a	6.02 (0.17) b	5.67 (0.47) a	8.2 (0.5) d	7.9 (0.5) d	5.8 (0.6) b
	6	3000	18 (0.4) e	56 (1.9) c	53 (1.8) c	2.65 (0.04) a	9.32 (1.46) a	6.17 (0.07) a	5.1 (0.0) e	5.1 (0.0) e	5.1 (0.0) b
40/60	45	400	72 (0.5) a	72 (0.3) a	72 (2.3) a	2.14 (0.01) c	3.14 (0.16) c	2.99 (0.31) d	26.1 (0.5) a	15.2 (1.5) a	11.9 (0.8) a
	20	900	71 (0.4) a	72 (0.8) a	69 (2.4) b	3.43 (0.02) b	4.46 (0.16) bc	4.00 (0.20) c	17.6 (0.5) b	11.6 (1.3) b	9.9 (0.5) b
	12	1500	53 (0.8) b	72 (1.6) a	69 (2.5) b	3.94 (0.16) a	6.02 (0.20) b	5.68 (0.42) b	11.6 (0.5) c	9.1 (0.5) c	8.8 (0.5) bc
	9	2000	40 (0.3) c	68 (0.9) b	59 (0.9) c	3.87 (0.01) a	6.89 (0.36) b	5.17 (0.14) b	8.5 (0.0) d	7.7 (0.0) d	7.7 (0.0) c
	6	3000	27 (0.2) d	68 (0.4) b	56 (0.3) d	3.97 (0.04) a	12.17 (2.10) a	9.58 (1.11) a	5.1 (0.0) e	3.1 (1.0) e	3.1 (1.0) d
50/50	45	400	71 (0.1) a	68 (0.6) a	68 (1.2) a	2.61 (0.02) e	3.32 (0.17) d	3.18 (0.08) d	19.3 (0.5) a	13.3 (1.0) a	11.6 (1.3) a
	20	900	72 (0.1) a	69 (0.6) a	64 (0.6) a,b	4.16 (0.04) d	5.49 (0.40) c	4.22 (0.15) cd	12.8 (0.0) b	7.1 (0.5) b	9.1 (0.5) b
	12	1500	67 (0.4) b	68 (1.4) a	64 (1.8) a,b	5.24 (0.05) c	7.00 (0.47) b	5.77 (0.72) bc	10.5 (0.5) c	6.0 (0.8) bc	7.7 (0.8) b
	9	2000	59 (0.4) c	67 (0.9) a	61 (2.3) b	5.86 (0.04) b	7.97 (0.21) ab	7.43 (0.59) b	8.2 (0.5) d	6.0 (0.0) bc	3.7 (0.5) c
	6	3000	40 (0.3) d	61 (2.2) b	61 (2.4) b	6.03 (0.05) a	8.73 (0.73) a	10.86 (1.07) a	5.1 (0.0) e	4.8 (0.5) c	2.0 (0.5) c
60/40	45	400	66 (0.5) a	62 (0.6) a	61 (1.7) a	2.44 (0.03) e	3.21 (0.09) e	3.47 (0.11) e	16.7 (0.5) a	10.5 (0.5) a	10.2 (0.9) a
	20	900	65 (0.6) a	61 (1.8) a	60 (1.6) a	3.85 (0.02) d	4.71 (0.30) d	5.62 (0.33) cd	11.1 (0.0) b	7.4 (0.5) b	4.5 (0.5) b
	12	1500	64 (0.6) b	62 (1.4) a	61 (2.1) a	4.89 (0.21) c	6.80 (0.55) c	6.79 (0.38) c	9.1 (0.1) c	5.4 (0.5) c	4.3 (0.8) b
	9	2000	60 (0.5) c	61 (0.4) a	61 (0.3) a	5.74 (0.03) b	8.13 (0.18) b	8.23 (0.23) b	7.7 (0.0) d	3.7 (0.7) d	3.4 (0.0) b
	6	3000	50 (0.4) d	61 (0.2) b	60 (0.9) a	6.40 (0.07) a	10.62 (0.38) a	9.98 (1.49) a	5.1 (0.0) e	3.0 (0.5) d	3.4 (0.0) b
70/30	45	400	61 (0.5) a	54 (1.1) a	52 (2.0) a	2.34 (0.08) e	2.89 (0.26) e	2.39 (0.07) e	14.46 (0.9) a	5.4 (0.5) a	6.2 (0.5) a
	20	900	60 (1.2) a	54 (0.3) a	52 (1.6) a	3.65 (0.12) d	4.29 (0.28) d	4.32 (0.38) d	8.5 (0.9) b	5.1 (0.9) a	4.8 (0.5) ab
	12	1500	60 (0.7) a	55 (0.8) a	52 (1.5) a	4.97 (0.10) c	5.76 (0.38) c	5.75 (0.38) c	7.1 (0.5) c	4.8 (0.5) a	5.1 (0.8) ab
	9	2000	57 (0.2) b	54 (0.3) a	52 (1.0) a	5.72 (0.06) b	7.05 (0.04) b	7.16 (0.62) b	6.2 (0.5) d	2.8 (0.5) b	3.4 (0.9) bc
	6	3000	52 (0.8) c	52 (0.5) a	50 (1.5) a	7.40 (0.81) a	8.21 (0.06) a	8.98 (0.81) a	5.1 (0.0) e	2.3 (0.5) b	2.3 (0.5) c

Table 2.2.1: DC^{Max}, R_p^{Max} and Time to R_p^{Max} of experimental formulations. Similar letters within columns indicate no significant differences (p<0.05).

2.2.5 Discussion

The first hypothesis (that for similar radiant exposure, if reciprocity did not hold for the parent resin then it should not hold for its composite counterpart) must be rejected as reciprocity was not verified in any of the tested resins but was for some composites. However, the differences in DC^{Max} between the extremes of the irradiation modes for the resins became less marked for high viscosity resins (9% for 7:3 compared with 55% for 3:7 Bis-GMA/TEGDMA ratio). Previous work has indicated that reciprocity can be verified above a specific (unfilled) resin viscosity ($\geq 6:4$, Bis-GMA/TEGDMA) (Feng and Suh, 2007) which plays a significant role in the polymerization rate and the onset of reaction-diffusion controlled termination, with the highest maximum rates observed for mixtures containing 50-75 wt% Bis-GMA (Emami and Söderholm, 2003). As suggested in the previous section, low viscosity resins have high system mobility and thus a bimolecular termination (radical-radical) pathway dominates during the initial stages of polymerization. Therefore, loss of radical growth centres is greater for high irradiance protocols (where a multitude of radical growth centers co-exist), since at low conversion termination is proportional to the squared concentration of free radicals (Leprince et al., 2010b). Consequently, at constant radiant exposure, with higher irradiance (and shorter time), a reduction in growth centres during auto-acceleration is expected and DC^{Max} will be significantly decreased compared with similar radiant exposures achieved using lower irradiance for longer time. Furthermore, the initiation rate (R_p) has a square root dependency on intensity (Decker, 2002). This relationship only holds in the early stages of polymerization where termination occurs by bimolecular reactions between polymer radicals (i.e. in low viscosity resins). As polymerization proceeds, mobility restrictions increase as the density of the polymer increases and sole segmental diffusion of cross-linked polymer radicals become less likely and the radical site will mainly move by reacting with neighboring functional groups until termination occurs. A more linear dependence of R_p on light intensity is observed after gelation as the polymer chain stops growing and corresponds to a monomolecular termination process (Decker, 2002). For high viscosity resins, polymerization exhibits reaction-diffusion-controlled termination

from the onset due to its very low mobility (Lovell et al., 2001; Emami and Söderholm, 2003) and thus reaction rate increases significantly with light intensity (Table 2.2.1). The significant increase in the rates at high intensity allows for a reduced cure time (reduced time to R_p^{Max}) and thus the differences in DC^{Max} between the extremes of the curing protocols are lower.

The second hypothesis (an increase in the filler content will result in a decrease in the degree of conversion regardless of curing mode) must also be rejected. Indeed, the addition of filler proportionally decreases the concentration of carbon double bonds available for conversion and thus conversion is expected to be lower. When compared to highly filled systems, vitrification is expected to be delayed in unfilled or low filled systems which will allow further radical propagation and improved conversion in such systems (Feng and Suh, 2007; D'alphino et al., 2007; Feng et al., 2009). This may well explain why resins and composites made with ($\leq 6:4$, Bis-GMA/TEGDMA) always show the highest DC^{Max} for the reference protocol (45s at $400\text{mW}/\text{cm}^2$). However, the increased viscosity due to the addition of filler served to decrease the differences in DC^{Max} between the extremes of the irradiation modes for the filled systems (7:3 Bis-GMA/TEGDMA ratio: LF= 2%, HF= 2%; 3:7 Bis-GMA/TEGDMA ratio: LF= 20%, HF= 24%). Filled systems are expected to approach a more reciprocal relationship with respect to irradiance and time at constant radiant exposure by reducing the mobility of the propagating species. However, the applicability of the exposure reciprocity law may not be related to the bulk or macroscopic viscosity (Section 2.1). In filled systems, the diffusion of radicals will be limited to a relatively small volume as the organic fractions will be located in the small interstices between the inorganic particles and thus a monomolecular termination pathway would be favoured due to the reduced mobility (i.e. increased local viscosity). The local viscosity will increase as the contact area between the resin and the filler increases (due to interlocking and interfacial interactions between filler and resin), i.e. increasing the filler load, using smaller particle sizes or a combination of both (Leprince et al., 2009; Beun et al., 2009).

The kinetics of polymerisation is complex and changes in viscosity in the early stages differ with material composition. The effect of changing rheology throughout cure is an essential consideration, which is needed to further understand the applicability of

the exposure reciprocity law. Additionally, the optical properties are also an important consideration. Most notably the photoinitiator type may have a profound impact on reciprocity (Section 2.3) due to greater efficiencies caused by greater molar absorptivity (molar extinction coefficient). Since the decomposition of photoinitiators is directly related to its absorbance optical properties are likely to be affected throughout cure and will subsequently be investigated in Chapter 3.

Finally, temperature measurements should not be ignored as both irradiance and time will influence the total heat energy delivered and consequently will affect the kinetics and reaction exotherm (Section 2.1). Further work regarding the synergistic effects of rheology, optical properties and temperature should be a subject of further investigation.

2.2.6 Conclusion

Exposure reciprocity relies on the constituents of the resin matrix as well as the filler percentage. From this work it is evident that local viscosity (i.e. that between the fillers) is responsible for whether exposure reciprocity law is upheld. When a bi-radical termination pathway dominates, reciprocity is unlikely to be verified when the material contains only a camphoroquinone/amine photoinitiator system. However, when a monomolecular termination pathway dominates, degree of polymer conversion is more likely to be equal regardless of curing protocols with similar radiant exposure. This further underlines the necessity to adapt curing regimes according to material filler content, resin matrix composition and the photoinitiator system. This is of prime clinical relevance as the development of composite materials may be optimized to ensure equivalent polymer conversion at reduced curing times.

2.3 Photoinitiator Type and Applicability of the Exposure Reciprocity Law in Filled and Unfilled Photoactive Resins^{†‡}

2.3.1 Abstract

The aim of this section is to test the influence of photoinitiator type and filler particle inclusion on the validity of exposure reciprocity law. 50/50wt% Bis-GMA/TEGDMA resins were prepared with equimolar concentrations of camphorquinone/DMAEMA (0.20/0.80 mass%) (CQ) or Lucirin-TPO (0.42 mass%), and were used either unfilled or filled to 75 mass%. Specimens were cured with a halogen Swiss Master Light (EMS, Switzerland) using four different curing protocols: 400 mW/cm² for 45 s as reference protocol (18 J/cm²), 1500 mW/cm² for 12 s (18 J/cm²), 3000 mW/cm² for 6 s (18 J/cm²) and 3 s (9 J/cm²). Degree of conversion (DC) was measured in real time for 70 s by FT-NIRS and temperature rise using a thermocouple. Depth of cure was determined with a penetrometer technique. With respect to DC and depth of cure, exposure reciprocity law did not hold for any tested material, except for the depth of cure of filled CQ-based materials. At similar radiant exposure, DC was significantly higher ($p < 0.05$) for all unfilled and filled TPO-based materials compared with CQ-based materials. As exposure time was reduced and irradiance increased, TPO-based materials exhibited higher DC whilst an opposite trend was observed for CQ-based materials ($p < 0.05$). For similar curing regimes, depth of cure of CQ-based materials remained significantly greater than that of TPO-based materials. Adding fillers generally reduced DC, except at higher irradiance for CQ-based materials where a positive effect was observed ($p < 0.05$). The validity of exposure reciprocity law was dependent on several factors, among which photoinitiator type and filler content were important. Lucirin-TPO is a highly reactive and efficient photoinitiator, which may allow the potential for a reduction in curing time of TPO-based photoactive materials in thin sections due to its reduced depth of cure.

[†] Leprince J, Hadis MA, Ferracane JF, Devaux J, Leloup G, Shortall ACC, Palin WM. Photoinitiator type and applicability of exposure reciprocity law in filled and unfilled photoactive resins. *Dental Materials*, 2011; 27: 157-164

[‡] Winner of the Paffenbarger Award, Academy of Dental Materials, Portland, 2009.

2.3.2 Introduction

The demand for the reduction in curing time and minimized chairside procedure is un-questionable. Combined with the increasing popularity and demand for photoactive resin-based restorations as a replacement for dental amalgam, major financial implications have been suggested for dentists who chose to use high power, fast curing light sources (Christensen, 2000). Such light curing units may allow the possibility of shorter exposure times (<10 s) whilst maintaining optimum material properties. This is based on the assumption that the total energy of the irradiation, i.e. radiant exposure (J/cm^2), is the main determining factor in degree of conversion and mechanical properties of the photoactive material. As radiant exposure is the product of irradiance (I , mW/cm^2) and irradiation time (t , s), it supposed that comparable material properties result from similar radiant exposure, no matter how it is obtained (by different combinations of irradiance and time). While in certain cases this law is upheld (Sakaguchi and Ferracane, 2001; Halvorson et al., 2002; Emami and Söderholm, 2003; Price et al., 2004), other studies have reported that although radiant exposure plays an important role, irradiance and time independently influence polymer chain length, extent of crosslinking and mechanical properties, and as such, exposure reciprocity does not hold under all conditions (Musanje and Darvell, 2003; Asmussen and Peutzfeldt, 2005; Peutzfeldt and Asmussen, 2005; Feng and Suh, 2007; Dewaele et al., 2009).

Sections 2.1 and 2.2 of this thesis demonstrated the importance of monomer composition and filler percentage on reactivity and degree of conversion of RBC materials. Reactivity and degree of conversion are also dependent upon the efficiency of the photoinitiator system (Ogunyinka et al., 2007). Furthermore, alternative photoinitiators have been integrated in commercially available materials in the last few years, mostly for aesthetic reasons, due to yellowing by camphorquinone/amine systems (Price and Felix, 2009). As some of these alternative photoinitiators are also interesting in terms of photopolymerization efficiency, they might also have an impact on the validity of exposure reciprocity law. Notably, Lucirin-TPO, a monoacylphosphine oxide, seems to be a promising molecule, as it presents higher molar absorptivity and curing efficiency (Neumann et al., 2005; Neumann et al., 2006). However, considering the absorption

characteristics of Lucirin TPO and the increased scattering of light of shorter wavelengths, inferior curing depth of materials that replace camphoroquinone with Lucirin TPO may be a cause for concern (Price and Felix, 2009; Leprince et al., 2010a). Although recent work has highlighted the potential of Lucirin TPO-based (unfilled) polymers to be cured in thick layers (Kenning et al., 2008) and for use in dental adhesives (Illie and Hickel, 2009), its application in filled resin composites has not been fully described. To date, the curing efficiency of Lucirin TPO in filled resin-based composite has only been studied in commercial materials (Palin et al., 2008; Price and Felix, 2009; Leprince et al., 2010a), in which Lucirin TPO is not used by itself but in combination with camphoroquinone. Lower quantum yields have been measured for polymerization initiated by a mixture of Lucirin TPO and camphoroquinone compared to systems in which Lucirin TPO was used alone, probably due to a transfer from the more efficient photoinitiator (Lucirin TPO) to the less efficient one (camphoroquinone) (Neumann et al., 2009). In addition, the exact photoinitiator concentration and ratio in commercially available composite materials usually remains proprietary.

Consequently, the aim of this work was to investigate the impact of photoinitiator type on the polymerization efficiency in unfilled and filled experimental resins. Degree of conversion (*DC*) and depth of cure were measured to investigate the applicability of exposure reciprocity law on the photoinitiator type and also the presence or absence of fillers.

2.3.3 Materials and Methods

50/50 mass% Bis-GMA/TEGDMA resins were prepared with two different photoinitiator systems at equimolar concentration ($0.0134 \text{ mol.dm}^{-3}$): either camphorquinone, and the co-initiator, dimethylaminoethyl methacrylate (0.20/0.80 mass%) or Lucirin-TPO. The more bulky Lucirin TPO structure has a molecular weight of almost twice that of camphorquinone (348 and 166, respectively; Figure 2.3.1) and an equimolar concentration of TPO corresponded to 0.42 mass%.



Figure 2.3.1: The chemical structures of the photoinitators tested in the present study.

The resins were tested either unfilled or filled with silanated barium glass fillers ($0.7 \mu\text{m}$; Esschem Europe LTD, Seaham) and fumed silica (14 nm ; Aerosil 150, Evonik Industries, Germany) to 65/10 mass%, respectively. Specimens were cured with a 11 mm diameter tip halogen Swiss Master Light (EMS, Switzerland), chosen for its broad spectrum, overlapping the absorption spectrum of both photoinitiators (albeit only partially with Lucirin TPO). Emission spectra were determined using a method used previously (Chapter 2.1 and 2.2) The absorption spectra of the photoinitiators were determined in methyl methacrylate (Sigma-Aldrich, UK) using the UV-Vis spectrometer coupled to a standard cuvette holder in absorbance mode. Three different curing protocols were applied to compare the applicability of exposure reciprocity law and another one to assess the potential of time reduction at the highest irradiance. 400 mW/cm^2 for 45 s was chosen as the reference protocol and compared to protocols of the same radiant exposure (18 J/cm^2): 1500 mW/cm^2 for 12 s and 3000 mW/cm^2 for 6 s. The last irradiation protocol, 3000 mW/cm^2 for 3 s, was chosen to assess the effect of half the radiant exposure on Lucirin TPO-based materials. The irradiance values were based on

the in-built radiometer of the curing-unit and verified by the spectrometer. Degree of conversion (DC) and rate of polymerization (R_p) were measured in real-time by Fourier Transform near infrared spectroscopy (FT-NIRS; Figure 2.1.1) by monitoring the height of the peak at 6164 cm^{-1} . The method was previously described in Chapter 2.1 and 2.2. Though DC is known to slightly evolve up to 24 h after irradiation, the major part of the polymerization occurs during irradiation (Truffier-Boutry et al., 2006). For that reason, DC at 70 s after the start of irradiation was chosen for comparison and will be referred to as DC^{Max} . A real time measurement of the maximum temperature rise (ΔT^{Max}) during the polymerization was performed using a similar set up as for FT-IR measurements, except that the mould was split on one side to insert a thermocouple. The latter was glued in place so that the bead of the thermocouple (1.5 mm diameter) was in contact with the outer edge of the sample. The thermocouple was calibrated against a standard mercury thermometer. The depth of cure was measured by an penetrometer technique (Shortall et al., 1995) after curing the samples in a cylindrical Teflon mould (5 mm diameter) (n=3).

Statistical analyses were performed on the DC^{Max} , R_p^{Max} , depth of cure and ΔT^{Max} data sets using a general linear model (GLM) three-way analysis of variance (ANOVA) with photoinitiator (2 levels), filler inclusion (2 levels) and curing mode (4 levels) as the independent variables. Supplementary two-way ANOVA and post-hoc Tukey multiple comparison tests (p=0.05) were also performed.

2.3.4 Results

Figure 2.3.2 presents a comparison between the absorption spectra of both photoinitiators (Camphorquinone and Lucirin TPO) and the emission spectra of the curing-unit at 400, 1500 and 3000 mW/cm². An increased power of the curing unit resulted in a greater effective irradiance at shorter wavelengths and thus the overlap with Lucirin TPO absorption was improved. However, spectral emission was better adapted to camphorquinone absorption, as a greater overlap of the two spectra was evident.

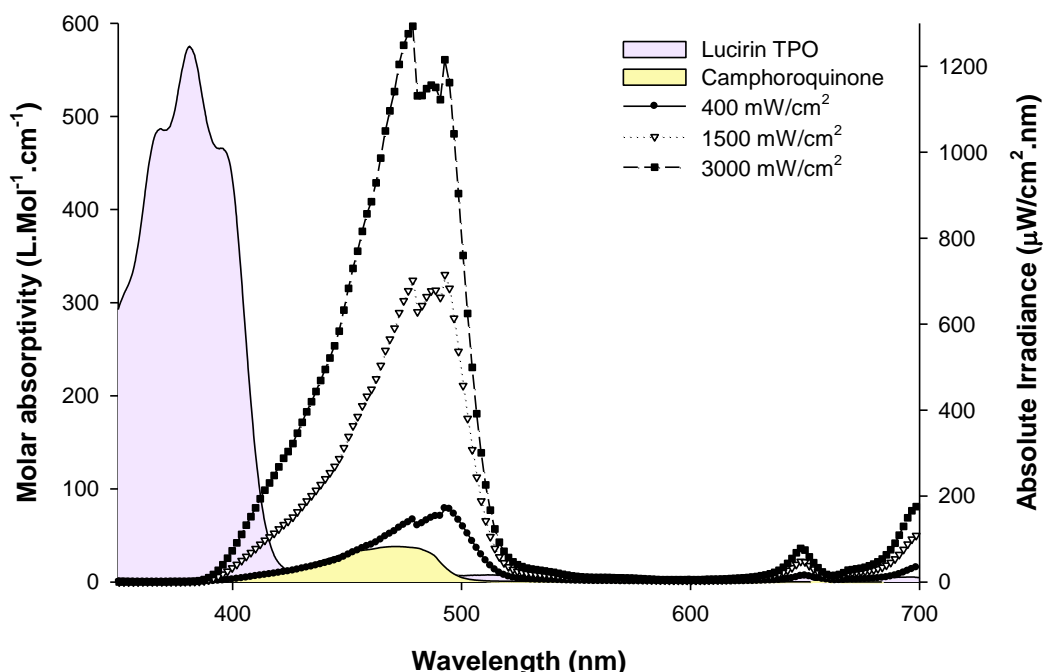
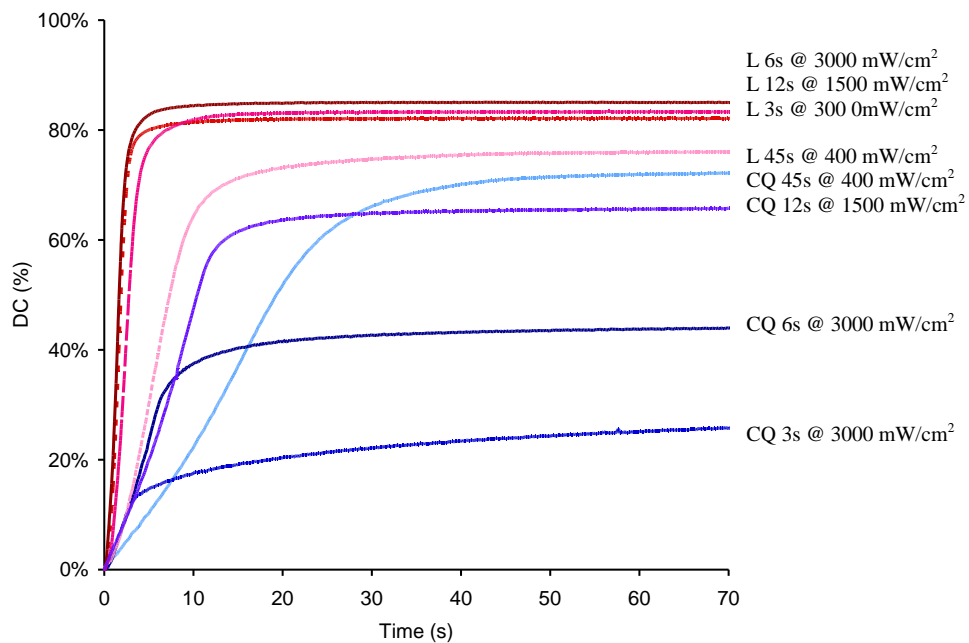


Figure 2.3.2: Absorption spectra (left axis, in $\text{L.Mol}^{-1}\text{cm}^{-1}$) of Camphoroquinone (yellow) and Lucirin TPO (purple) compared to the emission spectra of the Swiss Master Light (right axis, in $\mu\text{W}/\text{cm}^2.\text{nm}$) at different irradiances, i.e. 400, 1500 and 3000 mW/cm^2 . An Improvement of the overlap with Lucirin TPO absorption can be observed with increased power of the curing unit.

The GLM of all independent variables revealed significant differences ($p < 0.001$) where photoinitiator ($\text{df}=1$; $F=5971$) > filler inclusion ($\text{df}=1$; $F=1302$) > curing mode ($\text{df}=3$; $F=809$). A two-way ANOVA test of DC^{Max} results revealed significant overall differences for both photoinitiators and curing mode ($p < 0.001$). Comparing irradiation modes of similar radiant exposure ($18 \text{ J}/\text{cm}^2$), DC^{Max} was significantly higher for all resin and resin composite Lucirin TPO-based materials compared with Camphoroquinone-based materials ($p < 0.05$; Figure 2.3.3(a) and Figure 2.3.4(a)). For Lucirin TPO, an increase in irradiance resulted in higher DC^{Max} regardless of exposure duration, while a reverse trend was observed for camphoroquinone-based materials where longer exposure led to higher DC^{Max} regardless of irradiance. DC^{Max} significantly decreased in most cases with addition of fillers ($p < 0.05$), except for camphoroquinone-based materials cured with the highest irradiance where filled resins exhibited a significantly increased DC^{Max} compared with the unfilled parent resin ($p < 0.05$).

(a)



(b)

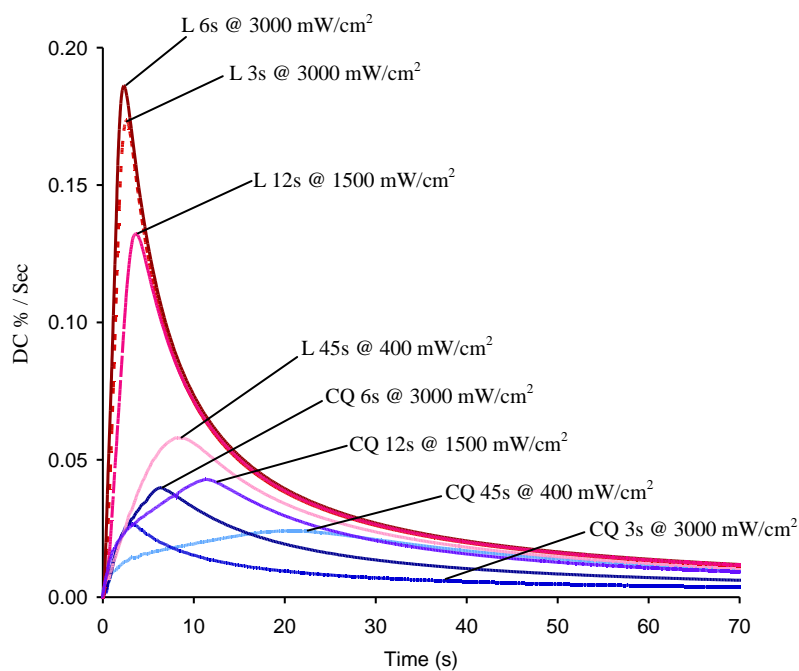
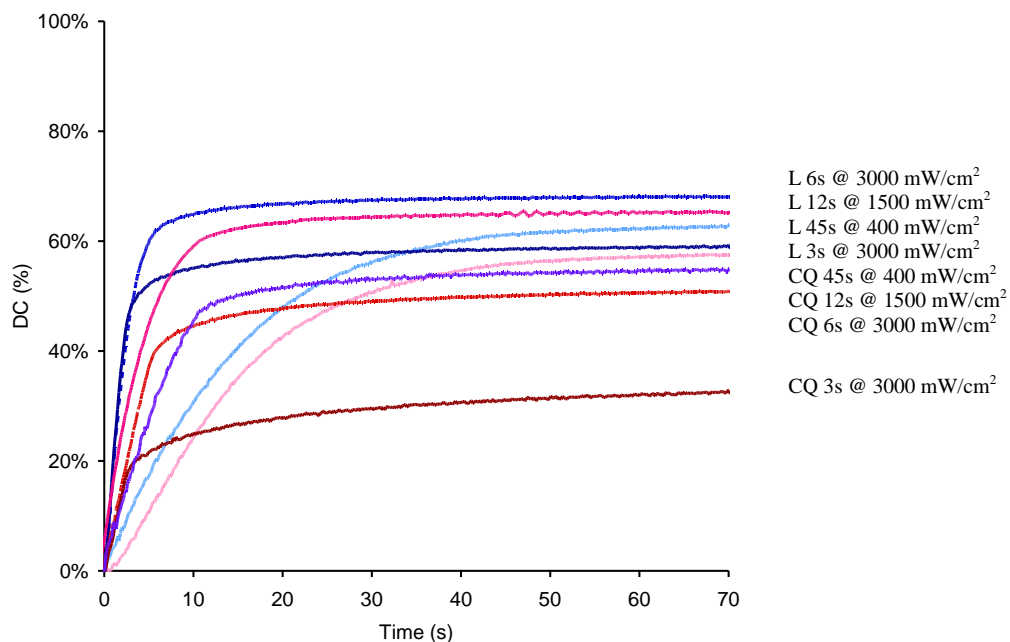


Figure 2.3.3: DC (a) and R_p (b) curves in real time during 70 s for 50/50 mass% Bis-GMA/TEGDMA resin. Blue and red curves correspond respectively to Lucirin TPO-based and CQ-based resins. Irradiation modes are indicated on the charts.

(a)



(b)

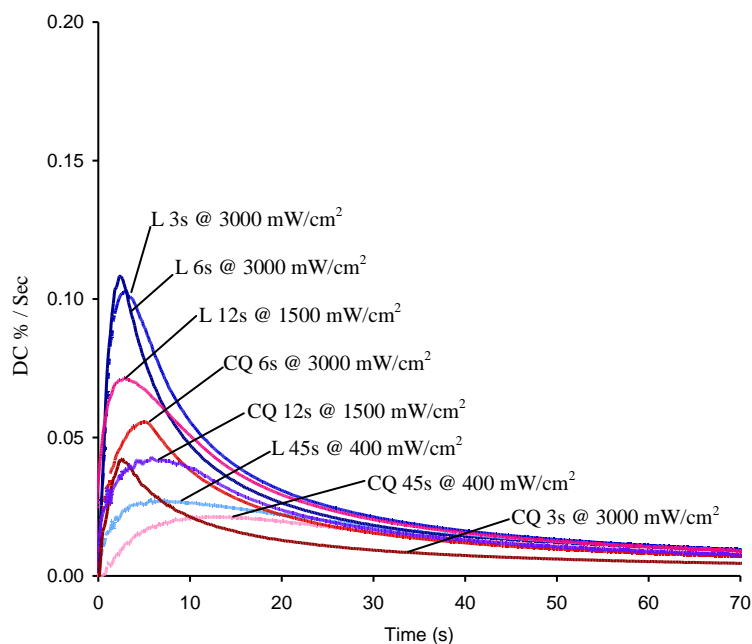


Figure 2.3.4: DC (a) and R_p (b) curves in real time during 70 s for 50/50 mass% Bis-GMA/TEGDMA resin filled with 75 mass% fillers. Blue and red curves correspond respectively to Lucirin TPO-based and CQ-based resins. Irradiation modes are indicated on the charts.

For each curing regime, all Lucirin TPO-based materials exhibited a significantly increased R_p^{Max} ($p < 0.05$) compared with camphoroquinone-based materials, except for filled Lucirin TPO and camphoroquinone-based composites irradiated with the reference curing mode (400 mW/cm^2 for 45 s) ($p > 0.05$) (Figure 2.3.3(b) and Figure 2.3.4(b); Table 2.3.1 and Table 2.3.2). Regarding ΔT^{Max} (Table 2.3.1 and Table 2.3.2), values were significantly higher for Lucirin TPO-based unfilled resins compared with camphoroquinone-based unfilled resins, except for the reference curing protocol. A comparable trend was observed for filled materials, although the differences between similar curing modes were not statistically significant ($p > 0.05$). Moreover, multiple regression analysis revealed that DC^{Max} and ΔT^{Max} were highly correlated for unfilled ($R^2 = 0.95$; $p < 0.0001$) and to a lesser extent for filled resins ($R^2 = 0.74$; $p < 0.0001$). Contrary to the observations of DC^{Max} and R_p^{Max} , depth of cure was significantly higher for camphoroquinone-based compared with Lucirin TPO-based materials (~1.5 to 3 times for unfilled and filled materials, respectively; Table 2.3.1 and Table 2.3.2).

Comparing modes of different radiant exposure, both DC^{Max} and depth of cure were significantly lower ($p < 0.05$) for camphoroquinone-based materials using the lower radiant exposure (9 J/cm^2). Interestingly, in the case of Lucirin TPO-based materials, DC^{Max} at low radiant exposure is comparable in some cases or at least close to DC^{Max} at high radiant exposure, especially for unfilled resins. However, depth of cure at 9 J/cm^2 was significantly reduced ($p < 0.05$) compared to all specimens cured using 18 J/cm^2 (Table 2.3.1 and Table 2.3.2).

Material	Cure Protocol		Radiant Exposure (J/cm ²)	DC^{Max} (%)	R_p^{Max} (s ⁻¹)	T^{Max} (°C)	DOC (mm)
	Time (s)	Irradiance (mW/cm ²)		Mean (+/- SD)	Mean (+/- SD)	Mean (+/- SD)	Mean (+/- SD)
Unfilled Camphoroquinone	45	400	18	72.1 (0.7) d	0.024 (0.000) e	33.6 (2.0) b,c	18.2 (0.2) a
	12	1500	18	65.8 (0.5) e	0.043 (0.001) c,d	29.1 (1.4) c	16.9 (0.1) b
	6	3000	18	43.9 (1.2) f	0.040 (0.001) d	19.5 (3.9) d	14.4 (0.1) c
	3	3000	9	25.7 (1.3) g	0.028 (0.001) d,e	10.6 (1.4) e	9.5 (0.1) f
Unfilled Lucirin TPO	45	400	18	76.0 (0.1) c	0.058 (0.001) c	29.5 (1.9) b,c	9.6 (0.3) f
	12	1500	18	83.3 (0.5) a,b	0.133 (0.005) b	41.5 (3.2) a	11.8 (0.2) d
	6	3000	18	85.0 (0.3) a	0.186 (0.006) a	41.1 (1.8) a	10.4 (0.2) e
	3	3000	9	82.1 (0.2) b	0.177 (0.014) a	36.4 (3.3) a,b	4.4 (0.2) g

Table 2.3.1: DC^{Max} , R_p^{Max} , T^{Max} and depth of cure (DOC) of unfilled resins. Similar letters within columns indicate no significant differences (p<0.05).

Material	Mode		Radiant Exposure (J/cm ²)	DC^{Max} (%)	R_p^{Max} (s ⁻¹)	T^{Max} (°C)	DOC (mm)
	Time (s)	Irradiance (mW/cm ²)		Mean (+/- SD)	Mean (+/- SD)	Mean (+/- SD)	Mean (+/- SD)
Filled Camphoroquinone	45	400	18	57.6 (0.7) c,d	0.022 (<0.001) e	14.0 (0.3) c	10.6 (0.2) a
	12	1500	18	54.9 (0.9) d	0.043 (0.001) d	17.6 (0.5) a,b,c	10.8 (0.2) a
	6	3000	18	50.8 (0.2) e	0.056 (0.001) c	17.0 (1.4) a,b,c	10.1 (0.3) a
	3	3000	9	32.5 (1.0) f	0.043 (0.001) d	7.6 (0.1) d	7.3 (0.2) b
Filled Lucirin TPO	45	400	18	62.7 (0.5) b	0.028 (0.001) e	14.5 (0.4) b,c	3.2 (0.2) d
	12	1500	18	65.2 (2.5) a,b	0.071 (0.01) b	18.9 (1.0) a,b	3.7 (0.1) c,d
	6	3000	18	68.0 (0.3) a	0.104 (0.005) a	19.9 (0.2) a	3.9 (0.3) c
	3	3000	9	59.0 (1.7) c	0.109 (0.002) a	13.6 (4.2) c	2.5 (0.3) e

Table 2.3.2: DC^{Max} , R_p^{Max} , T^{Max} and depth of cure (DOC) of filled resins. Similar letters within columns indicate no significant differences (p<0.05).

2.3.5 Discussion

Both photoinitiator type and filler content significantly influence material properties cured with similar radiant exposure by varying irradiance and/or exposure time, which contests the applicability of the exposure reciprocity law. However, in certain instances comparable DC^{Max} and depth of cure were observed regardless of curing protocol. Whilst reciprocity is observed for depth of cure of the camphoroquinone-based composite in accordance with Nomoto *et al.* (1994), it is not the case for DC^{Max} of the same material tested here. The trend observed for camphoroquinone-based materials is in agreement with previous literature (Peutzfeldt and Asmussen, 2005; Dewaele *et al.*, 2009) supporting that for these systems, time has more impact than irradiance on the material properties. As explained in recent work on trapped free radicals (Leprince *et al.*, 2010b), there is probably a loss in radical growth centres by translational diffusion and bimolecular (radical-radical) termination during the early stages of polymerization wherein system mobility is high. This loss (or early termination) is greater for high irradiance protocols, since at low conversion termination is proportional to the squared concentration of free radicals. As a result, at a radiant exposure with higher irradiance (and shorter time), the amount of growth centres created during auto-acceleration are reduced and DC^{Max} is significantly decreased compared with similar radiant exposure achieved using lower irradiance for longer time. However, while variations in DC^{Max} of camphoroquinone-based materials are obvious for unfilled systems, the differences appear more subtle in filled materials and approach a reciprocal relationship with respect to irradiance and time. This effect is supported by previous findings (Leprince *et al.*, 2009) demonstrating that trapped free radical concentrations were nearly 3 times higher in the organic matrix of a filled composite than in a corresponding unfilled resin and was related to spatial confinement of radicals in filled resins, preventing their recombination. It could indeed be assumed that the organic fraction is located in small interstices between the inorganic particles, which could restrict the diffusion of radicals to a relatively small volume. But, it could also be related to an increase in local viscosity as reported in the previous sections (Section 2.1 and 2.2), reducing molecular mobility, thereby favoring a monomolecular termination pathway. Feng and Suh (2007)

demonstrated that above a specific Bis-GMA/TEGDMA ratios ($\geq 6:4$, Bis-GMA/TEGDMA) reciprocity holds (Feng and Suh., 2007). Interestingly, Beun et al. (2009) suggested that the large contact area between fillers and resin probably results in a significant local increase in viscosity. Accordingly in section 2.1 and 2.2 it was shown that for low viscosity resins, the addition of fillers significantly improved conversion and it was suggested that this was due to a significant increase in local viscosity.

Consequently, at the highest irradiance, the filled camphoroquinone-based materials exhibited a significant increase in DC^{Max} and R_p^{Max} compared with the unfilled parent resin (50.8 and 43.9 %, 0.056 and 0.040 s⁻¹, respectively). Since the polymerization kinetics of camphoroquinone-based materials at high irradiance are limited, inclusion of filler particulates may have also affected light scattering properties through the limited specimen thickness (1.4mm) and/or provided some thermal insulating effect which improved final conversion of the filled composite.

By comparing camphoroquinone- with Lucirin TPO-based materials, deviations from the simple assumption of exposure reciprocity law can be highlighted. The significant increase in DC^{Max} of Lucirin TPO- compared with camphoroquinone-based materials (Figure 2.3.3(a) and Figure 2.3.4(a); Table 2.3.1 and Table 2.3.2) can be explained by an increased generation of free radicals resulting from the vastly higher molar absorptivity (Figure 2.3.2) and quantum yield efficiency of Lucirin TPO (Neumann et al., 2005; Neumann et al., 2006) compared with camphoroquinone. Photochemical cleavage of Lucirin TPO (a Norrish Type I photoinitiator) results in two free radical intermediates compared with one active radical in camphoroquinone/amine systems, and the absence of co-initiator for Lucirin TPO reduces the intermediate steps required for radical production. Nevertheless, even if it can be assumed that an increased number of free radicals led to higher DC^{Max} , the increased efficiency of higher irradiance modes achieved by curing Lucirin TPO-based materials with the halogen curing unit is interesting, especially in unfilled resins. Firstly, the positive impact of higher irradiance on DC^{Max} is difficult to estimate due to a greater effective irradiance and complex overlap of Lucirin TPO-absorption with increasing power of the curing unit (Figure 2.3.2). Secondly, the quantum yield efficiency was reported to be 5 times higher for Lucirin

TPO than for camphoroquinone with a halogen light (Neumann et al., 2006). Hence, despite some radical loss by early recombination, there is probably a much quicker build-up of the polymer network into a gel-phase, favouring radical entrapment and leading to an earlier onset of autoacceleration (Figure 2.3.3(b) and Figure 2.3.4(b)). During autoacceleration, which is characterized by a low termination rate, Lucirin TPO-based materials generate many more polymer growth centers. At high irradiance, this results in a significantly increased R_p^{Max} compared with camphoroquinone-systems, which can compensate time reduction and lead to high DC^{Max} . Although the results are presented as a function of the total power of the irradiating light source rather than the total number of photons absorbed directly by the photoinitiator, the conclusions remain valid since both quantities remained proportional for each curing protocol. In addition, the influence of temperature should not be ignored. As high R_p^{Max} is associated with a greater and quicker reaction exotherm, it further improved DC^{Max} by increasing ΔT^{Max} , hence monomer mobility and increasing vitrification temperature

Despite their obvious higher curing efficiency highlighted above, Lucirin TPO-based materials exhibit inferior curing efficiency through depth, for unfilled and even more for filled materials. Since Lucirin TPO is a colourless photoinitiator and has been shown to exhibit efficient photo-bleaching properties with no absorbance of photolysis products above 360 nm (Stephenson Kenning et al., 2008), negligible light attenuation through an unfilled Lucirin TPO-based resin might be assumed although high exotherm and increasing UV radiation may cause some oxidation and discolouration of the polymer resin. The more likely explanation for inferior cure depth of Lucirin TPO-based materials is related to the much higher molar absorptivity of Lucirin TPO (Neumann et al., 2005). At similar photoinitiator concentration, Lucirin TPO utilizes many more photons than camphoroquinone due to the much higher molar extinction coefficient (Neumann et al., 2005). Therefore, the penetration of appropriate wavelength light (i.e. <420nm; Figure 2.3.2) needed to activate Lucirin TPO is reduced rapidly through depth. As such, the absorption properties of Lucirin TPO and camphoroquinone are interesting in terms of photoinitiator optical properties of RBC materials, which will be further investigated in Chapter 3.

In the present study, the ratio of camphoroquinone:amine was selected as an optimum combination to produce maximum DC (Yoshida and Greener, 1994) and the camphoroquinone concentration was similar to that used in commercial composites (Taira et al., 1988), and an equimolar concentration was used for Lucirin TPO for comparative purposes. However, given the very high polymerization efficiency of Lucirin TPO, its concentration could probably be reduced, thereby improving depth of cure. As for the additional decrease of depth of cure in filled materials, it is probably related to an increased light scattering at the shorter wavelengths that correspond with the effective absorption of Lucirin TPO (Price and Felix, 2009). However, the effective spectral irradiance of the curing-unit in the absorption range of Lucirin TPO is several times lower than that of camphoroquinone. As there is a correlation between irradiance and depth of cure (Lindberg et al., 2004), a higher depth of cure for Lucirin TPO-based materials might be achieved using a higher intensity source emitting specifically around 400 nm.

2.3.6 Conclusion

In most cases, the implicit assumption of exposure reciprocity law was not upheld. Exposure reciprocity relies on the extent of monomolecular termination of the curing material, which is directly affected by molar absorptivity, quantum yield efficiency, material viscosity and temperature. However, while the validity of this law appears far from systematic, there is a clear interest to consider its global relevance to clinical practice and the potential of significant time reduction using modern high-irradiance curing lights. From this work, it appears that this potential clearly depends on the photoinitiator type, as Lucirin TPO-based materials reach higher DC^{Max} than camphoroquinone-based materials at similar or even half radiant exposure. However, the problem of low depth of cure of Lucirin TPO-based materials still has to be addressed. For camphoroquinone-based systems, filler content has an impact on the efficiency of high irradiance modes, probably due to decreased local mobility of monomers. This underlines the necessity to adapt curing protocols according to the material filler content, which is of prime interest for clinicians.

2.4 Summary

The aim of this chapter was to investigate the reason why in some cases exposure reciprocity is verified whilst in others it is not and to investigate the possibility of reduced curing time based on radiant exposure.

Manufacturers of RBCs typically produce flowable materials either by reducing the filler percentage used in the high viscosity versions or by increasing the amount of diluent monomer used, both of which are likely to affect the kinetics of polymerization. Five commercial resin composite materials (Venus; Heraeus Kulzer, Synergy D6; Coltene, Premis; Kerr, Grandio; Voco and Gradia; GC Corp) and their counterpart flowable versions were tested (Section 2.1). The materials were irradiated at various combinations of irradiance and time so that they received a total radiant exposure of 18 J/cm^2 . The majority of the higher viscosity resin composite paste materials exhibited similar degree of conversion regardless of curing protocol. However, a significant decrease in degree of conversion for specimens cured at 3000 mW/cm^2 for 6 s compared with 400 mW/cm^2 for 45 s was observed for flowable materials, Grandio (41 ± 0.36 and 62 ± 1.15 %, respectively) and Venus (44 ± 0.44 and 67 ± 0.44 % respectively). Conversely, other flowable materials exhibited little or no significant difference between curing modes. The validity of exposure reciprocity and the extent of conversion depend on several factors amongst which resin viscosity and filler content are important factors in determining radical termination pathway. Bi-radical termination pathways are expected to be dominant when system mobility is high and therefore reciprocity is expected to fail. However, as reciprocity was verified for some low viscosity materials, it was suggested that the extent of bi-radical terminations depends upon the local viscosity (i.e. the resins between the fillers) and not the macroscopic viscosity. Consequently, the problem of partly unknown composition limits the analysis of commercial materials. Therefore, to further understand such an effect, exposure reciprocity was investigated in model systems comprising of Bis-GMA/TEGDMA ratios of 70/30, 60/40, 50/50, 40/60 and 30/70 mass%, respectively (Section 2.2). The resins were tested either unfilled (UF) or filled with silanated barium glass fillers and fumed silica to 30/10 vol% (Low Fill, LF), and 50/10 vol% (High Fill, HF) respectively. The model composites were tested under the

hypothesis that if reciprocity held for the parent resins then it should also hold for the composites, i.e low viscosity resins are expected to fail reciprocity (Feng and Suh, 2007) regardless of whether they are filled or not. For UF resins, the differences in degree of conversion between the highest and lowest irradiation modes were greater for the low viscosity resins and degree of conversion was more dependent upon irradiation mode rather than resin type. The opposite occurred for the LF and HF counterparts where resin type became more influential. Therefore, it was suggested that the viscosity of the system determines how radicals are terminated. When the system viscosity is low, bi-radical termination occurs and polymer conversion becomes more dependent on exposure time but when the viscosity is high, a monomolecular termination pathway dominates and polymer conversion can rely, to a greater extent, on a reciprocal relationship between irradiance and time. Although the constituents of the resin matrix play a major role in the polymerization kinetics, it is evident that the fillers have a significant affect as they restrict mobility to a relatively small volume and increase the probability of a monomolecular termination pathway. Evidently bi-radical terminations even occur in highly filled composites with low local viscosities and therefore reciprocity cannot be verified.

However, not all flowable composites are manufactured in the same way and there exists scenarios where reciprocity can be verified even with low resin viscosities (Section 2.3). The use of alternative photoinitiators such as Lucirin TPO seems an interesting approach in terms of controlling polymerization kinetics as reactivity and efficiency is significantly greater when compared to camphoroquinone. Although inferior depth of cure may be a problem with such initiators, this may be overcome by improving optical properties with appropriate filler adaptation or by using camphoroquinone synergistically. As such, the optical properties of RBC materials will be investigated in Chapter 3 of this thesis.

In conclusion, the current chapter highlights the importance of manufacturers to divulge the composition of their materials to some extent in order to select adequate curing protocols which may reduce the occurrence of under-cured restorations. Also, the findings presented here may help the design of novel materials that allow high irradiance curing at reduced time which will benefit both the dentist and the patient.

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CHAPTER 3

Optical Phenomena of Photoactive Dental Resins

3.0 Optical Phenomena of Photoactive Dental Resins

The demand for filling materials to match the aesthetic quality of natural teeth has vastly increased in popularity over the last few decades. However, the durability of large ‘white’-fillings or resin-based composites (RBCs) in posterior load bearing applications is generally considered to be inferior to ‘silver’-filling types (dental amalgams). This shortcoming may be related to both the increased number of technique sensitive application steps the dentist has to perform when restoring the tooth as well as any inherent shortcomings of a quasi-brittle solid compared with amalgam alloy materials.

The placement of modern light activated RBCs is limited to approximately 2 mm increments which are cured ‘on command’ with high intensity blue light following the application and polymerisation of a bonding system. Therefore, to fill large cavities, adhesives and multiple increments are required, which increases procedure time and cost when compared with placing dental amalgam which can be filled in bulk. Chapter 2 of this thesis investigated the potential of high irradiance curing protocols to reduce procedure time and cost. Another interesting approach to reduce procedure time is to improve depth of cure, which will reduce the number of technique sensitive steps. Consequently, the development of RBCs with improved light transmission through depth is highly desirable which may improve depth of cure and reduce placement time.

The dynamic change in optical properties of RBCs is not well understood. Optical properties can be controlled through photoinitiator chemistry (Section 3.1) and also by altering the components of the resin matrix (Section 3.2), and therefore it may be possible to alter the extent to which the curing light is transmitted through the material. As RBCs cure, the photoinitiator decomposes and the resin components become denser and shrink. Consequently changes in optical properties occur. Shrinkage vectors are dependent upon several factors, amongst which material composition and cavity constraint are important. As shrinkage will affect light transmission, dynamic optical properties of the curing light may differ according to specimen geometry (Section 3.3). Therefore chapter aims to gain a better understanding of optical properties, which may provide a platform with which to improve cure depth, reduce procedure time and potentially improve clinical durability and success of such materials.

3.1 Competitive Light Absorbers in Photoactive Resin-Based Materials^{††‡‡}

3.1.1 Abstract

The aim of this Section is to characterise real time photoinitiator absorption properties of “white” Lucirin TPO and “yellow” camphoroquinone and to assess their influence on polymer discolouration. Resin formulations (50:50 Bis-GMA/TEGDMA mass %) contained low (0.0134 mol/dm³), intermediate (0.0405 mol/dm³) or high (0.0678 mol/dm³) concentrations of the photoinitiators camphoroquinone or Lucirin TPO and the inhibitor BHT at 0, 0.1 or 0.2% by mass. Disc shaped specimens (n=3) of each resin were cured for 60 s using a halogen light curing unit (~800 mW/cm²). Dynamic measurements of photoinitiator absorption, degree of conversion (DC)/ curing kinetics and reaction temperature were performed on curing resins. A spectrophotometer was used to measure CIELAB co-ordinates of each specimen before and after cure. Statistical analyses were performed using a general linear model (GLM) three-way analysis of variance (ANOVA) with photoinitiator (2 levels), photoinitiator concentration (3 levels) and BHT percentage (3 levels) as the independent variables. Supplementary one-way ANOVA and post-hoc Tukey tests (p=0.05) were also performed. The GLM of all independent variables revealed significant differences (p<0.001) where photoinitiator concentration (df=2; F=618.83) > photoinitiator type (df=1; F=176.12) > % BHT (df=2, F=13.17). BHT affected kinetics but consequently produced lower conversion in some of the camphoroquinone-based resins. Significant differences between photoinitiator type and concentrations were seen in colour (Lucirin TPO resins became yellow and camphoroquinone became less yellow). Reaction temperature, kinetics and conversion were also significantly different for both initiators (P<0.001). Despite Lucirin-TPO-based resins producing a visually perceptible colour change on cure, the colour change was significantly lower than that produced in camphoroquinone-based resins.

^{††} Hadis MA, Shortall AC and Palin WM. Competetive light absorbers in photoactive dental materials. Submitted to Acta Biomaterialia.

^{‡‡} Winner of the Heraeus Kulzer Travel award for innovation in materials testing. International Association of Dental Research, Barcelona, 2010.

3.1.2 Introduction

Over the last decade, the significant increase in the use of photo-active RBCs over dental amalgams in restorative dentistry is primarily due to patients aesthetic demands, the dentists material preference to use RBCs and environmental concerns and government legislations (Burke, 2004). The development of RBCs has continued since they were first introduced several decades ago with much research being carried out in order to improve the setting reaction and photoinitiator chemistry (the primary absorber of light in RBCs) (Yoshida and Greener, 1994; Rueggeberg et al., 1997; Halvorson et al., 2002; Jakubiak et al., 2003; Neumann et al., 2005; Neumann et al., 2006; Ogunyinka et al., 2007; Musanje et al., 2009) as-well as mechanical properties (Kahler et al., 2008; Curtis et al., 2009).

The α -diketone, camphoroquinone is the most widely and successfully used photoinitiator in RBC materials (Stansbury, 2000; Jakubiak et al., 2003) and has a distinct yellow tint, which does not bleach fully upon irradiation. Such discolouration may reduce aesthetic quality (Rueggeberg et al., 1997; Suh, 1999; Brackett et al., 2007) and affect light transport through material thickness. Furthermore, the use of camphoroquinone also requires amine co-initiators such as dimethylamino ethyl methacrylate (DMAEMA) to efficiently produce free radicals to initiate polymerisation. The use of such co-initiators is known to produce coloured reaction by-products in the presence of light or heat (Asmussen, 1985). Any colouring by the photoinitiators will inevitably reduce aesthetic quality and cause problems in colour matching to natural teeth (Alvim et al., 2007). Despite this, the camphoroquinone/amine system remains the most widely used photoinitiator system in photoactive dental RBCs although more recently alternative “colourless” photoinitiators have been introduced which may be based on either iodonium salts (Weinmann et al., 2005), onium compounds (Shin and Rawis, 2009) or Norrish Type I photoinitiator systems such as acylphosphine oxides (Neumann et al., 2005; Illie and Hickel, 2005; Neumann et al., 2006). Some of these types of photoinitiators have recently been used by manufacturers either synergistically with camphoroquinone or as a stand alone photoinitiator system, which may improve polymerisation kinetics (Chapter

2), mechanical properties and aesthetic quality (Park et al., 1999; Weinmann et al., 2005; Neumann et al., 2006; Shin and Rawis, 2009).

Free-radical polymerisation of methacrylate monomers used in dental RBCs is initiated by the excitation of suitable photoinitiator systems by light despite the monomer itself being capable of forming reactive species although this is deemed inefficient (Sun and Chae, 2000; Shin and Rawis, 2009). The photoinitiator chemistry affects light transmission, polymerisation kinetics and overall conversion of monomer to polymer thus affecting material properties (Yoshida and Greener, 1994; Rueggeberg et al., 1997; Ogunyinka et al., 2007; Musanje et al., 2009). It is well known that an increase in the amount of photoinitiator in composite resins or an increase in the spectral overlap between the light curing unit and the initiating system (i.e. synergistically using more than one photoinitiator) improves polymerisation kinetics, aesthetics, mechanical and biological properties (Rueggeberg et al., 1997; Park et al., 1999). However, there exists an optimal level above which any increase in photoinitiator does not benefit final conversion (Yoshida and Greener, 1994; Musanje et al., 2009) and aesthetic quality may be compromised in the case of pigmented photoinitiators.

The increased demand for aesthetic quality has introduced the use of alternative colourless photoinitiators, which may provide the potential for improved aesthetic (Price and Felix, 2009) and curing efficiency (Chapter 2; Neumann et al., 2005; Neumann et al., 2006). Furthermore, photoinitiators based on acylphosphine oxides such as Lucirin TPO (2,4,6-trimethylbenzoyl-diphenylphosphine oxide) do not require the use of amine accelerators, which further increases the potential for aesthetic quality. The photochemical cleavage of the carbon phosphorus bond in such molecules leads to the formation of two free radical intermediates, which can initiate polymerisation without the need for additional amine accelerators (Neumann et al., 2006).

Camphoroquinone has an absorption range from 400-500nm with a maximum excitation wavelength of 470nm. Consequently, most light curing units are adapted for the excitation of camphoroquinone. Lucirin TPO seems to be an interesting molecule despite its shorter wavelength absorption range since its higher molar absorptivity may provide greater efficiency to cure dental resins (Chapter 2; Neumann et al., 2005; Neumann et al., 2006). Moreover, recent work has highlighted the potential of Lucirin

TPO-based (unfilled) polymers to be cured in thick layers (Stephenson Kenning et al., 2008). However, in Section 2.3 it was suggested that inferior cure depth in resins and RBCs containing Lucirin TPO when compared to camphoroquinone based materials may be related to inefficient light transmission at depth. Additionally, the use of this initiator is not compatible with most conventional light emitting diode (LED) light curing units, which have an emission range between 450-500nm. Several manufacturers have now introduced so-called third generation dual wavelength LED units by combining LEDs with two different wavelengths that are capable of exciting alternative photoinitiators (Illie and Hickel, 2008).

Consequently, the aim of this study was to analyse the photo-absorption characteristics and photo-bleaching characteristics of camphoroquinone and Lucirin TPO throughout cure using a dynamic spectroscopy technique, which may have some significance to improving light transmission and the perceived improvement of aesthetic quality using alternative, non-yellowing photoinitiator systems. The two hypotheses that were tested were:

- i) Camphoroquinone will exhibit a decrease in yellowing throughout and following cure
- ii) Lucirin TPO will not result in a significant change in colour during or following cure.

3.1.3 Materials and Methods

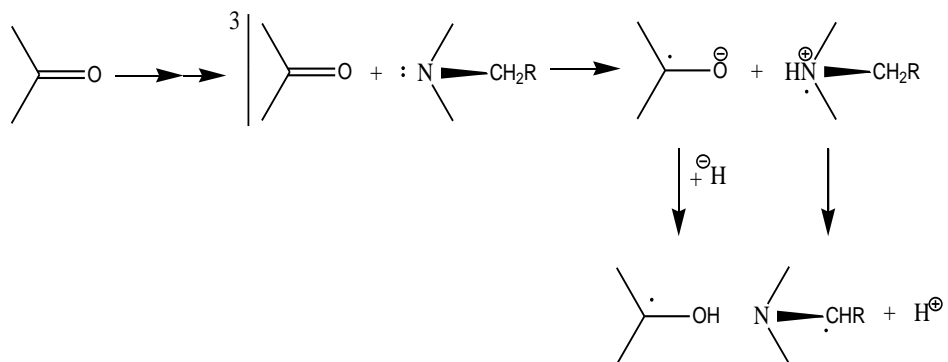
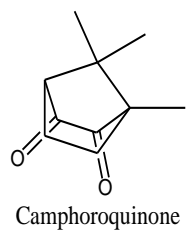
Resin formulations were prepared using a 50/50 Bis-GMA and TEGDMA mixture containing low (0.0134 mol/dm^3), intermediate (0.0405 mol/dm^3) and high (0.0678 mol/dm^3) concentrations of the photoinitiators camphoroquinone or Lucirin-TPO (Table 3.1.1).

[PI]	Weight %		
	Camphoroquinone	DMAEMA	LUCRIN TPO
Low	0.2	0.8	0.4212
Intermediate	0.6	1.2	1.2644
High	1	1.6	2.1062

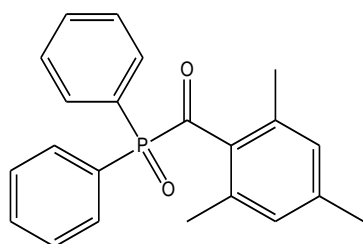
Table 3.1.1: Photoinitaitor concentration as percentage of total weight of resin having equimolar concentrations of camphoroquinone and Lucirin TPO.

Resins were obtained by Sigma-Aldrich (Gillingham, UK) and either used as received or contained the inhibitor butylated hydroxytoluene (BHT) at 0.1% and 0.2% (by mass). The chemical structures of the photoinitiators and the mechanism by which free radicals are produced and terminated are shown in Figure 3.1.1. A resin containing no photoinitiator and no inhibitor was used as the control. A quartz-tungsten halogen light curing unit (XL2500, 3M) having a total irradiance of $\sim 800 \text{ mW/cm}^2$ measured on a handheld radiometer (Cotolux light meter; Cotelene Whaledent, UK) and verified using a UV-Vis spectrometer (USB400, Ocean Optics, Dunedin, USA) was used to cure the specimens for 60 s. The light curing unit was chosen specifically for its broad spectrum irradiance which overlaps the absorption spectrums of both photoinitiators (albeit only partially with Lucirin TPO). The emission spectrum was measured by calibrating the UV-Vis spectrometer using a deuterium tungsten light source (according to NIST standard) with known spectral output in the UV-Vis and NIR range (Mikropack DH2000-CAL, Ocean Optics, Dunedin, USA). During the measurement of absolute irradiance, the irradiance of the light curing unit was controlled by stacking reflective neutral density filters (Thorlabs UK) which produced a 6-fold reduction in irradiance over the whole spectrum. Static absorption spectra of the photoinitiators and the inhibitor BHT were determined in methyl methacrylate (Sigma-Aldrich, UK) using UV-Vis spectroscopy.

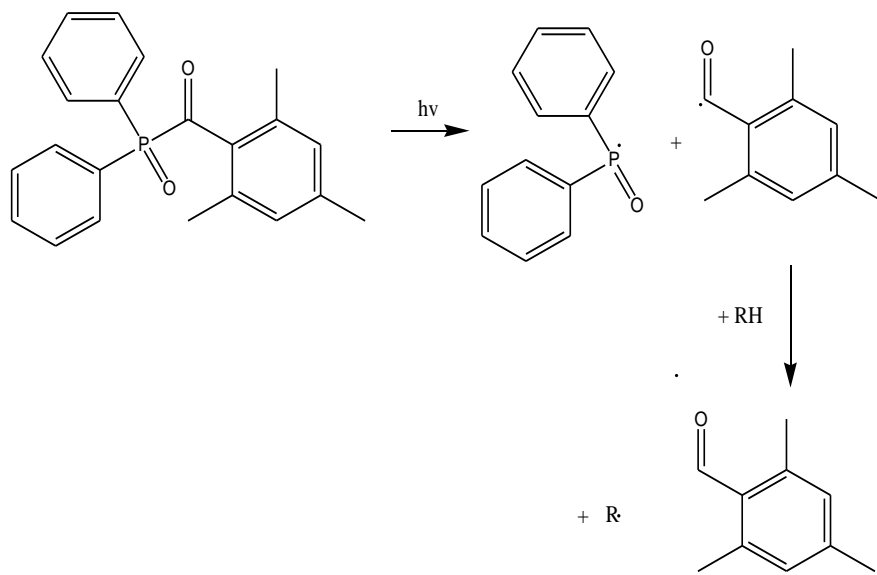
(a)



(b)



2,4,6-trimethylbenzoyl-diphenylphosphine oxide
(Lucirin TPO)



(c) *continued...*

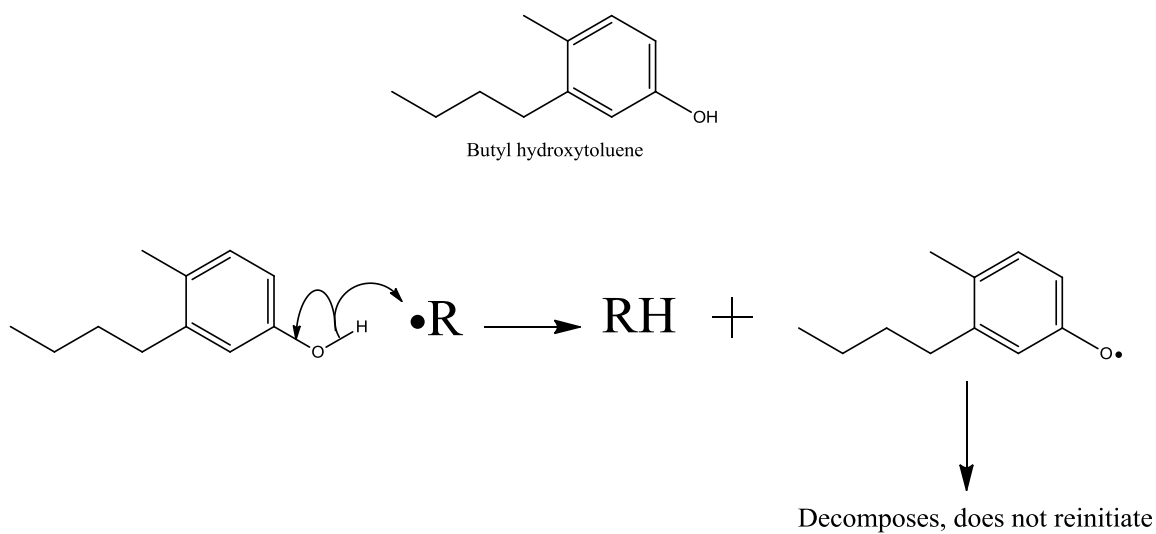


Figure 3.1.1: Free radical (•) production with the photoinitiator system camphorquinone/amine (a) Lucirin TPO (b) and the termination of free radicals by the inhibitor BHT (c).

Degree of Conversion and Kinetics

Dynamic measurements of DC/ curing kinetics were conducted in real time for curing resins (n=3) through cylindrical white Teflon moulds (12 x 1.4 mm) covered by two glass microscope slides (Figure 3.1.2). The FT-NIR method employed is described previously in Chapter 2 of this thesis.

Temperature Measurements

Temperature measurements were also conducted in real-time during light activation (n=3) through cylindrical white Teflon moulds (12 mm central circular cavity diameter x 1.4 mm thickness) covered by two glass microscope slides. The method employed was previously described in Section 2.3 where a thermocouple bead was inserted and glued into a mould which had been split on one side (Figure 3.1.2). The thermocouple was calibrated against a standard mercury thermometer. The temperature measurements were then used to calculate the maximum temperature rise (ΔT^{\max}) during irradiation.

Photoinitiator Absorption

Additionally, photoinitiator absorption was measured in real time during irradiation at the maximum point of overlap between the photoinitiators and the light curing unit (camphoroquinone: 470nm; Lucirin TPO: 420nm) by a technique adapted from Chen et al., (2007) (Figure 3.1.2). The UV-Vis spectrometer (USB4000, Ocean Optics) was set up so that a 200 μm fibre optic cable, which had a cosine corrector attachment (CC3-UV; Ocean Optics, USA; light sensing diameter: 3.9 mm) was able to measure light of the appropriate wavelength from the centre of the light curing unit. Due to the relatively high intensity of the light curing unit, reflective neutral density filters were stacked directly above the cosine corrector. The spectrometer was then calibrated so that the 'light spectrum' was taken with a sample within a cylindrical teflon mould (12 x 1.4 mm) containing no photoinitiator (covered with two microscope slides). The 'dark

spectrum' was then taken by blocking any light reaching the optical fibre. The distance between the top of the sample and the light curing unit was kept to 1 mm, which facilitated the use of a shutter. A shutter was used as preliminary studies identified fluctuations in light intensity during the first 10 s of activation and therefore the light was activated for 10 s before the shutter was removed to start both calibration and sample measurements. For sample measurements, real time absorbance was recorded using Spectrasuite software during irradiation of the sample.

Colour Analysis

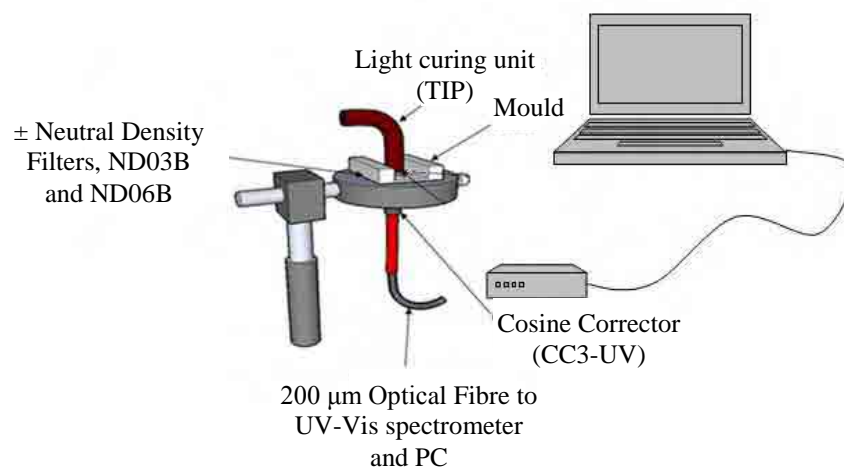
A spectrophotometer (CM-2600d; Konica, Minolta, UK) was used to measure the tristimulus values of the International Commission on Illumination (CIELAB) values based on D65 standard illuminant and a 10-degree observer. The tristimulus values represent the amount of three primary stimuli (red, green and blue) required to give colour match with the colour stimulus. The spectrophotometer was first calibrated against a standard white tile provided by the manufacturer of the spectrophotometer. Colour measurements were made before and after real time photoinitiator absorption measurements (described above). The measurements were made whilst the samples were contained within the moulds and the glass microscope slides but were placed on top of a white ceramic backing tile before each measurements. The CIELAB values obtained included specular gloss (SCI) as the specular port was covered with a diffuse white material and specular gloss was reflected back into the the integrating sphere. From these values, the colour change was calculated according to the following equation:

$$\Delta E = \sqrt{(L_2^* - L_1^*)^2 + (a_2^* - a_1^*)^2 + (b_2^* - b_1^*)^2}$$

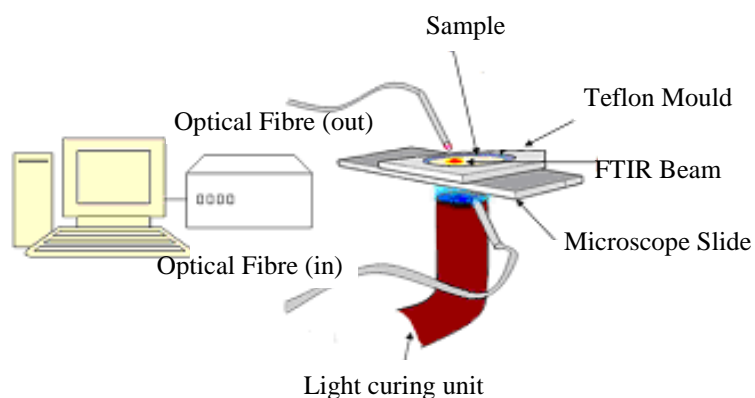
Statistical Analysis

Statistical analyses were performed on the DC^{\max} , R_p^{\max} , ΔT^{\max} , ΔE (change in colour) and Δb (the change in b-value of the CIELAB co-ordinates) and data sets using a general linear model (GLM) three-way analysis of variance (ANOVA) with photoinitiator (2 levels), photoinitiator concentration (3 levels) and BHT concentration (3 levels) as the independent variables. Supplementary one-way ANOVA and post-hoc Tukey tests ($P=0.05$) were also undertaken.

(a)



(b)



(c)

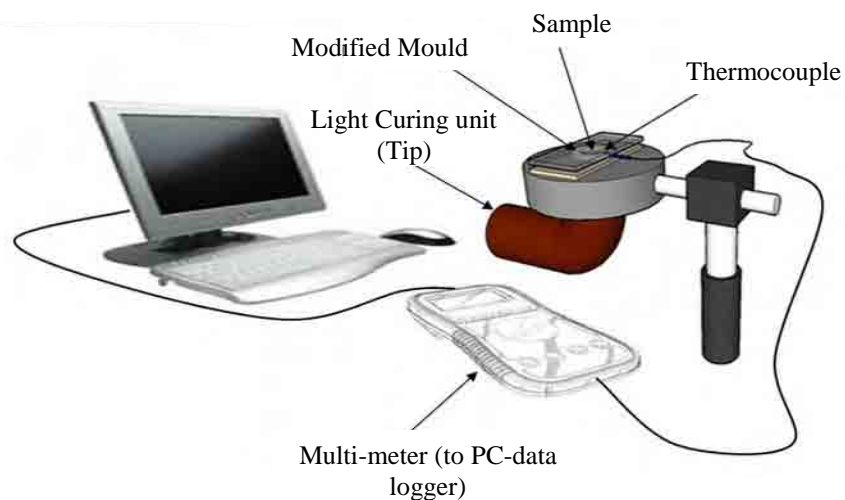


Figure 3.1.2: Experimental set up for the measurement of real time photoinitiator absorption (a), degree of conversion (b) and temperature (c).

3.1.4 Results

Figure 3.1.3 presents a comparison between the absorption spectra of both photoinitiators (camphoroquinone and Lucirin TPO), the inhibitor BHT and the emission spectra of the XL2500 light curing unit. The spectral emission of the light curing unit was better adapted to camphoroquinone absorption, as a greater overlap of the two spectra was evident although Lucirin TPO shows a significantly higher molar absorptivity at λ^{\max} . The absorption range of BHT fell outside the emission range of the light curing unit and absorbed in the UV region of the electromagnetic spectrum.

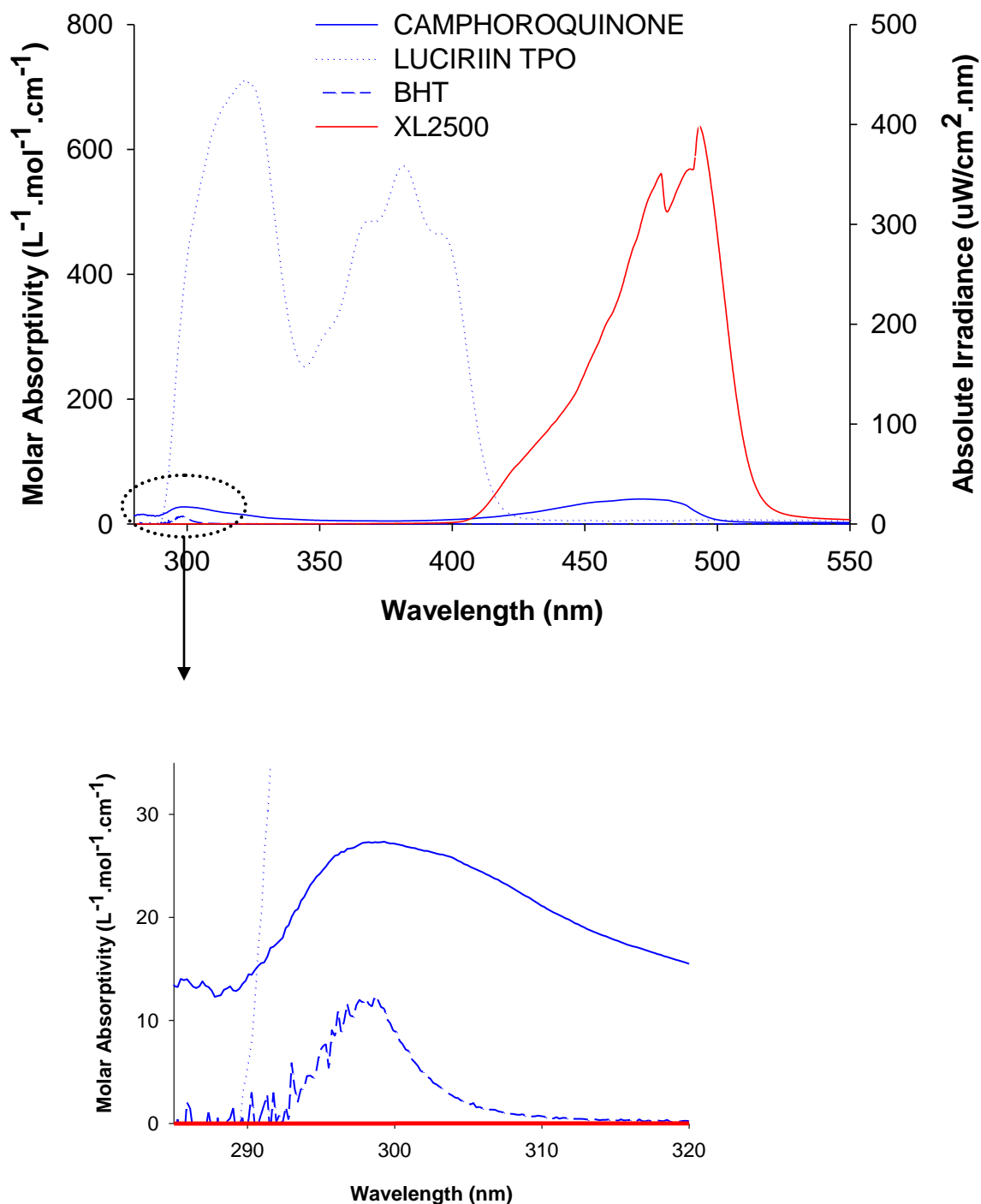
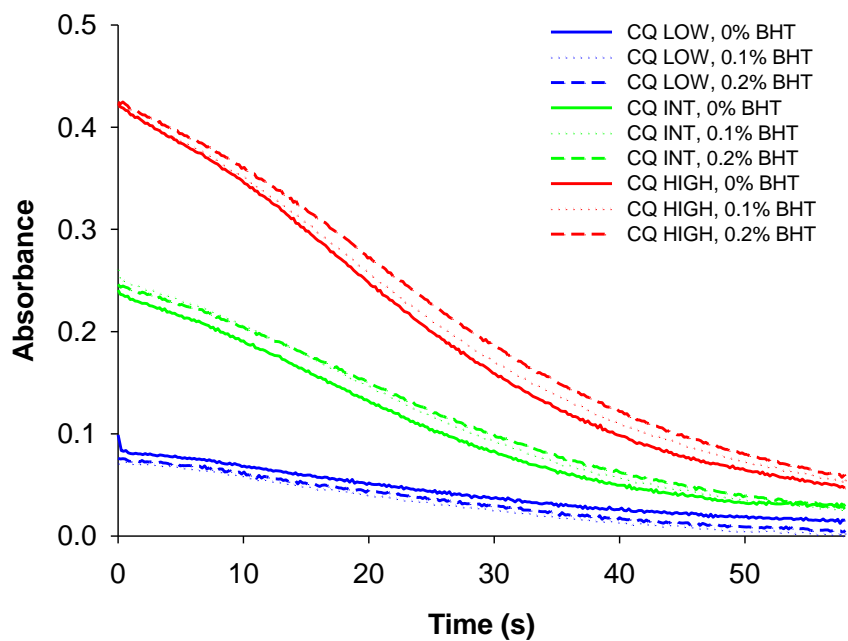


Figure 3.1.3: Absorption spectra (left axis, in $\text{L} \cdot \text{Mol}^{-1} \cdot \text{cm}^{-1}$) of camphoroquinone (solid blue line), Lucirin TPO (broken blue line) and BHT (dotted blue line) compared to emission spectra of the XL2500 (right axis, in $\mu\text{W}/\text{cm}^2 \cdot \text{nm}$) (a). Magnified peak for BHT (b).

Figure 3.1.4 represents the change in photoinitiator absorption for (a) camphoroquinone and (b) Lucirin TPO at the emission-absorption overlap. The concentration of photoinitiator significantly affects both initial and final absorbance in both camphoroquinone and Lucirin TPO ($P < 0.001$). BHT did not significantly affect initial and final absorbance values for any of the camphoroquinone or Lucirin TPO although kinetics was affected.

(a)



(b)

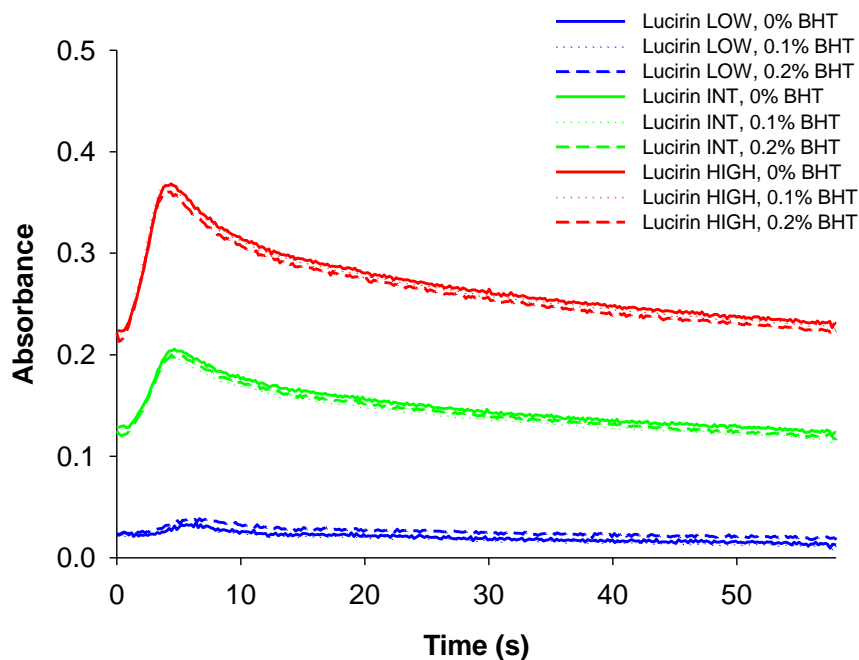


Figure 3.1.4: Real time change in PI absorption of (a) camphoroquinone (CQ) (b) and Lucirin TPO, measured at the maximum point of overlap of the LCU (camphoroquinone: 470nm; Lucirin TPO: 420nm) by UV-Vis spectroscopy.

The GLM of all independent variables revealed significant differences ($P < 0.001$) where photoinitiator concentration ($df=2$; $F=618.83$) $>$ photoinitiator type ($df=1$; $F=176.12$) $>$ % BHT ($df=2$, $F=13.17$). Separate one-way ANOVAs and post hoc-Tukey tests ($p < 0.05$) revealed significant overall differences for each separate response to compare the effect of photoinitiator concentration and BHT concentrations within each photoinitiator group (Table 3.1.2). Camphoroquinone resins with similar BHT concentrations revealed significantly higher ΔE and Δb values as photoinitiator concentration was increased which was related to the significantly higher *b-values* ($p < 0.05$) for both pre- and post irradiation (although *b-values* decreased upon irradiation resulting in a negative Δb). Similar trends were also observed for Lucirin TPO-based resins although ΔE and Δb (also related to the *b-values*) values significantly lower ($p < 0.05$) for all Lucirin TPO resins compared with camphoroquinone-based resins. For Lucirin TPO resins, Δb values were positive as the *b-value* increased upon irradiation (Table 3.1.3). Resins containing no photoinitiator showed no significant colour change so ΔE and Δb should be taken as zero.

PI	BHT	COLOUR CHANGE				TEMPERATURE	CONVERSION
		ΔE	<i>b-value</i> pre-cure	<i>b-value</i> post-cure	Δb	ΔT^{\max}	DC ^{MAX}
NO PI	0%	0.56 (0.38)	3.42 (0.25)	3.29 (0.01)	-0.13 (0.25)	0.43 (0.12)	-
Camphoroquinone LOW	0%	8.64 (0.50) d	15.61 (0.33) d	7.10 (0.50) c	-8.51 (0.51) d	20.17 (1.62) a	70.53 (0.37) d
	0.10%	7.62 (0.38) d	15.50 (0.13) d	7.93 (0.51) c	-7.57 (0.38) d	19.70 (1.01) a	70.95 (0.27) d
	0.20%	8.41 (0.23) d	15.41 (0.19) d	7.15 (0.05) c	-8.26 (0.18) d	17.50 (0.72) b	69.57 (0.45) d
Camphoroquinone INT	0%	24.37 (0.29) c	35.25 (0.17) c	11.10 (0.28) b	-24.16 (0.35) c	23.67 (0.67) a	77.28 (0.32) b
	0.10%	24.47 (0.74) c	35.38 (0.47) c	11.15 (1.14) b	-24.23 (0.74) c	22.70 (0.36) a	76.07 (0.32) c
	0.20%	23.74 (0.38) c	35.16 (0.36) c	11.62 (0.48) b	-23.54 (0.41) c	23.47 (1.93) a	75.45 (0.13) c
Camphoroquinone HIGH	0%	36.61 (0.63) a	51.35 (0.08) a	14.94 (0.64) a	-36.40 (0.60) a	22.47 (0.40) a	80.36 (0.40) a
	0.10%	33.47 (0.84) b	49.08 (0.35) b	15.79 (0.91) a	-33.29 (0.86) b	22.67 (2.54) a	77.97 (0.21) b
	0.20%	33.26 (0.84) b	49.44 (0.74) b	16.37 (1.11) a	-33.08 (0.81) b	23.63 (0.80) a	77.95 (0.27) b

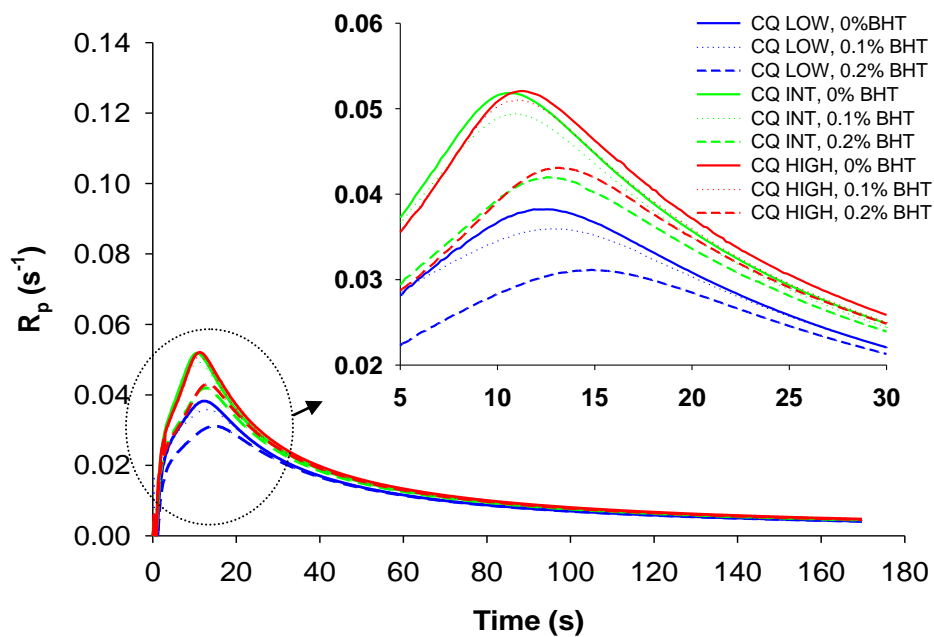
Table 3.1.2: Colour change, temperature rise and DC for resins containing the photoinitiator Camphoroquinone. Similar letters within the same column represents no statistical significance (p<0.05).

PI	BHT	COLOUR CHANGE				TEMPERATURE	CONVERSION
		ΔE	<i>b-value</i> pre-cure	<i>b-value</i> post-cure	Δb	ΔT^{\max}	DC ^{MAX}
LUCIRIN LOW	0%	1.15 (0.21) b	5.21 (0.19) c	6.13 (0.20) c	0.92 (0.03) b	28.40 (0.72) a	68.30 (0.27) c
	0.10%	1.03 (0.03) b	5.13 (0.13) c	6.10 (0.10) c	0.97 (0.09) b	30.97 (2.94) ab	68.78 (0.15) c
	0.20%	0.89 (0.12) b	5.18 (0.10) c	5.74 (0.41) c	0.56 (0.49) b	31.37 (1.12) ab	66.46 (0.19) d
LUCIRIN INT	0%	2.06 (0.33) ab	6.36 (0.08) b	8.38 (0.38) ab	2.01 (0.33) a	28.67 (2.06) b	74.24 (0.33) ab
	0.10%	1.85 (0.07) ab	6.17 (0.09) b	7.94 (0.10) b	1.78 (0.09) a	31.87 (2.54) ab	75.45 (0.13) ab
	0.20%	1.21 (0.81) b	6.22 (0.08) b	7.37 (0.85) b	1.15 (0.84) a	34.43 (2.50) a	73.84 (0.41) a
LUCIRIN HIGH	0%	2.25 (0.11) a	7.14 (0.06) a	9.29 (0.06) a	2.15 (0.50) a	34.47 (0.31) a	76.05 (0.70) a
	0.10%	2.12 (0.24) a	6.85 (0.10) a	8.91 (0.32) a	2.06 (0.24) a	35.33 (1.71) a	77.95 (0.27) a
	0.20%	2.08 (0.61) a	6.97 (0.12) a	8.97 (0.69) a	2.00 (0.59) a	33.57 (1.42) ab	76.09 (0.45) a

Table 3.1.3: Colour change, temperature rise and DC for resins containing the photoinitiator Lucirin TPO. Similar letters within the same column represents no statistical significance (p<0.05).

For Lucirin TPO based materials, each resin exhibited significantly higher R_p^{\max} ($p < 0.05$) than camphoroquinone based resins of similar concentrations although DC^{\max} was significantly lower ($p < 0.05$). For both camphoroquinone and Lucirin TPO, the addition of BHT significantly reduced R_p^{\max} ($P < 0.005$) which significantly affected DC^{\max} in some of the resins ($p > 0.05$) (Figure 3.1.5; Table 3.1.2 and Table 3.1.3). The temperature rise was significantly higher in all Lucirin TPO resins ($p < 0.05$) with the highest temperature rise being observed for high concentration. The addition of BHT did not significantly change ΔT^{\max} ($p < 0.05$) for any of the resins apart from the low concentration of camphoroquinone that contained 0.20 wt% BHT which showed significantly lower ΔT^{\max} ($p < 0.05$). The temperature rise in resins containing no photoinitiator was 0.43 ± 0.12 °C due to heat dissipated by the light curing unit and was used to correct resins containing photoinitiator to obtain temperature rise during polymerisation.

(a)



(b)

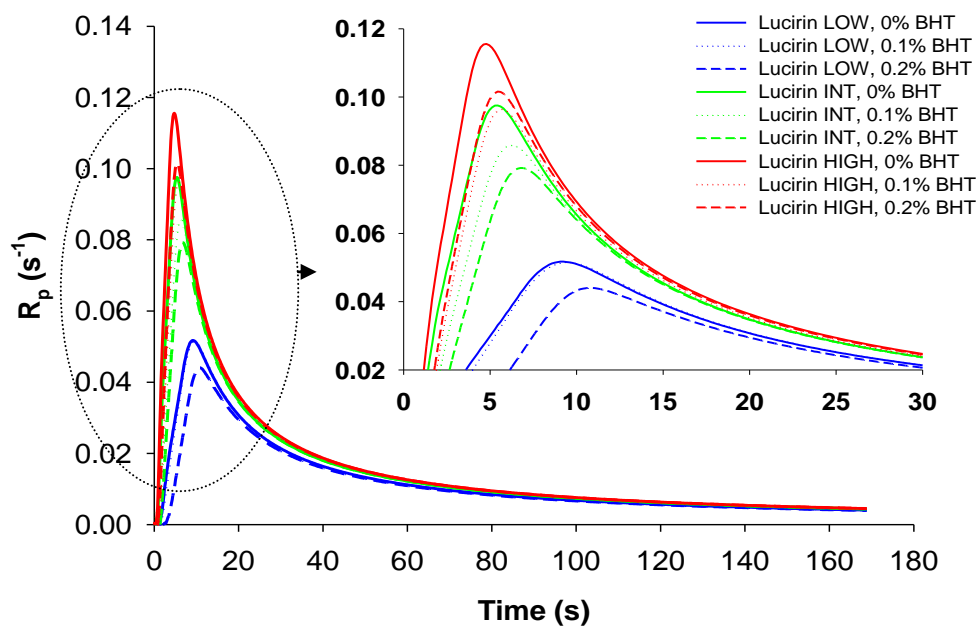


Figure 3.1.5: R_p curves for (a) camphoroquinone (CQ) and (b) Lucirin TPO measured in real-time using FT-NIR spectroscopy.

3.1.5 Discussion

Light cured photoactive resins are composed of photoinitiator systems that absorb light to form excited states that initiate polymerisation. The proper combination of photoinitiator and co-initiator as well as light source and exposure time are regarded as major factors to optimise light curing of dental composites (Emami and Söderholm, 2005) to improve conversion in order to attain better mechanical properties, clinical performance as-well as colour stability and biocompatibility (Shin and Rawls, 2009). The concentration of photoinitiator and co-initiator are also key factors that influence curing in dental composites where generally higher concentrations lead to better conversion (Musanje et al., 2009) although the use of high concentrations of coloured photoinitiators such as camphoroquinone may lead to undesirable photo-yellowing effects (Rueggeberg et al., 1997).

The rapid rate of polymerisation in Lucirin TPO resins compared to camphoroquinone-based resins (Table 3.1.2, Table 3.1.3 and Figure 3.1.) is due to the photochemical cleavage of the carbon-phosphorus bond resulting in the formation of two polymerisation initiating radicals (an acyl and phosphonyl radical) which are more reactive than the one active radical expected in the camphoroquinone amine system (Neumann et al., 2006; Figure 3.1.1a). Although Lucirin TPO is reported to be a more efficient photoinitiator due to its higher molar absorptivity, the phosphonyl radical is considered to be more efficient for polymerisation than the acyl radical (Neumann et al., 2006). In the case of camphoroquinone without amine co-initiators, homolytic cleavage of the C-C bond between the two carbonyl groups results in the formation of carbonyl radicals only. These radical pairs can undergo cage escape (where the radical site moves around the structure) to form photo-decomposed products which results in a decrease in the concentration of the active camphoroquinone molecules thus decreasing absorbance and the *b-value* of the CIELAB coordinates. However, the two carbonyl radicals in camphoroquinone remain structurally connected to each other and the probability of recombination to reform camphoroquinone molecules is great (Sun and Chae, 2000). Recombination can be greatly reduced by the addition of amine accelerators. Irradiation results in the formation of an exciplex (excited state complex) via electron transfer from

amine to the excited di-ketone. When amine accelerators are added, the formation of an aminyl radical (responsible for the initiation of polymerisation) and a hydryl radical (which is almost inactive) are a result of hydrogen transfer from amine to the excited diketone in the exciplex. (Figure 3.1.1b; Shin and Rawls, 2009). Irradiation leads to the decomposition of the photoinitiator system over time until the photoinitiator is consumed (Figure 3.1.4). The amine accelerator serves to increase the rate of photodecomposition (Sun and Chae, 2000) although the rate of decomposition is much slower than in Lucirin TPO (Figure 3.1.5). The rate of decomposition and absorbance for both photoinitiator is proportional to their concentrations as the likelihood of excitation is greatly increased when the molar concentrations of absorbing species (photoinitiator) are increased. In the present study, the ratio of camphoroquinone and amine was selected as an optimum combination to produce maximum *DC* (Yoshida and Greener, 1994) and camphoroquinone concentration was similar to the one used in commercial composites (Taira et al., 1988), and an equimolar concentration was used for Lucirin TPO for comparative purposes. Increasing photoinitiator concentration serves to accelerate the reaction, consequently leading to greater R_p^{\max} and ΔT^{\max} at a reduced time (Figure 3.1.5; Tables 3.1.2 and 3.1.3).

BHT is commonly used in RBCs due to its anti-oxidant properties which help improve shelf-life and operatory light stability during chairside placement but may also influence polymerisation kinetics and degree of cure (Rosentritt et al., 2010). Reduction in reaction speed by chemical inhibition occurs as free radicals are terminated by reacting with the phenolic hydrogens of the BHT molecule (Figure 3.1.1c) and thus increase the rate of recombination and prolongs absorption by the photoinitiators with no dramatic change in temperature or conversion (Table 3.1.2 and Table 3.1.3). The conversion of monomer to polymer proceeds at a reduced rate until the inhibitor is consumed (Figure 3.1.5). Additionally, it should also be mentioned that ‘off the shelf’ monomers typically contain inhibitors to improve storability and shelf life and in this work these monomers were used without purification. Therefore the subtle difference between 0% BHT and 0.1% BHT on cure kinetics, *DC* and colour (Figure 3.1.5; Tables 3.1.2 and 3.1.3) may relate to trace impurities of inhibitors. This may also account for the non-significant differences found in the same resins containing different BHT concentrations rather than

insensitivity of the analytical techniques employed (Tables 3.1.2 and 3.1.3). This is in agreement with previous work showing reaction time increasing significantly when BHT concentrations are above 0.25% (Rosentritt et al., 2010)

For RBCs it has been reported that a visually perceptible colour change occurs when ΔE exceeds 5.4 (Cook and Chong, 1985). However, a ΔE value that exceeds 2 is considered representative of a visually perceptible colour change for resins in the yellow-red range (*b-value*) of the colour spectrum (Ferracane et al., 1985). As the current investigation is based only on unfilled resins, it is more relevant to base the discussion around the latter value. For Lucirin TPO, visually perceptible colour changes occurred for the intermediate and high concentrations of Lucirin TPO, which may also occur in filled systems (Table 3.1.3). Therefore, where so-called colourless photoinitiators might be perceived to improve aesthetic quality of bleach white or translucent shades, they may actually cause significant colour contamination and reduce shade matching if used at high concentrations. Colour change may also be affected by several other factors such as temperature, moisture and the presence of oxygen (Barres, 1945; Shin and Rawls, 2009). One way to reduce the colour contamination is by the addition of inhibitors, which act as anti-oxidants and reduce the oxidative by-products. The addition of BHT did not significantly influence colour change for any of the tested resins although there was a noticeable trend of decreased yellowing with increased BHT concentration for most of the resins containing Lucirin TPO, which was not observed for camphoroquinone-based resins as camphoroquinone is strongly pigmented (Table 3.1.2 and Table 3.1.3). This may be due to the anti-oxidant properties of BHT (Lambert et al., 1996; Rosentritt et al., 2010), which may reduce any oxidative by-products that are produced during the reaction. The non-significant data may be due to the limited concentration of BHT used in the present study (≤ 0.2 wt%) and the trace amounts of inhibitors already present in the monomer. Nonetheless, chemical inhibition will reduce R_p^{\max} (Figure 3.1.5) which will reduce the rate of conversion and the rate of colour change therefore affecting the final colour change.

Yellowing of camphoroquinone resins is proportional to the concentration of camphoroquinone, however, photo-activation causes camphoroquinone to decompose and decreases the concentration of camphoroquinone causing a photo-bleaching effect

(decrease in *b-value*). Contrary to this, Lucirin TPO resins (initially colourless) show an initial increase in absorption (Figure 3.1.4b) as a consequence of an increase in the *b-value* of the CIELAB co-ordinates (Table 3.1.3) in the initial stages of polymerisation. Since Lucirin TPO is a colourless photoinitiator and exhibits efficient photo-bleaching properties with no absorbance of photolysis by-products above 360 nm (Stephenson Kenning et al., 2006; Stephenson Kenning et al., 2008), the increase in *b-value* may be related to monomer discolouration. The vinyl groups of the monomers will readily react with oxygen to form coloured peroxides (Barres, 1945). The yellowing of Bis-GMA monomer under UV illumination was also reported by Ferracane et al., (1985). Consequently, the presence of oxygen and the higher reaction temperatures and kinetics may have lead to the formation of coloured peroxides that absorb in the visible region of the electromagnetic spectrum which are the probable cause of the yellow discolouration observed in Lucirin TPO resins. Also, it should be noted that the amine (not used in Lucirin TPO resins) and monomers do not absorb significantly between 420-520nm (Asmussen and Vallo, 2009). It is also worth noting that BHT has an absorbance peak in the UV region (290-310 nm) of the electromagnetic spectrum and thus falls outside the measured range (470 nm for camphoroquinone and 420 nm for Lucirin TPO). Therefore the measured absorbance is unlikely to be related to BHT. Additionally, the light curing unit has an emission range between 400-550 nm and therefore it is expected that the addition of BHT will not significantly affect light transmission through absorption by BHT.

Polymerisation is usually accompanied by shrinkage and a decrease in volume and possibly results in a path length change of approximately 3-4% that occurs during the first 20 s of irradiation with a further 1-2% occurring thereafter (Mucci et al., 2009). As the measured absorbance is the product of the molar absorptivity and concentration of photoinitiator multiplied by the optical path length (Beer-lamberts law) any changes in the path length will affect absorbance (where an increase in path length, i.e. expansion will increase absorbance). Indeed, it may well be that for Lucirin TPO resins, volumetric expansion by high reaction temperatures caused an initial increase in absorbance (Figure 3.1.4b). However, at the first stage of fast reaction, the sample shrinkage and expansion is believed to be very slow and changes in sample thickness can be disregarded (Asmussen

and Vallo, 2009). This was further verified using a interferometry technique (Section 3.2), the results of which showed no such expansion during irradiation and can be found in the Appendix of this thesis. Nonetheless, the presence of Lucirin TPO serves to accelerate the colour change and this effect is related to photoinitiator concentration (Lucirin High > Lucirin Intermediate > Lucirin Low) where intermediate and high concentrations were deemed to be visually perceptible. Although the criteria for visual perception vary in the literature any discolouration as a result of the photo-polymerisation process of so-called “colourless” photoinitators may not only affect aesthetic quality but also negatively affect light transmission and resultant *DC*.

As photoactivation leads to the decomposition of the photoinitators, the first hypothesis was accepted as all camphoroquinone-based resins exhibited a decrease in yellowing throughout and following cure. Although Lucirin TPO has been considered a colourless photoinitiator, monomers may become discoloured upon irradiation due to high R_p^{\max} and temperatures, which may produce coloured oxidative by-products and therefore the second hypothesis was rejected. In the present study, the absorption characteristics and colour change of only two photoinitiator (camphoroquinone and Lucirin TPO) were studied in unfilled resins. Therefore the absorption characteristics and colour change of these initiators and others in filled composites should be the subject of further studies as any scattering and attenuation of light (Brie et al., 1999) will inevitably affect the absorption characteristics and colour change. Also, since optical properties are likely to be directly related to shrinkage of the curing polymer, the change in physical and optical properties of photoactive resins will be investigated in Sections 3.2 and 3.3 of this thesis. Here, the understanding of optical absorption properties is of prime importance in order to improve photoinitiator chemistry and clinical success of RBCs.

3.1.6 Conclusion

Although TPO, a “colourless” photoinitiator, exhibits greater absorption efficiency compared with camphoroquinone the formation of coloured peroxides was the likely cause of the yellow discolouration observed following increased rates of polymerisation.

The dynamic understanding of photoinitiator absorption is of fundamental importance as any unconverted photoinitiator can cause aesthetic imperfections and reduced colour stability. Furthermore, competitive light absorbers from reaction by-products from monomer or photoinitiator may reduce light transmission through depth and affect final aesthetic quality of dental composite restorations.

Dynamic measurements during photo-polymerisation yields important information on curing characteristics of photoactive resins which may be used in order to significantly improve light transmission. Not only is the photoinitiator important for optical properties, the resin matrix of RBC materials seems an important factor to control optical and physical change which will be investigated in Section 3.2 of this thesis. Improving knowledge of optical properties may ultimately improve curing depths, *DC*, colour stability and mechanical properties of RBCs as well as have potential for reduced procedure times.

3.2 Dynamic Monitoring of Refractive Index Change Through Photoactive Resins^{§§}

3.2.1 Abstract

The change in optical characteristics through the bulk of curing photopolymers is not fully understood. Photopolymerisation processes are accompanied by photoinitiator absorption, density changes and volumetric shrinkage, which alter optical properties and may affect curing efficiency through depth. This investigation demonstrates the use of a novel low coherence interferometry technique for simultaneous measurement of optical (refractive index) and physical (shrinkage) properties throughout curing of photoactive monomers containing various concentrations of bisphenol-A diglycidyl ether dimethacrylate and triethylene glycol dimethacrylate. Reliability of the interferometry technique was compared with an Abbé refractometer and showed a significant linear regression ($p < 0.001$; adjusted $R^2 > 0.99$) for both uncured and cured resins. The change and rate of refractive index and magnitude of shrinkage strain was dependent upon monomer formulation. The development of this interferometry technique provides a powerful non-invasive tool that will be useful for improving light transport through photoactive resins and filled resin composites by precise control of optical properties through material bulk.

^{§§} Hadis MA, Tomlins PH, Shortall AC and Palin WM. Dynamic monitoring of refractive index change through photoactive resins. *Dental Materials*, 2010; 26: 1106-1112

3.2.2 Introduction

The worldwide use of dental amalgam continues to decrease with several (Scandinavian) governments already prohibiting its application. The nearest chair side alternative in terms of mechanical and physical properties adequate for long-lasting, large stress-bearing restorations are resin-based composite (RBC) restorative materials polymerized by high intensity visible (blue) light. Unfortunately, there exist several shortcomings of RBCs in terms of technique sensitivity (partly a consequence of inadequate curing depths and multiple incremental placement steps) and dimensional change, which are partly a consequence of an incomplete understanding of the complex optical and physical material properties throughout cure.

Besides diffusion limited termination reactions (by free radicals trapped in the rapidly vitrifying matrix), inefficient light transmission and limited curing depths (generally less than 3 mm) are a result of surface reflection (Watts and Cash, 1994), photoinitiator and dye/pigment absorption (Ogunyinka et al., 2007), scattering by filler particles (Enami et al., 2005, dos Santos et al., 2008) and interfacial filler/resin refraction (Shortall et al., 2008, Feng et al., 2010). Consequently RBC materials are typically placed incrementally which may be time consuming and costly. Chapter 2 demonstrated the potential for reduced curing time with the photoinitiator, Lucirin TPO although inferior cure depths when compared to camphoroquinone based materials may be a cause for concern.

An interesting approach to reduce procedure time may be to increase cure depths of RBC materials. Whilst photoinitiator type and concentration may have some significance to improving light transmission (Section 3.1) through depth, other RBC components such as the resin matrix, filler, pigments and dyes may absorb, scatter and reflect light, as mentioned above. Although absorption characteristics of thin films are well characterised, the change in optical properties through the bulk of curing photopolymers (~1 cm) are not fully understood due to issues associated with light attenuation with increasing sample depth and inherently non-uniform and complex curing rate profiles (Stephenson et al, 2005; Kenning et al., 2006; Kenning et al., 2008). As the resin matrix (typically, co-monomer mixtures of Bis-GMA and TEGDMA) cures, its

optical properties change and refractive index rises due to a rapid increase in cross-link density and viscosity. As the refractive index of the resin approaches that of the filler, interfacial filler/resin scattering is reduced and light transmission is increased (Shortall et al., 2008). The refractive index (n) of a medium at optical wavelengths is defined as the ratio of the phase velocity of light in a vacuum (c) to the phase velocity, (v_p), in a given medium (Equation 3.2.1; Shoemaker et al., 1989).

$$n = \frac{c}{v_p} \quad \text{Equation 3.2.1}$$

In most materials the velocity with which light propagates is dependent upon the optical wavelength λ . Hence, when broadband, rather than monochromatic, light illuminates a sample, the speed with which it propagates will have an associated distribution, the peak of which is known as the group velocity v_g , with an associated group index n_g (Equation 3.2.2).

$$n_g = \frac{c}{v_g} \quad \text{Equation 3.2.2}$$

As light passes through the interface of one medium to another, the change in velocity is also accompanied by an optical path diversion. The relationship between the phase velocity in two media (V_A and V_B) the angle of incidence (θ_A) and refraction (θ_B) and the refractive index of the two media (n_A and n_B) can be predicted by Snell's law (Equation 3.2.3).

$$\frac{V_A}{V_B} = \frac{\sin \theta_A}{\sin \theta_B} = \frac{n_B}{n_A} \quad \text{Equation 3.2.3}$$

Therefore, by measuring the angle of refraction and knowing the index of refraction of the layer that is in contact with the sample, it is possible to determine the refractive index of an unknown sample accurately. The most commonly employed technique to measure refractive index uses a commercial Abbé refractometer. Typically,

the wavelength and temperature dependent refractive indices are measured at the standard yellow Fraunhofer line of sodium (589 nm) at 20 °C (nD₂₀). The measurement is based on observing the boundary line of the total reflection. Although critical angle refractometry techniques are well established for measuring the refractive index of homogenous transparent materials, these techniques only measure the index at a smooth, planar surface of a material. Having the capability to monitor the change in refractive index throughout the bulk of a curing of photoactive resin will allow for a greater insight into the effects of material properties on the complex setting reaction.

The use of low coherence interferometry (LCI) as a measurement method, has allowed the development of several techniques to obtain thickness and topography of multiple materials with excellent results and high accuracy (Tomlins et al., 2007; Morel and Torga, 2009). LCI is an optical technique that relies on a coherence grating to provide precise axial positioning of an object in the direction of light propagation and may be used for optical and physical characterisation of materials with accuracy in the micron range. LCI relies upon a broadband optical source, typically with a bandwidth of 50-100 nm, to determine the optical time of flight τ between material interfaces separated by a geometrical distance d (Equation 3.2.4).

$$\tau = \frac{d}{v_g} \quad \text{Equation 3.2.4}$$

The related group thickness t_g is defined (Equation 3.2.5)

$$t_g = n_g d \quad \text{Equation 3.2.5}$$

and the group index is defined from the spectrally dependent phase refractive index $n(\lambda)$ (Equation 3.2.6)

$$n_g = n(\lambda) - \lambda \frac{dn}{d\lambda} \quad \text{Equation 3.2.6}$$

From Equation 3.2.5 it is clear that LCI measurements can be used to determine sample group index, providing that geometric thickness is known a priori. Simultaneous determination of both n_g and d by LCI has previously been described (Sorin and Gray, 1992). However, inversion of Equation 3.2.6 to determine n from measurements of n_g is non-trivial (Rogers and Hopler, 1988). An alternative approach is to apply Equation 3.2.6 directly to phase refractive index measurements made at multiple discrete wavelengths using an Abbe refractometer. In this way LCI group index measurements may be compared with those made using critical angle refractometry.

This study aims to demonstrate in situ simultaneous measurements of refractive index and thickness change throughout polymerisation of model photoactive resin formulations using low coherence interferometry and to determine its reliability compared with a conventional Abbé refractometer technique.

3.2.3 Materials and Methods

Each experimental resin contained the same photoinitiator system, comprising 0.3% mass camphoroquinone and 0.6% mass amine co-initiator, dimethylaminoethyl methacrylate. Resins were prepared containing a mixture of bisphenol-A diglycidyl ether dimethacrylate (bis-GMA) and triethylene glycol dimethacrylate (TEGDMA) at varying mass ratios to control viscosity and refractive index (Table 3.2.1). All reagents (Sigma Aldrich, Gillingham, UK) were used as received.

Resin	TEGDMA (mass %)	Bis-GMA (mass %)
M1	15	85
M2	30	70
M3	59	41
M4	90	10

Table 3.2.1: Monomer composition of the experimental resins.

Refractive index (RI) and physical thickness measurements of each model resin formulation were obtained simultaneously throughout cure. The bespoke configuration comprises of a Michelson interferometer arrangement (Figure 3.2.1) with a super luminescent diode (L2KASLD16-015-BFA; Amonics, Kowloon, Hong Kong) broadband light source with a centre wavelength at 1624 nm and a full-width half maximum of 55 nm. An ‘off the shelf’ fibre coupled InGaAs photo detector (PDB120C; Thorlabs, Cambridgeshire, UK) was used to detect the interference fringes from the sample and the reference arm. A voice coil stage (VCS10-023-BS-1, H2W Technologies, Inc., USA) was used to change the position of the reference mirror in order to modulate an optical time delay between the reference arm and the sample arm.

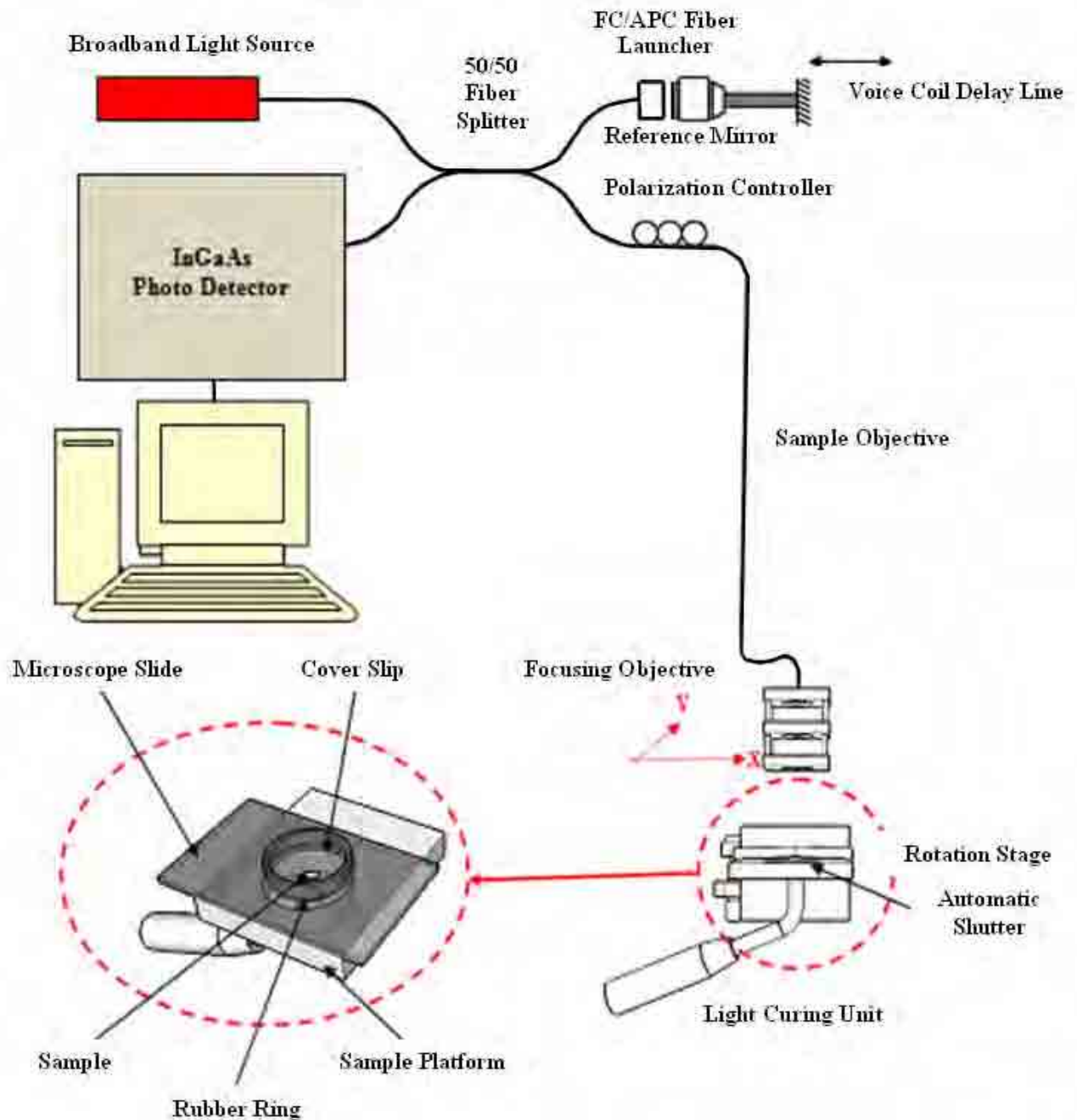


Figure 3.2.1: Experimental arrangement of the frequency domain low coherence interferometry system used for dynamic monitoring of refractive index and shrinkage of photoactive materials.

Nominally identical ($n=3$) cylindrical specimens were prepared using a white Teflon ring mould (12 mm diameter 1.4 mm thick). Resins were carefully pipetted into the mould cavity ensuring no inclusion of air bubbles and slightly overfilled. Opposing surfaces were covered with a glass microscope slide (1.10 ± 0.2 mm) on the irradiated surface and glass cover slip (0.15 ± 0.03 mm) on the non-irradiated surface to contain the resin within the mould and limit oxygen inhibition of the outer layers. Both slides and slips had an $RI = 1.518 \pm 0.001$ at 23°C . A black nylon ring mould was used to contain each specimen, minimise reflection and allow reproducible concentric alignment with the sample objective and the light curing unit (Figure 3.2.1). The light curing tip was held 5 mm away from the surface of the glass microscope slide in order to facilitate the use of a automatic shutter (aperture = 25 mm) operating at a shutter speed of $1/60$ s.

Each specimen was light irradiated for 90 s (after activation of the curing lamp for 10 s without irradiation of the sample controlled by the shutter) with a halogen light curing unit (XL2500, 3M ESPE, Dental Products, Seefeld, Germany) with a 12 mm diameter fibre-optic bundle at an ambient temperature of 23 ± 1 °C. Preliminary investigations identified the requirement for a 10 s non-irradiating lamp activation as fluctuations in light transmission were observed over this initial period. An irradiance of 733 ± 18 mW/cm² was measured using a handheld digital radiometer (Cotolux light meter, Cuyahoga Falls, OH, USA) held 5 mm normal to the light curing tip. Following cure, the samples were removed from the mould and the resultant flash was cut away using a sharp blade.

Spectral data was acquired continuously over a 110 s period, which included a 90 s cure duration. Averages ($n=3$) were calculated over the 90 s cure period and following 7 days post irradiation. A bespoke computer programme was used to automate the processing of the spectral data to obtain refractive index and shrinkage strain information in real-time.

A conventional Abbé refractometer (Abbe 60, Bellingham and Stanley Ltd, UK) coupled to an infra red camera (Alpha NIR, Indigo Systems, USA) was used to acquire static refractive indices at three wavelengths of 1500 1550 and 1620 nm for both uncured and cured phases of each resin. The sample was illuminated with a narrow-band tuneable laser (Tunics Plus CL/WB, Anritsu, Japan) at a power of less than 1 mW, spread over an

area of approximately 10x10 mm. The group index for each sample was determined by using Equation 6. The derivative of n with respect to λ was estimated from the slope of a linear fit through the phase refractive index measurements. This approach was shown to be valid by first measuring a pure silica calibration artefact (PT08002, Bellingham and Stanley, UK). Each specimen was prepared in a bar-shaped rectangular mould having similar dimensions to the calibration piece (25 x 10 x 2 mm). An alignment mould was used to ensure repeatable placement of the light-curing unit with respect to the specimen. Each portion of the bar-shaped specimen was cured for 90 s to establish a full and even cure across the sample length. Specimens were measured following cure and 7 days post irradiation.

As Kolmogorov-Smirnov normality tests on pooled refractive index data sets indicated variation from the pattern expected from a population with a normal distribution, data analysis was performed using distribution free test methods. Kruskal-Wallis and Mann-Whitney non-parametric tests were performed with a significance level of $p=0.05$. Where appropriate, parametric data was analysed using one-way analysis of variance (ANOVA) and Tukey post-hoc comparisons (SPSS v.15.0, Surrey, UK) at a confidence level of 95%.

3.2.4 Results

The effect of measurement method on the RI of each resin mixture prior to and following 7d after cure are presented in Table 3.2.2. Kruskal-Wallis tests revealed significant differences ($p<0.001$) in RI values obtained by LCI and phase index results measured with the Abbé refractometer, although RI measured by LCI and group indices were not significant ($p=0.261$). The phase index results were higher than those measured by LCI ($p<0.001$), however, a significant linear regression was observed between LCI and Abbé (phase) and LCI and Abbé (group) (Figure 3.2.2; adjusted $R^2>0.99$ for both regressions).

	LCI		Abbe (group index)		Abbe (phase index)	
Monomer	uncured	cured	uncured	cured	uncured	cured
M1 (15/85)	1.5336 (0.00153)	1.5598 (0.00432)	1.5373 (0.00027)	1.5630 (0.00016)	1.5132 (0.00027)	1.5387 (0.00016)
M2 (30/70)	1.5240 (0.00517)	1.5558 (0.00407)	1.5260 (0.00042)	1.5567 (0.00014)	1.5020 (0.00042)	1.5325 (0.00014)
M3 (59/41)	1.4987 (0.00613)	1.5490 (0.00570)	1.5008 (0.00015)	1.5407 (0.00014)	1.4775 (0.00015)	1.5168 (0.00014)
M4 (90/10)	1.4743 (0.00521)	1.5384 (0.00986)	1.4756 (0.00016)	1.5241 (0.00010)	1.4528 (0.00016)	1.5004 (0.00010)

Table 3.2.2: The mean refractive index for resins with various TEGDMA/Bis-GMA ratios prior to and following 7 days after initial polymerization using two methods; low coherence interferometry (LCI) and a conventional Abbé technique with group and phase index values. Standard deviations are displayed in parentheses.

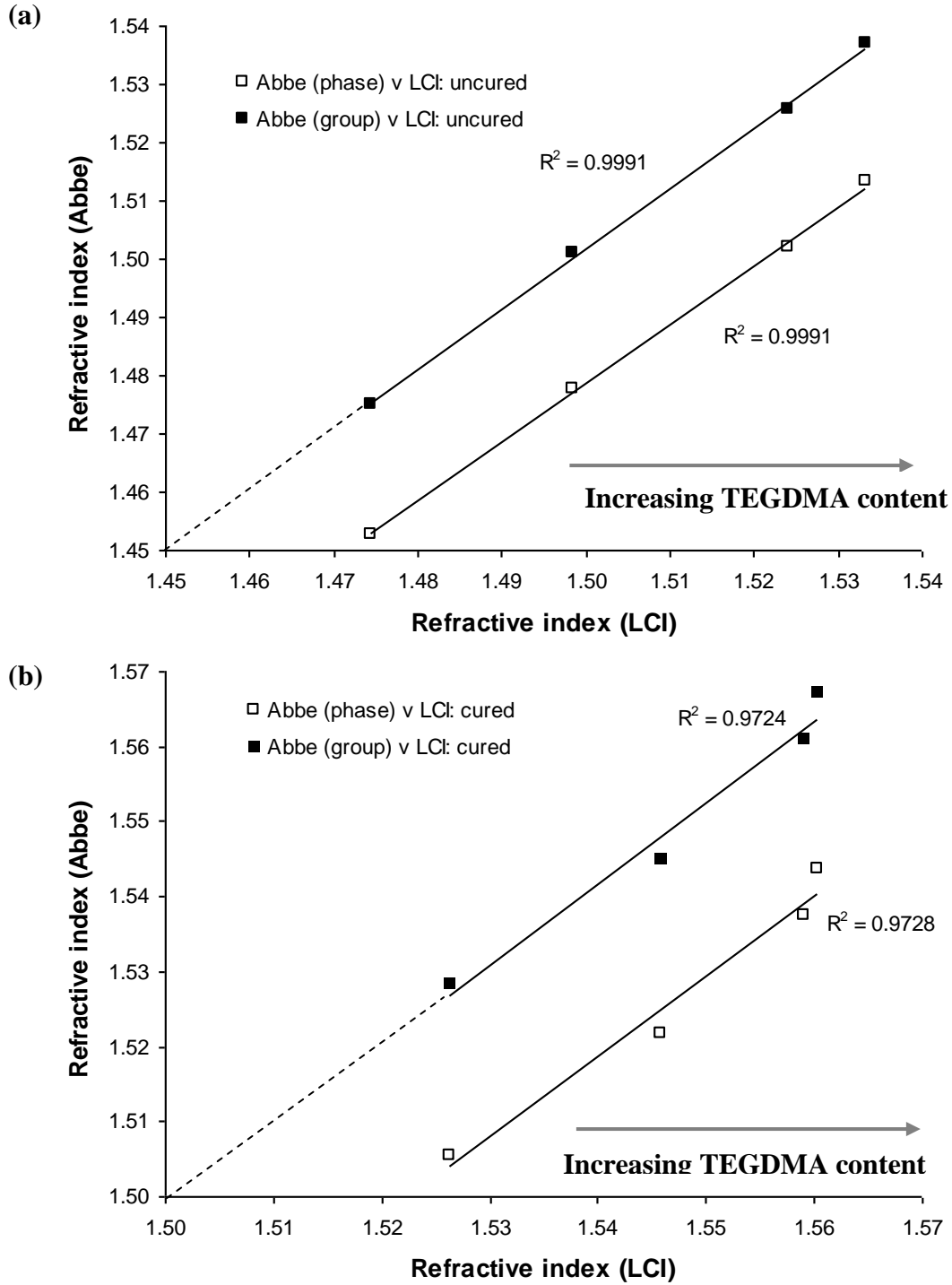
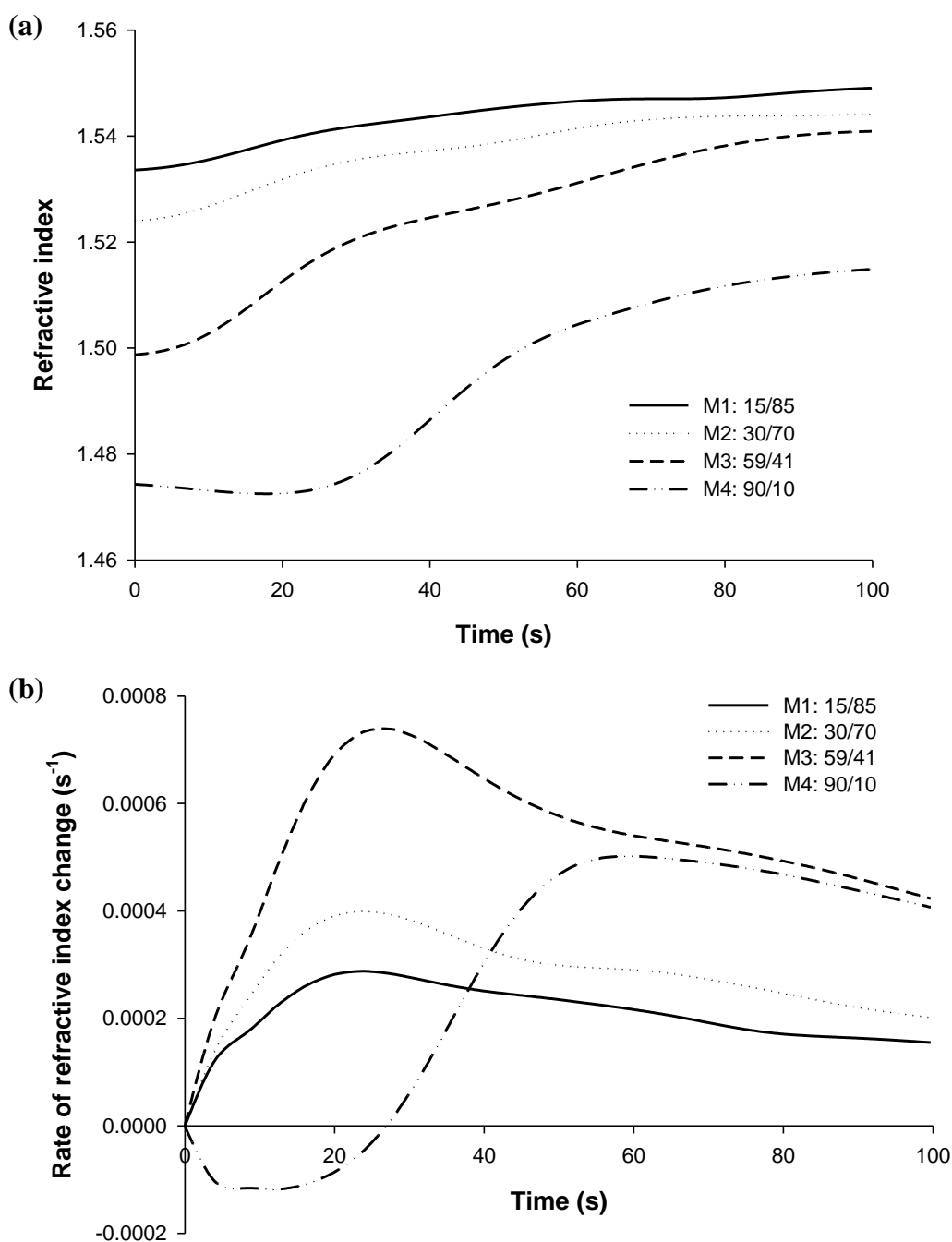


Figure 3.2.2: Correlation of the low coherence interferometry and conventional Abbé techniques for measuring phase and group refractive index of resins with increasing TEGDMA content (a) before and (b) following cure.

Regardless of measurement technique, RI significantly decreased as TEGDMA concentration increased for both pre-cure and post-cure measurements ($p < 0.05$) although the change in RI was greater for TEGDMA-rich resins following polymerization ($p < 0.001$; Table 3.2.2). For LCI measurements a gradual increase in RI of each resin mixture was observed throughout cure (Figure 3.2.3a) albeit at different rates (Figure 3.2.3b).



(c) *continued...*

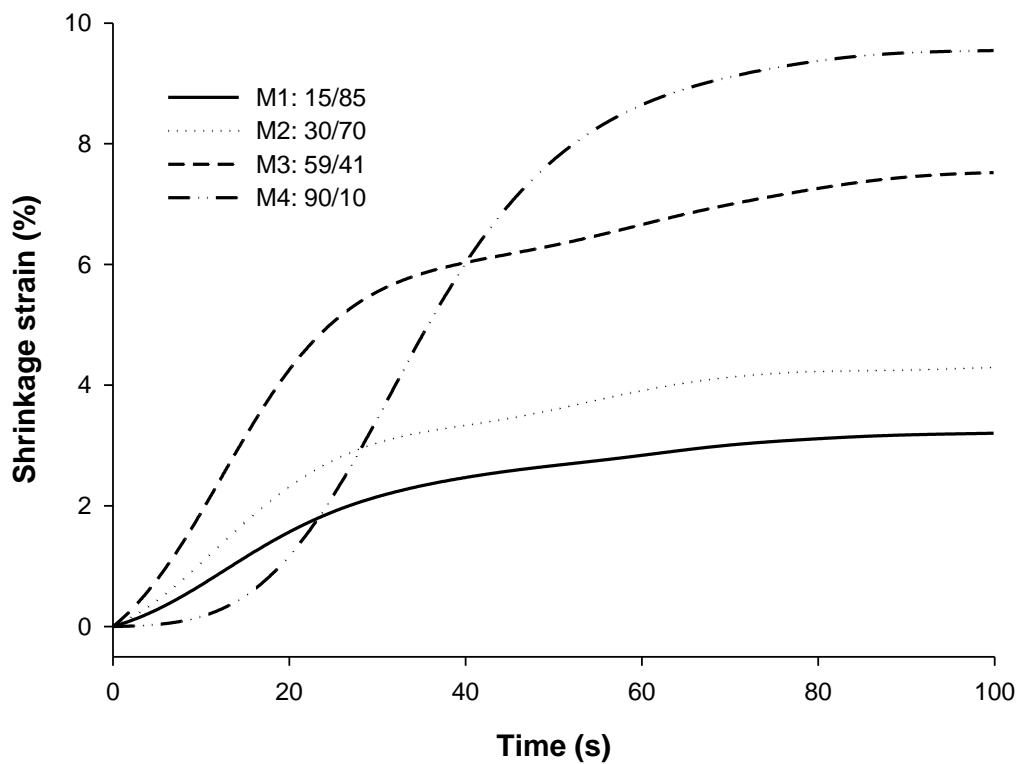


Figure 3.2.3: The dynamic and simultaneous measurement of (a) refractive index, (b) rate of refractive index change and (c) shrinkage strain of curing resins with various TEGDMA/Bis-GMA ratios.

The magnitude of shrinkage strain following cure was dependent upon monomer composition and one-way ANOVA and post-hoc comparisons revealed significant increases in shrinkage strain for increasing TEGDMA content where the shrinkage strain of $M4 > M3 > M2 > M1$ ($p < 0.05$; Figure 3.2.3c).

For M1, M2 and M3 resins at initial stages of polymerization, shrinkage strain increased proportionally with refractive index, however, for M4, RI appeared to decrease over the first 25 s of light irradiation (Figure 3.2.3b and Figure 3.2.4). A deviation from linearity occurred at different stages of cure depending on TEGDMA content, where development of shrinkage strain lagged behind the increase in RI.

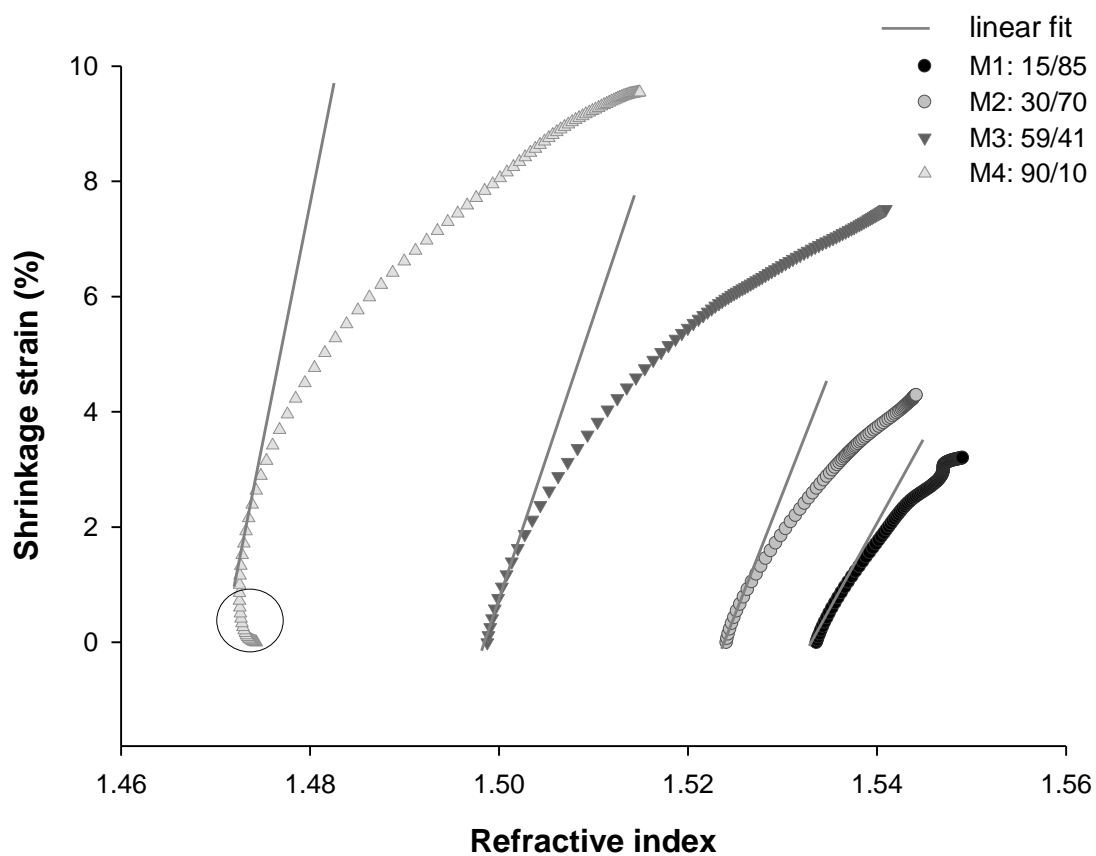


Figure 3.2.4: The relationship between increasing shrinkage strain and refractive index of curing resins with various TEGDMA/Bis-GMA ratios. The circled region highlights an inverse trend where refractive index appears to decrease as the material cures.

3.2.5 Discussion

The reliability of the LCI technique was tested against a conventional procedure using an Abbé refractometer and good agreement was established between the two methods. However, despite a >99% correlation, group refractive index values were found to be higher than phase refractive indices by approximately 2×10^{-2} (Figure 3.2.2). The relationship between group and phase refractive index is summarized by Equation 3.2.6, from which it is expected that the group index will be greater than the phase index when the phase index decreases as a function of increasing wavelength. This regime is consistent with normal dispersion, expected of the polymer matrix at near-infrared wavelengths. Calculated group refractive indices from phase refractive index measurements using Equation 3.2.6 show agreement with the LCI measurements, confirming the validity of this measurement modality.

Conventional refractometry is based on the subjective observation of a boundary line which may cause anomalies and specimen geometry is confined to the dimensions of the refractometer to minimise scattering. An alternative method to calculate RI of an amorphous polymer is based on the molar properties of a single polymer structural unit and for every bond that contributes to the structural unit, bond refraction can be defined (Lehtinen et al., 2008). However this method is only an approximation of fully polymerised materials and is restricted to tabulated values of bond refraction of the constituent bonds. With the technique described here, there is no sample restriction and appropriate specimen shape and thickness can easily represent those used in light-cured dental materials and other photoactive applications.

The decrease in refractive index with increasing TEGDMA content (Table 3.2.2) is a result of lower viscosity as molecular weight and density of the resin mixtures decreases (Shortall et al., 2008). As photo polymerization precedes a rise in RI (Figure 3.2.3a) is caused by decrease in molecular polarizability as monomer molecules connect to the growing polymer chain. The change in polarizability during cure is a result of double to single bond conversion and an increase in density due to volumetric shrinkage. Irradiation initiates an exothermic reaction characterized by gelation (liquid to rubber) and vitrification (rubber to glass transition) and the RI increases as the resin is converted

from a mobile system to a glassy solid, where initial mobility and viscosity are proportional to the TEGDMA concentration. During polymerisation, network formation and therefore the rate of RI change is significantly affected by the TEGDMA/Bis-GMA ratio (Figure 3.2.3b). The high mobility of TEGDMA rich resins dominate the reaction, resulting in a delay to reach the maximum change in RI and shrinkage strain. The low viscosity of TEGDMA increases the mobility of the propagating species resulting in a decreased local radical concentration which delays the onset of vitrification compared with more viscous mixtures. In resins containing lower concentrations of TEGDMA, higher viscosity results in the more rapid onset of auto acceleration as mobility is rapidly restricted and polymerization rate and termination become diffusion controlled (Lovell et al., 1999). Despite the lower initial rate of RI increase in the 90/10 (M4) resin, the extent of shrinkage strain following cure is increased (Figure 3.2.3c) due to the higher concentration of double bonds in TEGDMA compared with Bis-GMA (~7 and 4 mol/kg, respectively), which increases monomer to polymer conversion.

The Clausius-Mossotti relation or Lorentz-Lorenz equation defines the relationship between refractive index and molecular polarization where molar refraction changes proportionally with polarizability (Patel et al., 1992). Since a relative comparison exists between polarizability (and degree of polymer conversion) and density, a linear relationship was expected between refractive index and shrinkage (Kopietz et al., 1984). Indeed, throughout initial stages of polymerization there exists a linear relationship between the change in shrinkage strain and RI for each resin (Figure 3.2.4). It is well known that as the polymerization reaction reaches its maximum rate, the curing system vitrifies and shrinkage kinetics fall below that of conversion (Stansbury et al., 2005). Consequently, a deviation from the initial linear relationship of refractive index change and shrinkage strain is observed (Figure 3.2.4). This correlation is further complicated by the magnitude of shrinkage strain exhibited by each resin mixture. The resins containing higher Bis-GMA ratios exhibit less volumetric shrinkage and the departure from linearity is less pronounced compared with TEGDMA rich resins. The latter mixtures also exhibit an inverse trend where refractive index appears to decrease as the material cures. Since initial reaction rates are low in mobile curing systems, heating effects from the curing lamp may decrease viscosity, which outweighs any small decrease in volume at such

slow rates of polymerization resulting in the apparent decrease in RI at early stages of cure.

Specimen geometry is likely to significantly affect the shrinkage vectors of the curing specimen. A low aspect ratio (where the specimen thickness approaches its measured diameter) would be expected to cause multi-axial, rather than linear shrinkage (Watts and Marouf, 2000), which would significantly affect both the measured optical and physical thickness change of the curing specimen. Here, a large aspect ratio was chosen to confine shrinkage to one direction and allow reliable correlations with refractive index and thickness change. The effect of specimen geometry on dynamic optical and physical properties will be investigated in Section 3.3 of this thesis.

3.2.6 Conclusion

Filler particle refractive indices vary between filler type but remain constant throughout polymerization within a resin matrix (Shortall et al., 2008). The ability to measure refractive index change of curing photoactive resins will allow for the design of composite formulations, which exhibit optimum evolution of translucency (as the RI of the curing resin approaches that of the filler) and enable precise control of aesthetic quality, light transport through material bulk and curing depth. Furthermore, since polymer conversion relates directly to polarizability and refractive index, development of this technique (and miniaturisation of the components to a handheld device) will allow for in-situ monitoring of the curing extent of resin-based materials and therefore eliminate the risk of under cured restorations in clinical practice.

3.3 Specimen Aspect Ratio and Light-Transmission in Photoactive Resins

3.3.1 Abstract

This section tests the influence of specimen dimensions on light transmission and shrinkage strain properties of curing dental resins. Experimental photocurable resins (Bis-GMA/TEGDMA) were placed in cylindrical cavities in disc-shaped moulds with aspect ratios (AR) ranging from 2-24; 4x2mm (AR: 2); 4x1mm (AR: 4); 8x2mm (AR: 4); 8x1mm (AR: 8); 12x1mm (AR: 12); and 12x0.5mm (AR: 24) were cured using a quartz tungsten halogen light curing unit. Light transmission and shrinkage-strain data were recorded throughout curing for 70 s. Surface hardness measurements were taken to indirectly assess the extent of cure on the upper and lower surface of each specimen. One-way ANOVAs and post-hoc Holm-Sidak revealed significant effects of specimen dimensions on light transmission and Vickers hardness number ($p < 0.05$). Cavity height significantly affected initial and final light transmission in nearly all cases except for the final light transmission in specimens with 12 mm diameter moulds ($p < 0.05$). The lowest AR specimens showed an increase in transmission above 100%, which may be due to a possible 'lens' effect caused by specimen constraint. The extent of lower surface cure (as assessed by increasing micro-hardness) was principally affected by cavity height. Micro-hardness was significantly lower for thicker specimens except for specimens with 8 mm diameter ($p < 0.05$). At constant thickness, only the 2 mm thick specimens showed significantly higher lower surface cure with increasing diameter ($p < 0.05$). The highest AR specimen showed the significantly highest lower surface hardness percentage, which represents almost equivalent cure on either side ($p < 0.05$). Total shrinkage strain increased with AR but no noticeable trends were observed at constant diameter and varying height. Shrinkage strain per unit mass was significantly lower in most cases with increasing diameter at constant height ($p < 0.05$) except when diameter increased from 8 mm to 12 mm. Specimen constraint during polymerization in low AR cavities may compromise light transmission as unexpected light intensity variations may occur for low configuration factors. Improving light transmission through depth will ultimately improve conversion and clinical longevity of the restoration.

3.3.2 Introduction

Photoactive dental RBCs are now widely used in restorative dentistry due to their aesthetic properties and the choice of this material as an alternative for amalgam (Burke, 2004). However, polymerisation shrinkage of RBCs is a major concern amongst both clinicians and researchers. Polymerisation induced shrinkage stress can lead to failure of the interfacial bonds and can result in marginal leakage, premature failure of the restoration, and in some cases micro, or even macro-cracking of the surrounding tooth (Hübsch et al., 1999; Palin et al., 2005).

The shrinkage-strain induced during photo-polymerisation may be characterised by its magnitude, direction (or vector character of strain) and time dependence. The strain imposed during photo-polymerisation is not only influenced by the material factors such as composition (Anseth et al., 1996; Section 3.2) but also shape and the size of the cavity being restored as well as the compliance of the surrounding cavity (Feilzer et al., 1987; Watts and Marouf, 2000). When the material undergoing shrinkage is bonded by opposing cavity walls, strain becomes restricted and the final stress generated is given by the product of the elastic modulus and the strain vectors (Davidson and Feilzer, 1997; Sakaguchi and Ferracane, 1998). Not only is the stress magnitude determined by characteristics of composites such as filler content (Condon and Ferracane, 2000), curing rate (Bouschilcher and Rueggeberg, 2000) and degree of conversion (Braga and Ferracane, 2002; Lu et al., 2004) but many authors have also reported that the final stress is governed by the configuration factor (*C*-factor) of the cavity; the ratio of bonded to un-bonded surfaces (Equation 3.3.1), which may also be given by the aspect ratio (diameter/height), and the compliance of the system (Feilzer et al., 1987; Loguercio et al., 2004; Braga et al., 2006; Watts and Satterthwaite, 2008). In the equation, *r* is the cavity radius; *h* is the cavity height and ϕ is the cavity diameter.

$$C_f = \frac{2\pi r^2}{2\pi r h} = \frac{r}{h} = \frac{\phi}{2h} \quad \text{Equation 3.3.1}$$

Previous investigations have suggested that *C*-factors less than 1 result in the slow

development of shrinkage stress and the composite remains bonded to the cavity walls (Davidson and de Gee, 1984; Feilzer et al., 1987). The *C*-factor itself is governed by the geometry of the cavity, i.e. the diameter and the height which both independently influence the generation of shrinkage-stress. A recent investigation by Watts et al (2008) reported that at a constant height, with diameter variation, a *C*-factor increase from 0.6 to 6 gave an exact exponential decrease in shrinkage-stress, whereas at constant diameter with height variation, an increase in *C*-factor from 3-100 gave an increase in shrinkage-stress (Watts and Satterthwaite, 2008). It is clear that the stress generated during polymerisation cannot simply be extrapolated to *C*-factor and is dependent upon several other properties, amongst which volume and mass as well as the ratio of bonded to non-bonded surfaces will influence the final stress as-well as bonding or lack of such to mould or tooth and matrix surfaces and compliance of such (Watts and Marouf, 2000; Braga et al., 2006; Watts and Satterthwaite, 2008).

The transmission of light and number of photons absorbed by the photoinitiator is critically important to the extent of polymerisation and is a desired parameter to fully characterize since it ultimately relates to the final mechanical and physical properties of the polymer. Inefficient light transmission and limited cure depths are associated with surface reflection (Harrington and Wilson, 1995), absorption (Chen et al., 2007; Ogunyinka et al., 2007) and interfacial scattering (Shortall et al., 2008; Feng et al., 2010). Upon polymerisation, photoinitiator decomposition will reduce the concentration of light absorbers (Section 3.1) and polymerisation shrinkage (Section 3.2) will reduce the optical path length, which will increase the transmitted light through the sample according to Beer-Lamberts law (Equation 3.3.2).

$$A = \varepsilon Cl = -\log\left(\frac{I}{I_0}\right) \quad \text{Equation 3.3.2}$$

where *A* is the absorbance, *l* is the optical path length, *C* and ε are the concentration and molar extinction coefficient of the absorbing species (i.e the photoinitiator) and *I* and *I*₀ are the light transmission and initial light transmission, respectively.

Varying patterns of light transmission through filled and unfilled photocurable resins have been reported. Harrington et al. (1996) demonstrated a method, which allowed the actual consumption of light energy by the photoinitiator to be calculated and showed increasing light intensity with radiation time for commercial composites (Harrington et al., 1996). Using the same technique, Ogunyinka et al. (2007) reported decreasing transmission during polymerisation of experimental composite formulations. Shortall et al. (2008) showed the effect of monomer and filler composition on light transmission and reported that the refractive index mis-match between filler and resin will affect its light transmission profile where maximum transmission is observed when the refractive index of the resin and filler are equal. Section 3.2 also demonstrated the effect of monomer composition on the variation in optical properties and physical thickness change in unfilled resins during irradiation. Although shrinkage strain was dependent upon material composition, cavity constraint is likely to affect shrinkage strain development, which in turn may affect light transmission. Consequently, this Section tests the following hypotheses:

- i) For high specimen aspect ratio, increased shrinkage strain will increase final light transmission in curing photoactive resins.
- ii) If a higher final light transmission is observed, then curing extent will also improve.

3.3.3 Materials and Methods

Experimental resins containing a 50:50 mix of Bisphenol-A diglycidyl ether dimethacrylate (Bis-GMA) and triethylene glycol dimethacrylate (TEGDMA) were mixed. The photoinitiator system within each experimental resin contained the photoinitiator camphoroquinone (0.2 wt%) and the co-initiator dimethylaminoethyl methacrylate (DMAEMA; 0.8 wt%). A resin that contained no photoinitiator system was employed as a negative control for each of the moulds tested. All reagents were provided by Sigma Aldrich, UK and were used as received. Resins were light cured using a quartz tungsten halogen light curing unit (XL2500, 3M ESPE, Dental Products, Seefeld, Germany) in cylindrical cavities placed in black nylon moulds having C-factors (C_f ; Equation 1) ranging from 1-12 and aspect ratios (AR) ranging from 2-24; 4x2 mm (C_f : 1, AR: 2); 4x1 mm (C_f : 2, AR: 4); 8x2 mm (C_f : 2, AR: 4); 8x1 mm (C_f : 4, AR: 8); 12x1 mm (C_f : 6, AR: 12); and 12x0.5 mm (C_f : 12, AR: 24; Table 3.3.1). Aspect ratio was always twice that of C-factor (Table 3.3.1) and as they both represent the ratio of bonded to non-bonded surfaces, from herein only aspect ratio will be mentioned. Each specimen (n=3) was light irradiated for 60 s (after activation of the curing lamp for 10 s without irradiation of the sample) with a 12 mm diameter fibre-optic bundle at an ambient temperature of 23 ± 1 °C. A 10 s non-irradiating lamp activation period was required as fluctuations in light transmission occurred during this initial period. An irradiance of 733 ± 18 mW/cm⁻² was measured using a handheld digital radiometer (Coltolux light meter, Cuyahoga Falls, OH, USA) held 5 mm normal to the light curing tip.

Mould Dimensions ($\phi \times h$)	Volume (mm ³)	Mass (mg)	Aspect Ratio (ϕ / h)	C-factor
4x2	25.1	28.1	2	1
4x1	12.6	14.1	4	2
8x2	100.5	112.4	4	2
8x1	50.3	56.2	8	4
12x1	113.1	126.5	12	6
12x0.5	56.5	63.2	24	12

Table 3.3.1: Mould Properties

Resins were pipetted into the moulds with opposing surface covered with 1.10 ± 0.2 mm thick glass microscope slides (RI = 1.518, 23°C) to limit oxygen inhibition of the outer layers. The microscope slides were clamped into place with bulldog clips to form an air-tight arrangement before being placed onto the sample holder, in line with the optical fibre (8 µm diameter) and cosine corrector (CC3-UV; 3.9 mm diameter) and the light curing unit which was held 5 mm away from the irradiated surface. An optical rail ensured concentric alignment and rigidity throughout the experiment (Figure 3.3.1). Each specimen received 60 s irradiation controlled by an automated shutter (aperture 25 mm diameter) operating at a shutter speed of 1/60 s. Throughout irradiation light transmission was measured in real time using a USB4000 spectrometer and Spectrasuite application software (Ocean Optics, Duiven, The Netherlands) scanning the average from 400 nm to 550 nm with an integration time of 0.1 s. The 10 s non-irradiation period was then excluded from the charts and the initial transmission was taken as the point where transmission started to increase steadily ($t=0$). Transmission traces were then corrected accordingly for the decrease in light transmission through a non-curing resin in each mould that was tested, i.e. the control group of each mould. Firstly the initial light transmission values (the point at which maximum transmission occurs in non-curing resins) of the control group (N_i) was subtracted from the real time transmission (N_t) value of the control group to obtain the increase in transmission in real time for each mould. The corresponding values were then added onto real time measurements (LT) to exclude the affect of light intensity decreasing due to heating of the bulb which enabled corrected light transmission values (CLT; Equation 3.3.3) to be obtained. The control group was then taken as 100 % throughout the cure duration and CLT % of each curing resin was calculated as a percentage of each control group.

$$CLT = (N_i - N_t) + LT \quad \text{Equation 3.3.3}$$

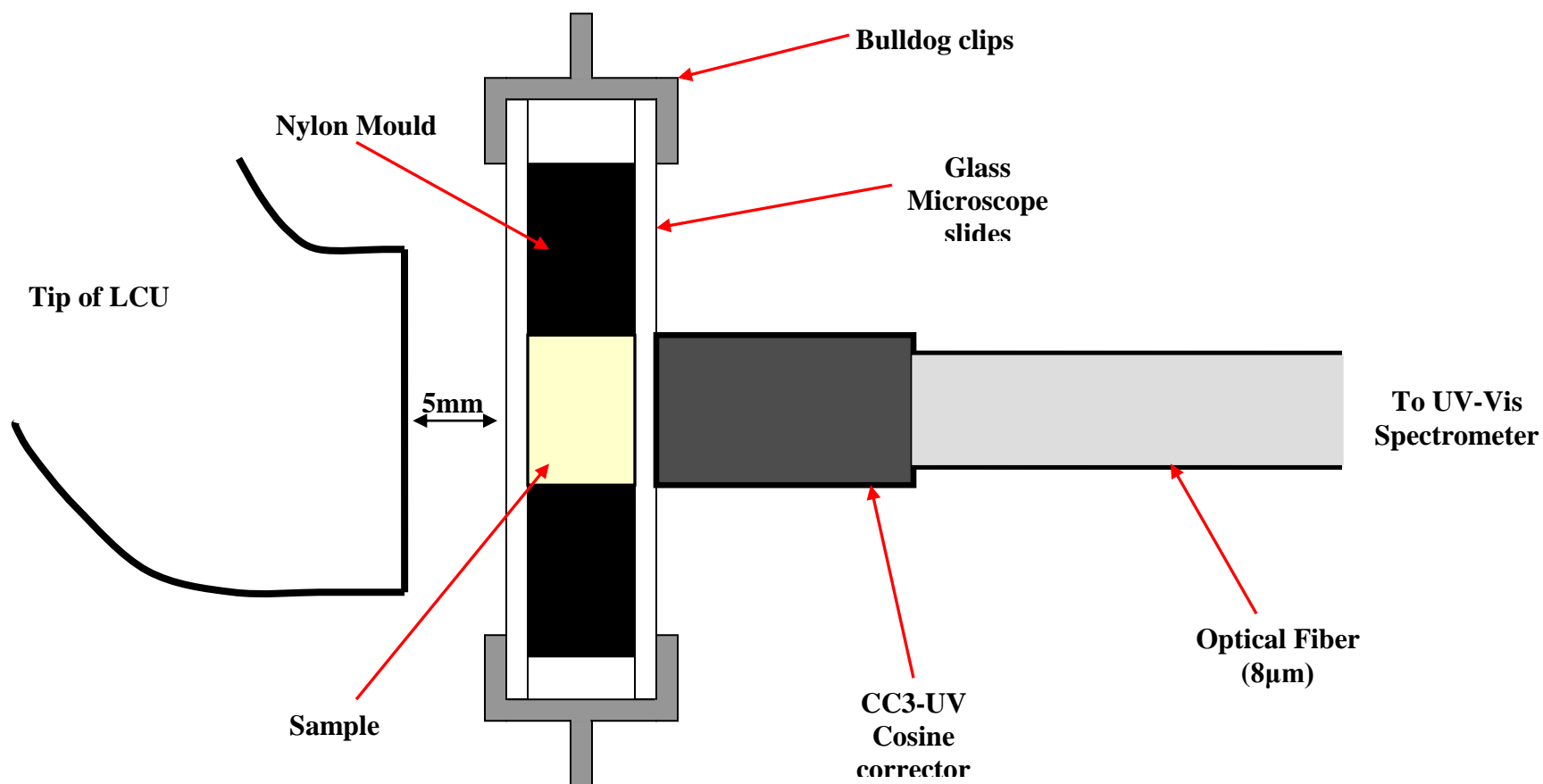


Figure 3.3.1: Experimental arrangement for the measurement of light transmission in real time of curing resins.

Shrinkage strain measurements were based on low coherence interferometry methods developed in Section 3.2. The bespoke configuration comprises of a Michelson interferometer arrangement, which has a super luminescent diode (L2KASLD16-015-BFA; Amonics, Kowloon, Hong Kong) broadband light source with a centre wavelength at 1624 nm which allows shrinkage-strain to be measured centrally in each sample (Figure 3.3.2). An optical time delay is modulated by changing the position of the reference arm with a voice coil stage (VCS10-023-BS-1, H2W Technologies, Inc., USA) which produces interference fringes. The ‘off the shelf’ fibre coupled InGaAs photodetector (PDB120C; Thorlabs, Cambridgeshire, UK) allows the detection of the interference fringes from the sample and the reference arm. Each specimen was polymerised using an identical curing protocol as that used for light transmission measurements. A bespoke computer programme was used to automate the processing of the spectral data to obtain shrinkage strain information in real-time. Since mass plays a dominant role (Watts and Satterthwaite, 2008), plots of strain variations per unit composite mass separated the effects which mould dimensions had on strain developed centrally at the specimen mid-point.

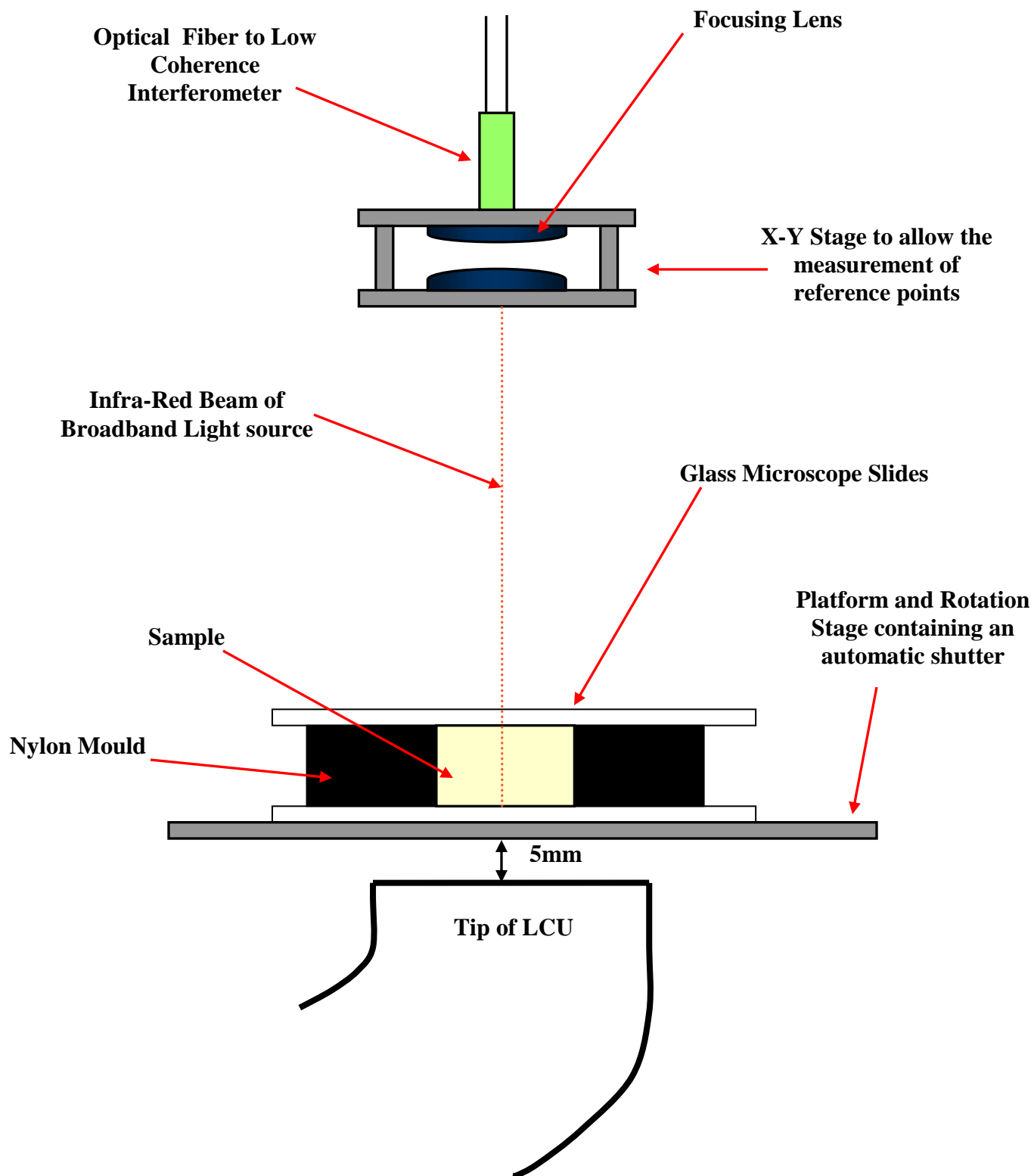


Figure 3.3.2: Experimental arrangement for the measurement of real time shrinkage-strain

Micro-hardness measurements were made on the directly irradiated surface and on the surface which was furthest away from the light curing unit using a Vickers Hardness tester (Struers, Glasgow, UK). Indents (load=1.96 N; dwell time= 10 s) were made (n=3) less than 1 mm away from the centre on each surface of each sample (n=3) after 24 h dark storage post irradiation. The averages were taken respectively and the mean was calculated as group. The percentage of lower surface hardness was then calculated with respect to the upper surface to give indications of lower surface cure.

In order to assess the effect of bonded versus non-bonded surfaces on light transmission, moulds were produced from pre-fabricated fine-particle feldspar ceramic blocks (Vitablocs Mark 2,, VITA, Germany) cut into 2 mm thick sections. The surfaces of each section were painted with black nail varnish (Rimmel, London) to minimize transmission and reflection of light through the moulds. Dental silicone impression material (Heraeus Kulzer, Germany) was used to hold the slabs centrally in an alignment ring. Cavities were prepared centrally (4 mm diameter) with a water cooled high speed dental hand piece which had a diamond bur attachment. The moulds were then split into two groups (n=3), a bonded group and a non-bonded group. For the first group, the cavity walls and glass slides were acid etched with hydrofluoric acid (5%) for 2 minutes and washed and dried several times before leaving in dry conditions at room temperature for 24 h. After 24 h, Scotchbond (3M ESPE, Adper Scotchbond) adhesive was applied to the cavity walls and glass slides before curing for 20 s from each side of the cavity with the light curing unit. For the second group, liquid separator based on hexane resin solution (DVA Very special separator, USA) was applied both to cavity walls and the glass slide before the resin was applied. Light transmission was recorded as described above.

Statistical analyses were performed on the initial and final light transmission values, the Vickers hardness number and strain percentage (or strain percentage per unit mass) data sets for the various moulds. Data sets were analysed using one-way analysis of variance (ANOVA) and post-hoc Holm-Sidak pairwise comparison tests in Sigmasat Version 3.5, where the effects of thickness at constant diameter and the effect of diameter at constant thickness were compared ($p < 0.05$).

3.3.4 Results

The one-way ANOVA's revealed a significant difference between the cavity dimensions for each of the responses. The initial light transmission was significantly affected by the cavity height ($p < 0.05$) where decreasing transmission was observed with increasing height (Table 3.3.2a and Figure 3.3.3). For increasing diameter at constant thickness only the 1 mm thick moulds showed a significant decrease in initial transmission ($p < 0.003$) for 4 to 8 mm and 12 mm wide specimens but not from 8 to 12 mm wide specimens (Table 3.3.2b and Figure 3.3.3). However, although some of the effects of diameter were insignificant, a noticeable decrease in initial light transmission was observed with increasing diameter and volume at constant height (Table 3.3.1; Table 3.3.2b and Figure 3.3.3). For final light transmission, thicker moulds at constant diameter showed significantly higher transmission ($p < 0.05$) except for the 12 mm diameter specimens where the effect of thickness was insignificant (Table 3.3.2a). At constant thickness, narrow specimens showed significantly higher final transmission although the change from 8x1 mm to 12x1 mm was not significant (Table 3.3.2b). Specimens cured in the thickest and narrowest mould (4x2 mm) showed the most significant increase in light transmission ($> 100\%$) during irradiation in comparison to the other specimen groups ($p < 0.05$; Table 3.3.2 and Figure 3.3.3).

(a)

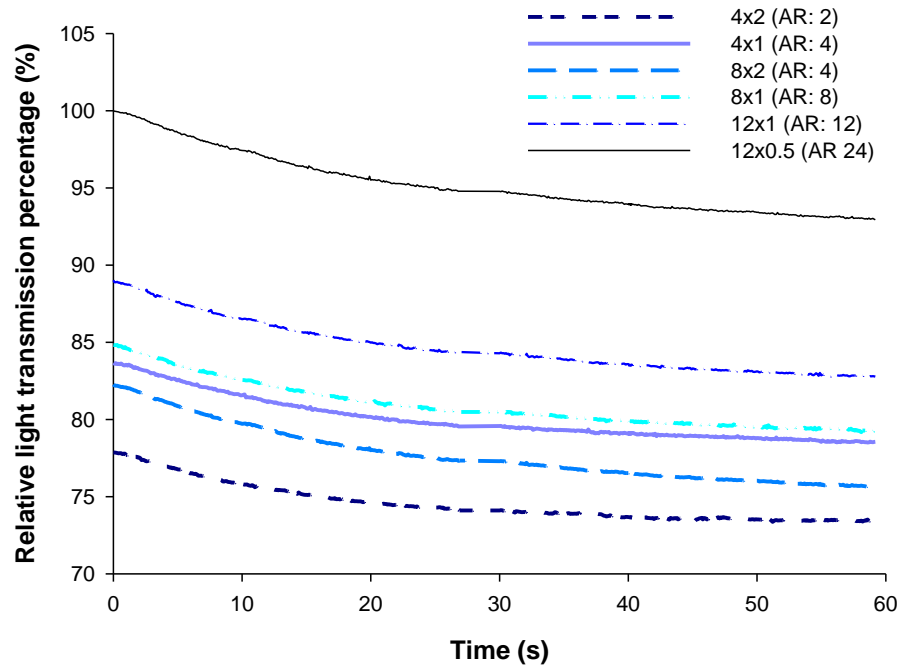
COMPARISON OF CHANGE IN THICKNESS					
Mould Dimensions (ϕ x h)	Initial relative light transmission percentage (t=0)	Final relative light transmission percentage (t=60)	Total strain % (S_e %)	Strain % Per unit mass (% S_o g ⁻¹)	Lower hardness percentage
4 x 1	93.8(0.5) a	101.5(0.5) b	4.9(0.4) b	0.346(0.031) a	93.2(1.3) b
4 x 2	89.5(1.6) c	107.7(2.4) a	4.1(0.7) b	0.145(0.025) b	85.6(1.1) c
8 x 1	91.3(0.1) b	99.8(0.2) c	5.8(0.8) a,b	0.103(0.013) c	92.1(2.0) b
8 x 2	87.9(0.6) c	101.4(0.9) b	6.4 (1.6) a,b	0.057(0.014) c	91.4(0.7) b
12 x 1	91.2(0.8) b	99.6(0.2) c	6.2(0.8) a,b	0.049(0.007) c	94.6(1.8) b
12 x 0.5	95.9(1.0) a	100.0(0.7) c	7.8(1.1) a	0.124(0.017) b	99.2(1.3) a

(b)

COMPARISON OF CHANGE IN DIAMETER					
Mould Dimensions (ϕ x h)	Initial relative light transmission percentage (t=0)	Final relative light transmission percentage (t=60)	Total strain % (S_e %)	Strain % Per unit mass (% S_o g ⁻¹)	Lower hardness percentage
4 x 2	89.5(1.6) c	107.7(2.4) a	4.1(0.7) b	0.145(0.025) b	85.6(1.1) c
8 x 2	87.9(0.6) c	101.4(0.9) b	6.4 (1.6) a,b	0.057(0.014) c	91.4(0.7) b
4 x 1	93.8(0.5) a	101.5(0.5) b	4.9(0.4) b	0.346(0.031) a	93.2(1.3) b
8 x 1	91.3(0.1) b	99.8(0.2) c	5.8(0.8) a,b	0.103(0.013) c	92.1(2.0) b
12 x 1	91.2(0.8) b	99.6(0.2) c	6.2(0.8) a,b	0.049(0.007) c	94.6(1.8) b

Table 3.3.2: Comparison of light transmission, strain and Vickers hardness data of each specimen group. The tables represent the comparison between thickness at constant diameter (a) and diameter at constant thickness (b). Similar letters within each column represents no significant differences.

(a)



(b)

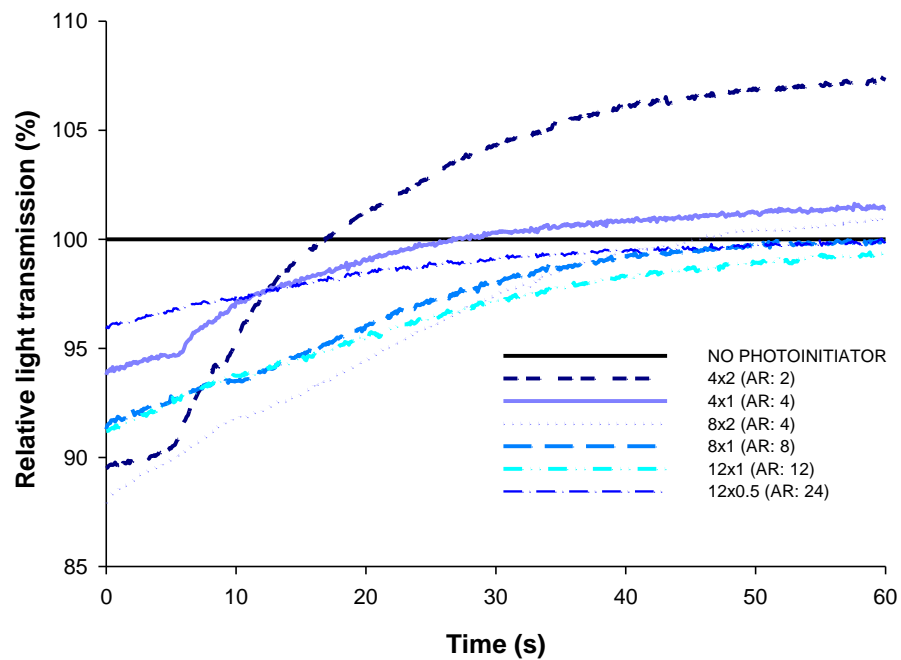


Figure 3.3.3: Relative light transmission percentage of resin contained in black nylon mould cavities having different aspect ratios ($n=3$) for the control group (a) and for curing resins (b), which are a percentage of their respective control group.

For the micro-hardness measurements, the one-way analysis of variance also revealed significant differences between cavity height and cavity diameter. Cavity height significantly reduced lower hardness percentage ($p < 0.05$) except for the 8 mm wide specimens, although a noticeable decrease in hardness was observed (Table 3.3.2 and Figure 3.3.4). Cavity diameter only significantly affected the percentage of lower surface cure for the 2 mm thick specimens ($p < 0.05$) but no significant difference was observed for the other specimen groups. The Vickers hardness value for the upper surface of each specimen remained consistent regardless of specimen geometry (Figure 3.3.4).

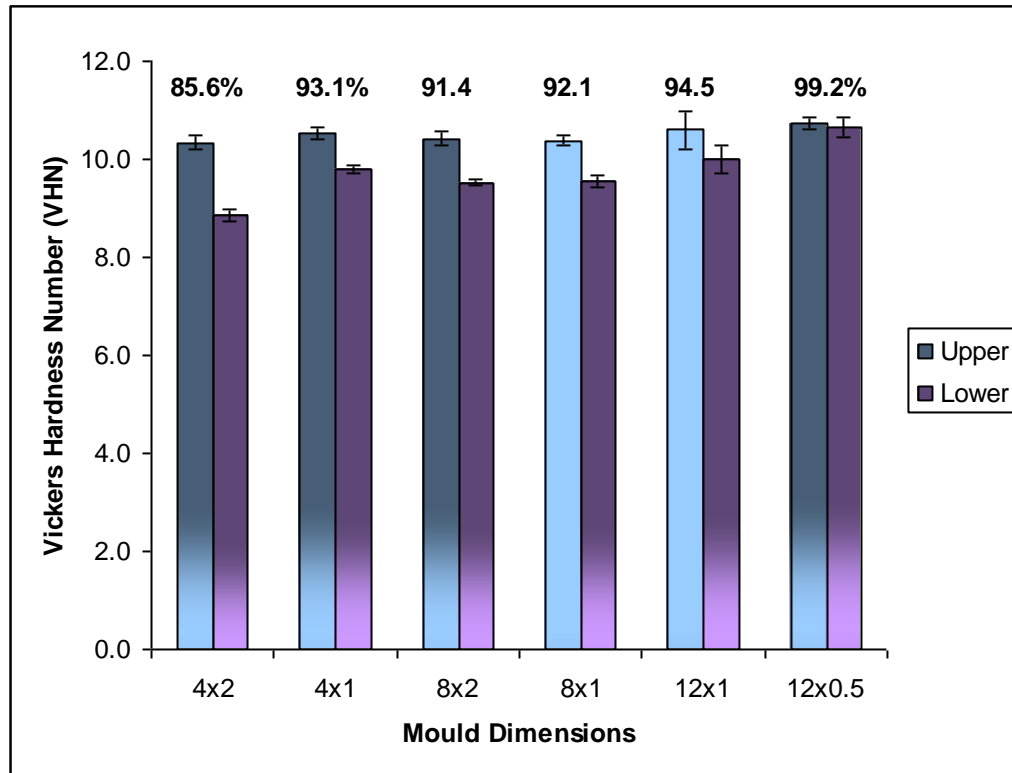
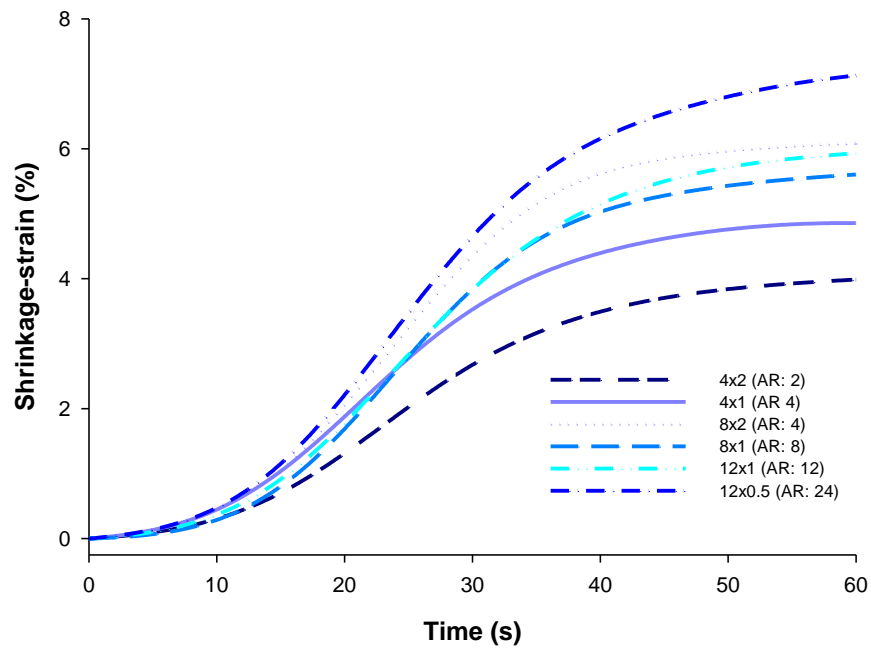


Figure 3.3.4: The Vickers hardness number of the upper and the lower surfaces of cured specimens. The error bars represent the standard deviations of the averages of three indents of each specimen (n=3).

Only the 1 mm thick specimens showed a significant effect with increasing diameter at constant height for total shrinkage strain, where total shrinkage strain was significantly higher for 8 and 12 mm specimens when compared to the 4 mm specimen. Although shrinkage strain increased with aspect ratio, no noticeable trend was observed for increasing height at constant diameter for the total shrinkage strain (Table 3.3.2a), but the effects may be separated by shrinkage strain per unit mass as the mass of the material played an influential role in determining the magnitude of the strain during polymerization (Table 3.3.1 and Table 3.3.2; Figure 3.3.5). Shrinkage strain per unit mass significantly decreased with increasing diameter at constant height ($p < 0.05$) for all the specimen groups although the effect was not significant between the 8 mm specimen group and the 12 mm specimen group. At constant diameter and increasing thickness, shrinkage strain per unit mass significantly decreased ($p < 0.05$) except for the 8 mm specimens, which showed a non-significant decreasing trend (Tables 3.3.2 and Figure 3.3.5).

(a)



(b)

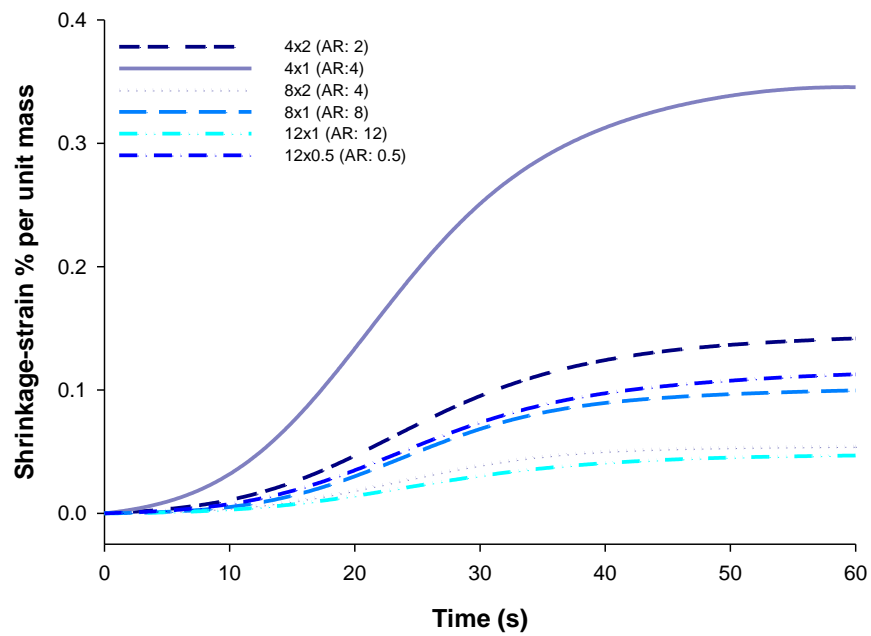


Figure 3.3.5: Shrinkage strain data for specimens cured in black nylon moulds; total measured shrinkage strain percentage (a) the shrinkage strain per unit mass (b).

Figure 3.3.6 represents the transmission profile and the digital image of a specimen bonded to its cavity walls in comparison to a non-bonded specimen. The difference in transmission profiles are obvious.

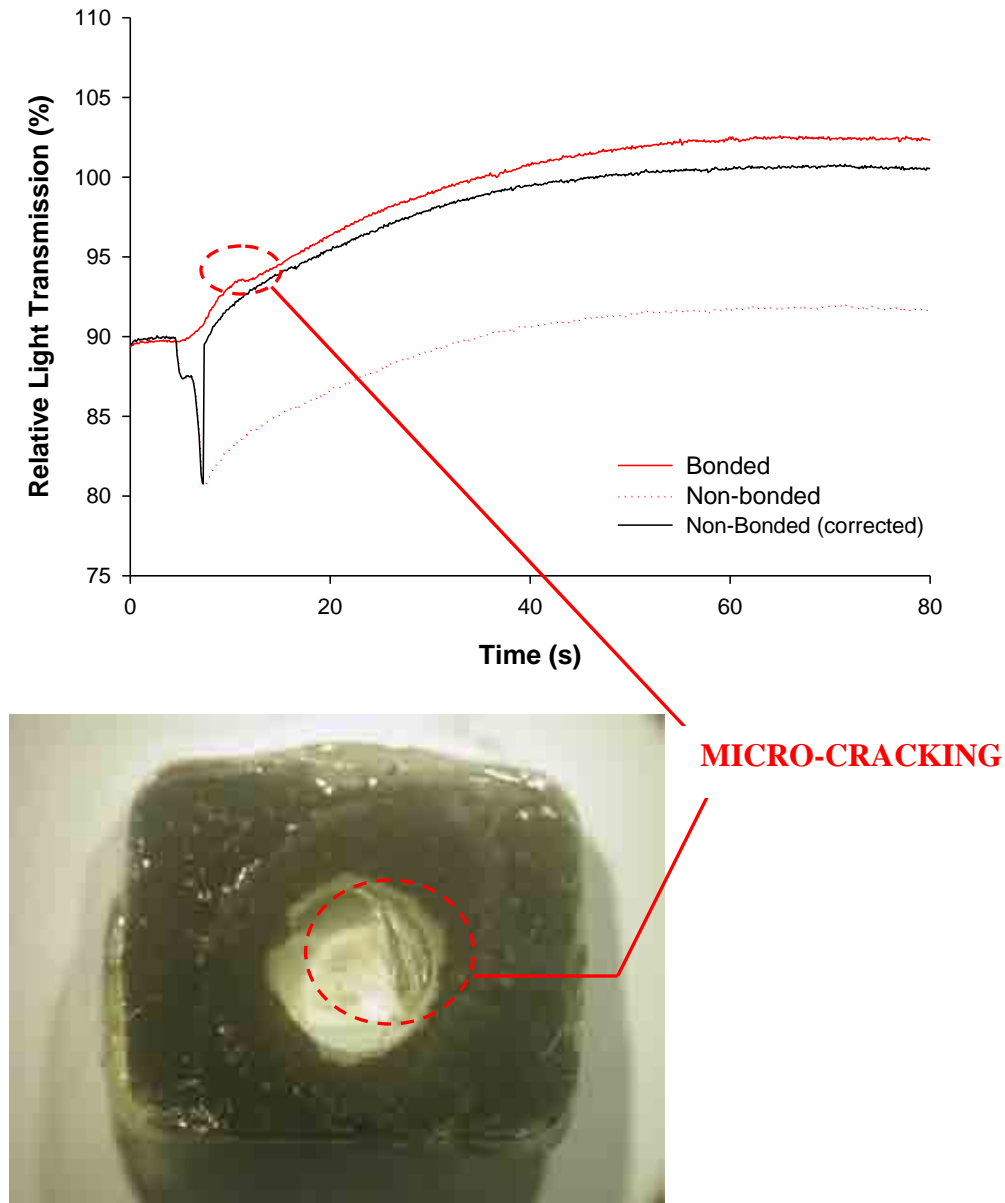


Figure 3.3.6: Light transmission through resins contained in fabricated ceramic disks which were either bonded or non-bonded (a). For the non bonded specimens, a decrease in transmission is indicative of de-bonding from the cavity walls. Representative digital image of micro-cracked specimens which were bonded to the cavity walls with Scotchbond adhesive (b). The visible micro-cracks were characterized by an audible cracking sound.

3.3.5 Discussion

The first hypothesis was accepted as mould cavity dimensions affected light transmission. The transmission of light is governed by Beer-Lamberts law (Equation 3.3.2) and as such it should be expected that only changes in specimen height (l) will affect light transmission. Therefore, it might have been expected that diametrical changes (i.e. increasing specimen diameter at constant height) would have had no affect on initial light transmission as light is only detected at the center of each specimen. However, this was not observed and interestingly at similar thickness, wider specimens showed less initial light transmission, significantly in the case of 1 mm thick specimens. This may seem unexpected as a previous study reported that narrow specimens at similar thickness exhibited reduced light transmission as a greater amount of light was attenuated due to absorption by the cavity walls of the black nylotron moulds (Harrington et al., 1993). Interestingly in the present investigation, this was only observed for the control groups (without photoinitiator), which showed improved light transmission with increasing diameter at similar thickness. The reduced initial light transmission in the curing resins for wider specimens may be related to the greater volume of resin (and increased photoinitiator molecules within the volume), which may increase the absorption and attenuation of light. The peripheral absorption and attenuation may be related to the fact that greater absorption occurs by the photoinitiator than the cavity walls of the mould, which resulted in the reduced transmission. Such an effect may be reduced in relatively wide specimens as the CC3-UV cosine corrector is only capable of registering directly transmitted light normal to the specimen surface with a field of view of 180° , which may account for the non-significant decrease in initial light transmission for the 12x1 mm specimen when compared to the 8x1 mm specimen (Table 3.3.2b). Furthermore, typical halogen lights such as the XL2500 have a Gaussian distribution of light irradiance across the tip diameter (Price et al., 2010) and therefore any edge effects on light transmission for wider specimens will be reduced which may also account for the non-significant data. Nonetheless, the highest initial light transmission was seen for the thinnest and widest specimen whereas the thickest and widest specimens showed the lowest initial light transmission (Tables 3.3.2 and Figure 3.3.3). This might be directly related to Beer-

Lamberts law as absorption increases proportionally with optical pathlength (i.e. thickness) and concentration of absorbing species (i.e. photoinitiator molecules as the given volume increases), which would reduce the intensity of the propagating light.

During polymerisation, the light transmission is also governed by the same law, and is affected by dynamic changes in pathlength and photoinitiator properties (Section 3.1 and 3.2). Consequently, the observed light transmission traces may not represent the ‘true’ transmission as correcting for such changes become complicated. Therefore, part of the observed increase in light transmission traces may be due to specimen shrinkage and photoinitiator decomposition (Chen et al., 2007). However, the effect of sample deformation to increase light transmission is also a possibility where a ‘lens effect’ due to shrinkage causes the surface of the sample to bow. Convex lenses typically converge light into a focused beam and will tend to improve the intensity of light detected by the cosine corrector. In contrast to this, concave lenses typically ‘diverge’ the light and spread the beam, which will reduce the intensity detected by the cosine corrector (Figure 3.3.7). It is interesting to note that Asmussen and Peutzfeldt (1999) reported finding a convexity of the sample surface closest to the light and concavity of the surface furthest away from the light for composites cured in 6 mm diameter moulds having thickness of up to 6 mm, which indicates a convex-concave scenario (Figure 3.3.7). Therefore the increased light transmission may well be due to the formation of a convex-concave lens, which will tend to focus the beam onto the cosine corrector. The deformation may be greater in narrow-thick specimens (i.e. 4x2 mm) due to shrinkage vectors occurring axially as well as non-axially as a consequence of the lower aspect ratio (Watts et al., 2003), which may account for the substantially high increase in light transmission (>100%). However, without surface profilometry (ideally 3-D), this may not be confirmed and remains speculative. Nonetheless, improved light transmission is associated with improved degree of conversion and improved mechanical properties such as strength, elastic modulus and hardness (Price et al., 2002; Palin et al., 2008), although some physical properties such as shrinkage may increase for conventional resins with improved degree of conversion (Braga et al., 2005).

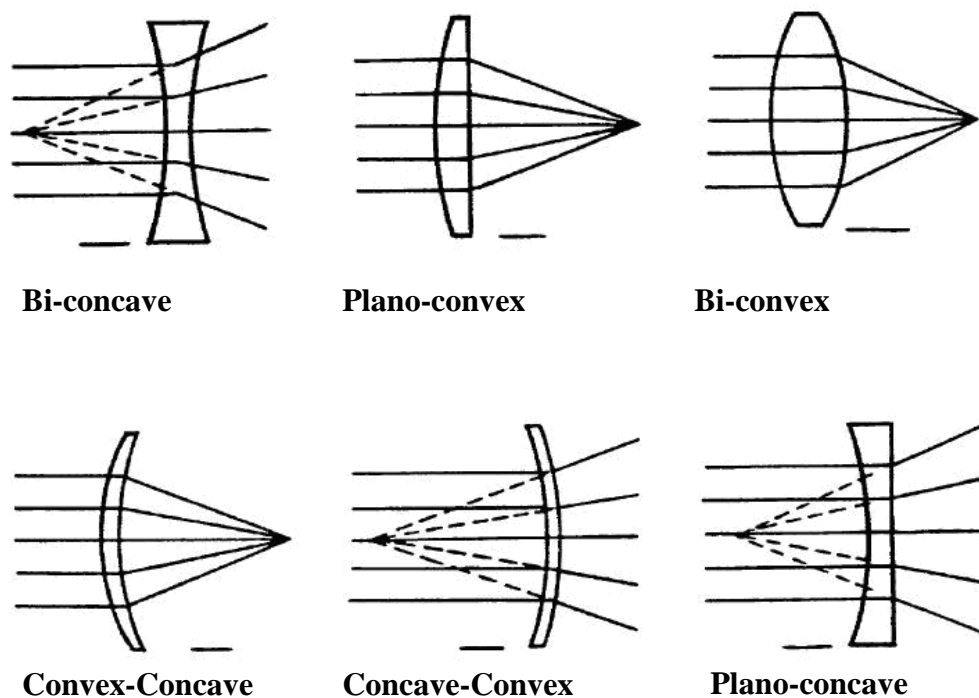


Figure 3.3.7: Lens characterization and the converging/diverging of light travelling from left to right.

The second hypothesis must be rejected as the specimen that showed the greatest improvement in light transmission (Table 3.3.2 and Figure 3.3.3) showed inferior lower surface hardness relative to its upper surface (Figure 3.3.4) in comparison to the other specimen groups. The hardness values for the upper surfaces remained comparable between the different specimen geometries whereas the hardness of the lower surfaces was dependent on the specimen geometry and the amount of light received initially at the lower surface. This might be supported by previous findings where a logarithmic correlation between intensity and depth of cure and a linear relationship between depth of cure and irradiating distance were reported (Aravamudhan et al., 2006). Not surprisingly, the 12x0.5 mm specimens showed significantly higher lower surface hardness in comparison to the other specimens, which is due to the relatively thin specimen height (0.5 mm) allowing equivalent cure on either surface due to the 96-100% light transmission throughout cure (Figure 3.3.3 and Figure 3.3.4). Interestingly, the specimen

group that showed the highest final light transmission had the lowest lower surface hardness. However, the substantially high increase in light transmission ($>100\%$; Figure 3.3.3) for the lowest aspect ratio mould may not reflect 'true' transmission as previously mentioned. The lower surface hardness for the 4 mm and 12 mm diameter moulds were significantly lower for the thickest mould but for the 8 mm diameter moulds, there was no significant difference in lower surface hardness percentage with height variation. This was not expected as the initial light transmission is significantly lower for the thicker of the two. This may be related to adequate light transmission through the unfilled resins, which may not occur in 4 mm diameter moulds due to absorption by the cavity walls.

The dimensions of the cavity will play a significant role in the viscous flow of the specimen during polymerisation and any developing stresses and strain, which may affect light transmission. During polymerisation, the resin transforms from a viscous resin to a viscoelastic phase (de Gee and Davidson, 1981). This transformation is characterized by the so called 'gel-point'. Prior to the gel-point, any developing stress can be compensated by the intrinsic flow of the resin (Sakaguchi et al., 2004). The shrinkage vectors (which are dependent upon height and diameter) will govern how light transmission is affected. When uni-axial shrinkage dominates and the shrinkage leads to a reduction in the height, it might be expected that the outcome would be improved light transmission and associated improvement in curing extent. For cavities with low aspect ratios, uni-axial shrinkage dominates with some non-axial shrinkage occurring due to the low aspect ratio (Watts et al., 2003). The constraint caused in such specimens may cause specimen deformation and a non-uniform transmission profile, which improves transmission significantly during polymerisation. The sudden increase in light transmission for the 4 mm diameter moulds may be due to the high shrinkage strain imposed by the low mass and the constraint of the specimen (Figure 3.3.5). A sharp increase in transmission is expected at the point where the resin can no longer relax due to pre-gel flow. Interestingly, transmission above hundred percent is observed which may be due to the shrinkage vectors dominating axially (with some non-axial contribution) and the constraint causing the samples to become convex on the irradiated side and concave on the non-irradiated side as previously mentioned (Figure 3.3.7). The increase in diameter from 4 mm to 8 mm will increase the aspect ratio, which will reduce the constraint,

allowing both axial and non axial shrinkage (i.e. bi-axial). Therefore, non-axial shrinkage will contribute more to the total shrinkage and the deformation should be reduced (i.e. less convexity and concavity) as the sample shrinks more freely. The mass of the specimen will determine the magnitude of the shrinkage strain. At similar aspect ratio, the strain per unit mass is significantly greater ($p < 0.05$) for the smaller specimen (4x1 mm) when compared to a larger one (8x2 mm), whilst the total strain increases with mass although non-significantly (Table 3.3.2 and Figure 3.3.5). As shrinkage occurs bi-axially in both 8 mm diameter moulds, the increase in light transmission occurs at a more steady pace than the 4 mm diameter specimens and the bi-axial shrinkage does not allow for an increase in transmission above hundred percent (Figure 3.3.3). Then, for the mould having geometry of 12x1 mm, the aspect ratio increases from 8 (for the 8x1 mm mould) to 12. As a consequence of the larger diameter, a greater mass of resin is introduced. This greater mass allows for a lower shrinkage strain per unit mass although the total shrinkage strain increases when compared to the 8x1 mm mould (Table 3.3.2 and Figure 3.3.5). Consequently, as constraint decreases and radial shrinkage (i.e. shrinkage away from the cavity walls) starts to dominate, a slower increase in light transmission is expected (Figure 3.3.3). For the highest aspect ratio specimen, the relatively thin specimen and low mass when compared to the specimen having dimensions of 12x1 mm, will only allow radial shrinkage. The magnitude of the shrinkage strain is high as the opposing glass covers will cause greater constraint than the constraint imposed by the cavity walls of the mould as the aspect ratio is high (Figure 3.3.5a) but the shrinkage strain per unit mass remains comparable to that of a specimen with similar mass (8x1 mm; Figure 3.3.5b). Radial shrinkage in the relatively thin specimen would only produce little changes in transmission, and it may be speculated that the little change in light transmission observed for the 12x0.5 mm mould may be mainly due to photoinitiator decomposition (Chen et al., 2007).

For low aspect ratio moulds, constraint is high and the bond between the cavity walls and the opposing surfaces might be compromised. When the shrinkage becomes restricted stress is generated. Stress generated after the gel-point competes with interfacial adhesion and can not be easily compensated. If this stress cannot be dissipated or relieved the sample may crack and introduce air-pockets which will affect light

transmission, as indicated by the inflection of the ‘bonded’ transmission profile and digital image in Figure 3.3.6, respectively. Free-shrinkage (by means of separation or due to low constraint) will allow stress dissipation or relief. Free shrinkage will also reduce the ‘lens’ effect as the specimen will shrink freely both axially and non-axially. Therefore, the lower transmission of the corrected ‘non-bonded’ transmission in comparison to the ‘bonded’ trace may support the greater deformation in the 4x2 mm specimens. However, in the ceramic moulds, the increase above 100% was lower than that observed in the nylon moulds (Figure 3.3.3 and Figure 3.3.6). This maybe related to the fact that acid etching ensured mechanical adhesion to both cavity walls and opposing surfaces and therefore deformation was almost impossible causing stress to build up and samples to crack. Nonetheless, for separated samples, de-bonding from the ceramic cavity walls and opposing surfaces resulted in a sharp decrease in light transmission following gelation of the curing resin as larger air gaps (than for sample micro-cracking) are introduced between the surfaces (Figure 3.3.6). This may be due to the much lower refractive index (RI) of air (RI~1) compared to that of resins cured or un-cured (RI~1.50-1.55) (Section 3.2). De-bonding will allow light to travel down the sides of the sample (through air) and through air gaps between the microscope slides and samples. The sudden mismatch in RI between glass, air and curing resin will increase scattering and account for the sudden decrease in light transmission that was observed (Figure 3.3.6). From the current investigation, it is evident that the cavity dimension influence how light is transmitted through curing specimens and the constraint controls strain vectors and magnitude, which may also affect light transmission. Further work is required to fully understand this phenomenon at similar material mass and how other material properties are affected by such changes in light transmission. In addition to correcting for absorption and shrinkage, surface profilometry and more specimen geometries should be used which may help confirm these speculations.

3.3.6 Conclusion

Specimen geometry and constraint affected light transmission throughout polymerization of dental resins. The light transmission was dictated by how the shrinkage vectors and strain developed during cure. Unexpected light intensity variations may occur for low configuration factors, which may increase or decrease curing extent of resin based restorations and ultimately compromise the success and longevity of the restoration. Characterizing fundamental aspects of photo-induced polymerization of dental resins such light transmission, with respect to shrinkage and shrinkage strain will provide a greater insight into the shrinkage behaviour of such materials and this will ultimately aid to improve longevity of such restorations. Furthermore, improving light transmission through depth may consequently improve degree of conversion.

3.4 Summary

The aim of this chapter was to investigate the changes in optical properties of curing dental resins and introduce a novel technique to dynamically measure refractive index change and physical thickness change of photoactive dental resins. Changes in optical properties in RBCs occur as light is absorbed, attenuated, scattered and refracted through the specimen. Not only does the material composition affect such properties, i.e. the resin matrix composition, filler type and percentage, photoinitiator type and concentration and the addition of pigments and dyes, the conversion from monomer to polymer is also expected to affect optical properties.

Along with dyes and pigments used to create varying shades of resin composite materials, the photoinitiator is a strong absorber of light in photoactive dental resin composites, which decomposes during activation therefore decreasing absorption and increasing the amount of light transmitted through the specimen. Alternative colourless photoinitators are mainly used for aesthetic reasons as camphoroquinone has a distinct yellow tint due to its chromophore group which is partly unbleachable (Jakubiak et al., 2003). As well as improved aesthetics, non pigmented photoinitators have been suggested to have additional benefits such as reduced curing time (Section 2.3) due to their higher efficiency. However, the use of such initiators is expected to have a significant influence on the optical and physical properties of the curing resin composite. Firstly, inferior depth of cure of Lucirin TPO-based resin composites was reported in Section 2.3 as a consequence of low wavelength scattering of light and high molar absorptivity causing in-efficient light transmission at depth. Secondly, such a high molar absorptivity is expected to affect reaction kinetics and increase the rate of polymerisation. Section 3.1 aimed to investigate some optical properties of resins made with camphoroquinone and Lucirin TPO with respect to reaction kinetics. Photoinitiator decomposition during cure is expected to increase light transmission proportionally according to Beer-Lamberts law (Chen et al., 2007). However, the inverse absorption observed for Lucirin TPO resins maybe related to monomer discolouration associated with high reaction exotherms by rapid rates of polymerisation. As such, the initially colourless resins become stained and an increase in yellowing was observed, verified by

the increase in b-value of the CIELAB co-ordinates. Conversely, for camphoroquinone resins, the initially yellow resins bleach upon irradiation and a decrease in yellowing is observed (decrease in b-value). Any discolouration during irradiation may potentially reduce light transmission through depth and affect depth of cure, even for photocurable resins that utilise so-called, “colourless” photoinitiators such as Lucirin-TPO. Characterising photoinitiator properties dynamically will ultimately improve photoinitiator chemistry in such materials.

The optical properties of resin-based composites are essential to optimise light transport through the material and to aid curing depth. Although refractive index of dental resins are well characterised statically, the change during cure is not well understood. Section 3.2 aimed to demonstrate a novel technique whereby refractive index can be measured during cure. The ability to measure refractive index change of curing photoactive resins will allow for the design of composite formulations, which exhibit optimum evolution of translucency (as the RI of the curing resin approaches that of the filler) and enable precise control of aesthetic quality, light transport through material bulk and curing depth.

The shrinkage of dental RBCs is also an important consideration. When composites shrink, the affect on light transmission may be predicted by Beer-Lamberts law, where a decrease in optical path length will increase light transmission. Not only is the shrinkage affected by material composition, i.e. filler type and percentage and resin matrix composition, specimen geometry is likely to significantly affect the shrinkage vectors of the curing specimen, which may affect light transmission. As such, in Section 3.3 it was observed that low aspect ratio moulds may result in unexpected light transmission profiles due to sample deformation caused by the relatively high constraint.

The current Chapter demonstrated for the first time the ability to simultaneously measure refractive index change and physical thickness change in photoactive dental resins during polymerisation and reported the effect of monomer composition and specimen constraint on optical and physical properties (Section 3.2 and 3.3). Such information may be useful in order to optimise the resin matrix / filler interaction to improve light transmission through depth, with the aim of improved depth of cure and clinical success of such restorations.

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CHAPTER 4

Recommendations for Further Work

4.0 Recommendations for Further Work

In order to compare the anomalies of the exposure reciprocity law, commercial and experimental materials were tested. Although static rheology measurements provide some insight into the mobility of propagating species and reaction kinetics, the increase in density increases the viscosity during cure. Consequently, dynamic rheological changes will differ remarkably according to the reactivity of the monomer, the filler percentage and the photoinitiator chemistry throughout and proceeding light irradiation. Therefore dynamic rheological measurements may give a better insight into mobility of propagating species. Furthermore dielectric spectroscopy may also be used to measure mobility within a curing resin directly through Stoke's law where ionic conductivity is inversely proportional to the viscosity, in which the ionic species is considered to be rigid spheres moving through the liquid monomers. As such, the utilisation of real time rheology measurements and dielectric spectroscopy would afford a more detailed insight into the differences in polymerisation kinetics and failures in the exposure reciprocity law and should be the subject of further studies.

Exposure reciprocity was verified even at low viscosity with the use of Lucirin TPO, however inferior cure depths were reported. Clinically, this may be problematic as an increased number of technique sensitive incremental steps will be required to fill a relatively large cavity, which may also increase procedural time rather than decrease it. The increased scattering at low wavelength may be responsible for the inferior cure depths. Optimising the filler particle and resin matrix according to photoinitiator absorption spectrum may reduce scattering and aid curing at depth. In addition to this, using camphoroquinone or another photoinitiator which absorbs at higher wavelengths synergistically with Lucirin TPO will allow light at higher wavelengths to penetrate deeper with less scattering and may improve curing depth. Therefore it is proposed that further investigations are conducted with different resin and filler constituents and a variety of photoinitiators such as camphoroquinone, phenyl propanedione and benzyl used synergistically with Lucirin TPO in order to try and improve cure depths.

Much of the current investigation has focussed on the applicability of the exposure reciprocity law based on degree of conversion data. However, degree of conversion data should not be considered the gold-standard of material properties. The polymer chain length and cross linking may also be influenced by irradiation time and intensity and thus even at similar degree of conversion reported in the current investigation, the mechanical properties such as flexural strength, flexural modulus, water sorption and cure depths will be affected. Therefore, in order to better understand how the reduction in curing time will influence clinically important parameters, it is proposed that the materials tested in the current investigation are tested for their performance under mechanical testing regimes, such as dynamic mechanical thermal analysis which may be useful for real time development of mechanical properties or static load-to-failure tests such as three-point bend to assess flexural strength and modulus.

The introduction of a novel low coherence interferometer used to assess optical and physical change dynamically in the present study may provide valuable information on optical and physical change in resin-based composites whilst under cure. The development of such a technique may allow for the design of composite formulations, which exhibit optimum evolution of translucency to enable precise control of aesthetic quality and light transport through material bulk and curing depth. Consequently, it is proposed that further work utilising model composites with different resin matrix constituents/ratios and filler type/percentage be tested for a more detailed insight into the evolution of translucency and its affect on depth of cure.

Additionally, although 2-D measurement of refractive index and physical thickness change provided useful information on the optical properties and shrinkage behaviour of experimental resins, such changes are not uniform throughout the specimen. Precise control of optical properties such as refractive index will optimise light transmission at depth and enable greater depths of cure. However, to fully understand bulk optical change, a spatial, 3-D measurement of refractive index and physical thickness change would be more appropriate. Modification of the bespoke software used to control the measurement parameters in the current study to allow axial measurement during cure will enable precise mapping of the sample under measurement. The mapping

of which will further optimise the design of composites to allow optimal light transmission at depth by better design of composite formulations.

Light transmission is an important factor for degree of conversion and depth of cure. As refractive index changes and approaches that of the filler, interfacial scattering decreases and light transmission increases. Therefore, light transmission is an important measurement parameter to optimise curing of resin-based composites. However, as optical measurements are sensitive to the experimental set-up, it seems more appropriate to measure them simultaneously. Thus, developing a device which allows the simultaneous measurement of refractive index/shrinkage and light transmission will reduce experimental errors and allow a better analysis of the data obtained.

Finally, to truly optimise the experimental set up and procedure, degree of conversion and temperature measurements should also be made simultaneously. Experimentally, this is problematic as there will be many probes and sensors to operate simultaneously. Consequently, the ‘one-hit’ measurement device seems almost impossible to construct for the measurement of multiple parameters in a relatively small sample. However, both the low coherence interferometer and the FT-IR spectrometer are based on the Michelson interferometer and the broadband light source used in the low coherence interferometer has a centre wavelength of 1624 nm, which corresponds with the vinyl carbon double bond absorption wavenumber at 6164 cm^{-1} (as the wavenumber is given by the reciprocal of wavelength). Consequently, it is possible to apply a Fourier Transform function on the low coherence interferometer data to obtain the absorption spectra for measuring conversion and curing kinetics. This would reduce number of probes and sensors in the ‘one-hit’ measuring device and may allow for real time simultaneous measurements of refractive index, shrinkage, degree of conversion, temperature and light transmission which will ultimately aid material development by providing a greater insight into the curing behaviour of resin based composites.

As mentioned previously, 3-D mapping of dental resins and dental resin composites may optimise the design of composites to allow optimal light transmission at depth by better design of composite formulations. However, such mapping may also be used to validate the ‘lens’ effect theory discussed in Section 3.3.3. Currently, 2-D mapping may be used to map the thickness across the specimen on one axis but the data

remains limited as it only indicates the physical thickness and not how the sample shrinks (i.e. the type of lens formed). The most appropriate currently available solution to validating the 'lens' effect theory may well be profilometry, which will directly indicate how the samples shrink and their resultant geometry.

APPENDIX

APPENDIX

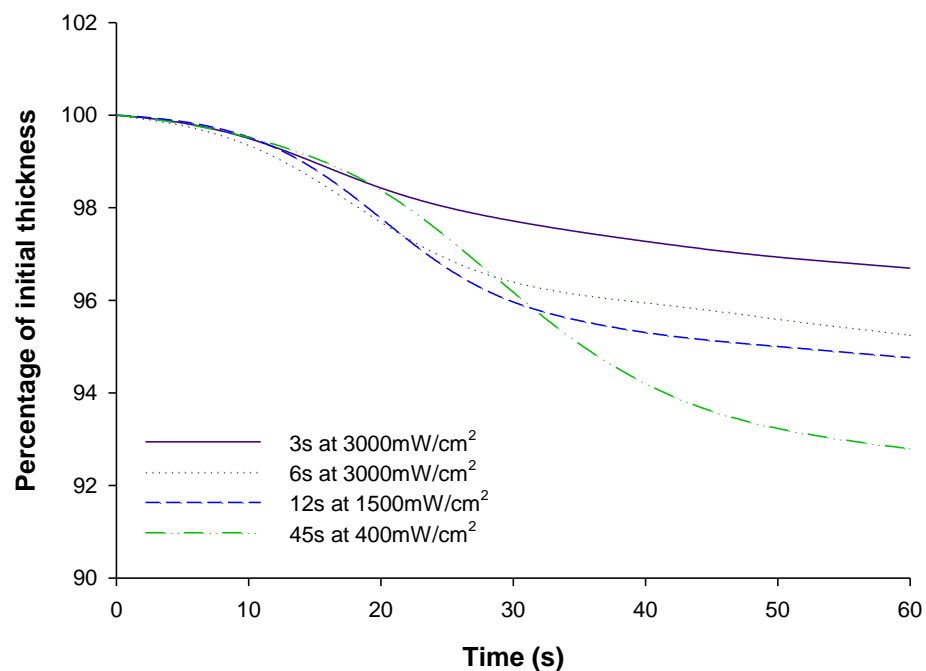
Appendix 1 Physical thickness change of Camphoroquinone resins and Lucirin TPO resins.

Camphoroquinone and Lucirin TPO resins were tested to verify that no expansion occurred during irradiation and reject the possibility that the initial increase in absorbance for Lucirin TPO resins was due to expansion by the high reaction exotherm (Section 3.1). The resins were tested with a Swiss Master halogen Light (EMS, Switzerland) under the hypothesis:

- i) Lucirin TPO resins will show initial expansion due to the high reaction exotherm whereas camphoroquinone resins will not.
- ii) The expansion would be greater for high irradiance protocols as rates of polymerisation and exotherm are expected to increase.

Figure A1 represents the physical thickness change of camphoroquinone resins and Lucirin TPO resins measured using the low coherence interferometry method described in Section 3.2. It is evident from the charts that no expansion occurred during irradiation, even for the highest irradiance protocols. Interestingly however, the highest irradiance protocol showed the greatest shrinkage for Lucirin TPO resins and an opposite trend was observed for camphoroquinone resins. Consequently, both hypotheses must be rejected and the possibility of increased absorption due expansion in Lucirin TPO resins must be ruled out.

(a)



(b)

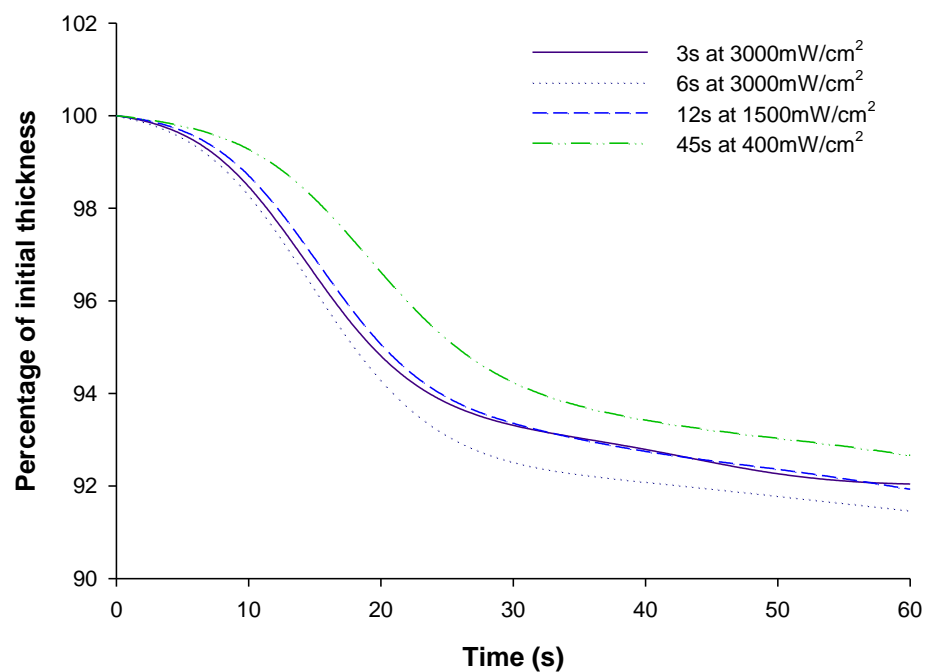


Figure A1: The physical thickness change of unfilled resins containing the low photoinitiator concentration (0.0134 mol/dm^3) for camphoroquinone (a) and Lucirin TPO (b) cured with a Swiss Master Halogen light at similar radiant exposure (18 J/cm^2).